Contact Mechanics and Elements of Tribology

Lecture 3.
Contact and Mechanics of Materials

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@ Centre des Matériaux (& virtually)
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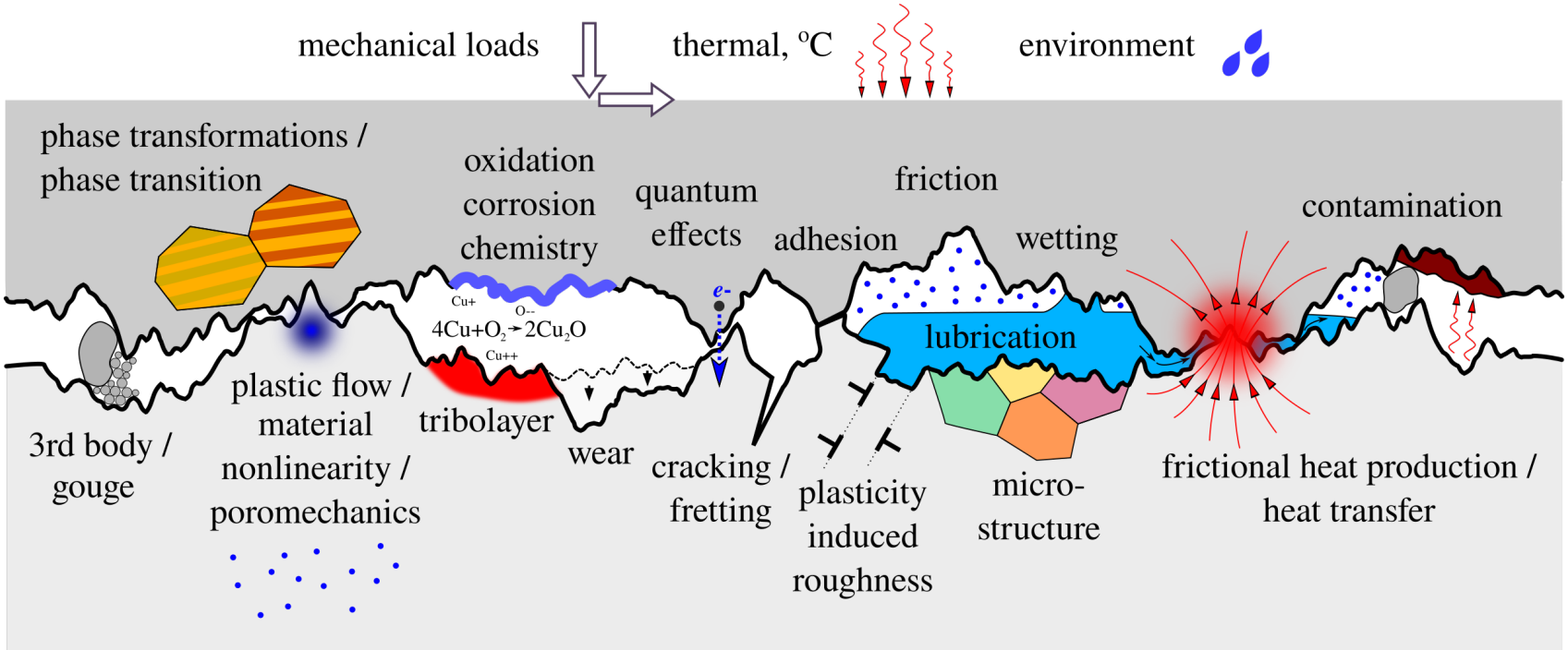
Interface physics



Tribological properties stem from this complex multi-scale physics of interfaces and near-surface layers[1]

[1] Vakis, Yastrebov, Scheibert, Nicola, Dini, Minfray,Almqvist, Paggi,Lee, Limbert, Molinari, Anciaux, Aghababaei, Echeverri Restrepo, Papangelo, Cammarata, Nicolini, Putignano, Carbone, Stupkiewicz, Lengiewicz, Costagliola, Bosia, Guarino, Pugno, Müser, Ciavarella, 2018. **Modeling and simulation in tribology across scales: An overview**. Tribology International, 125:169-199 (2018).

Outline

- Elastic-plastic contact
- Viscoelastic contact
- Frictional strengthening or time-dependent contact

Plasticity

Onset of plastic yielding

Hertz contact: body of revolution

- Onset of plasticity for pressure

$$p_Y = 1.6\sigma_Y$$

- Associated force

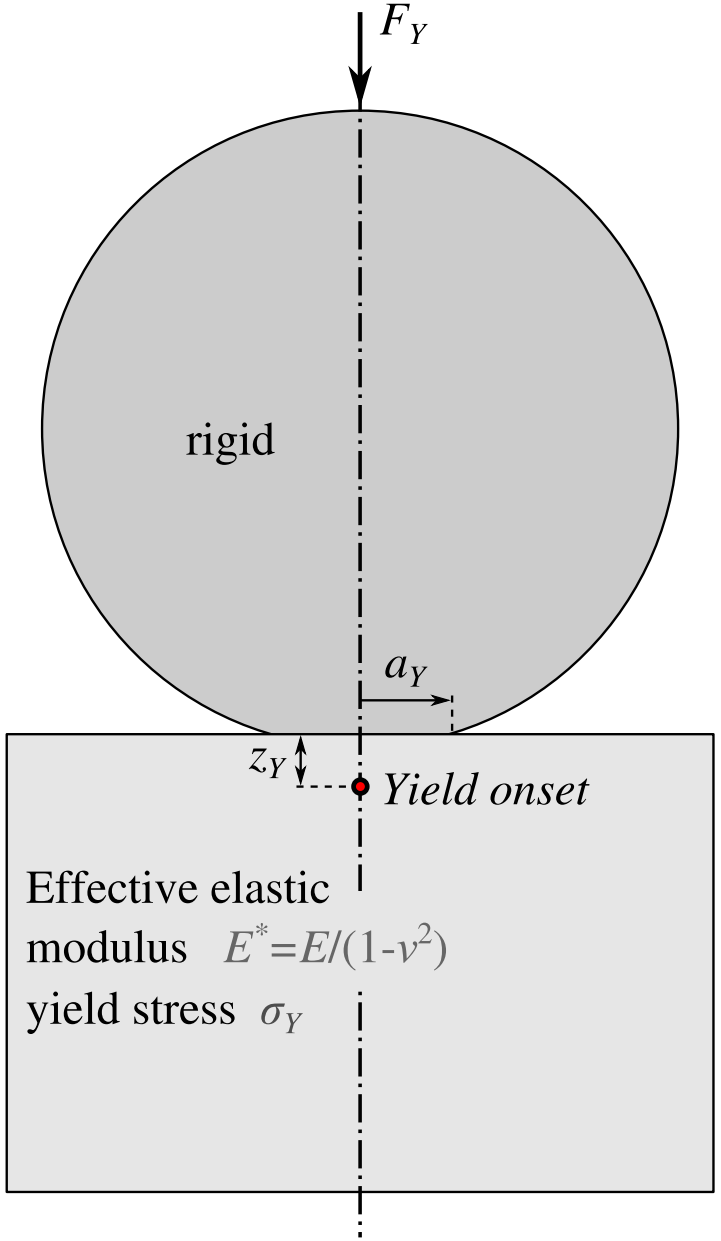
$$F_Y = \frac{1.6^3 \pi^3 R^2}{6} \left(\frac{\sigma_Y}{E^*}\right)^2 \sigma_Y$$

- Associated contact radius

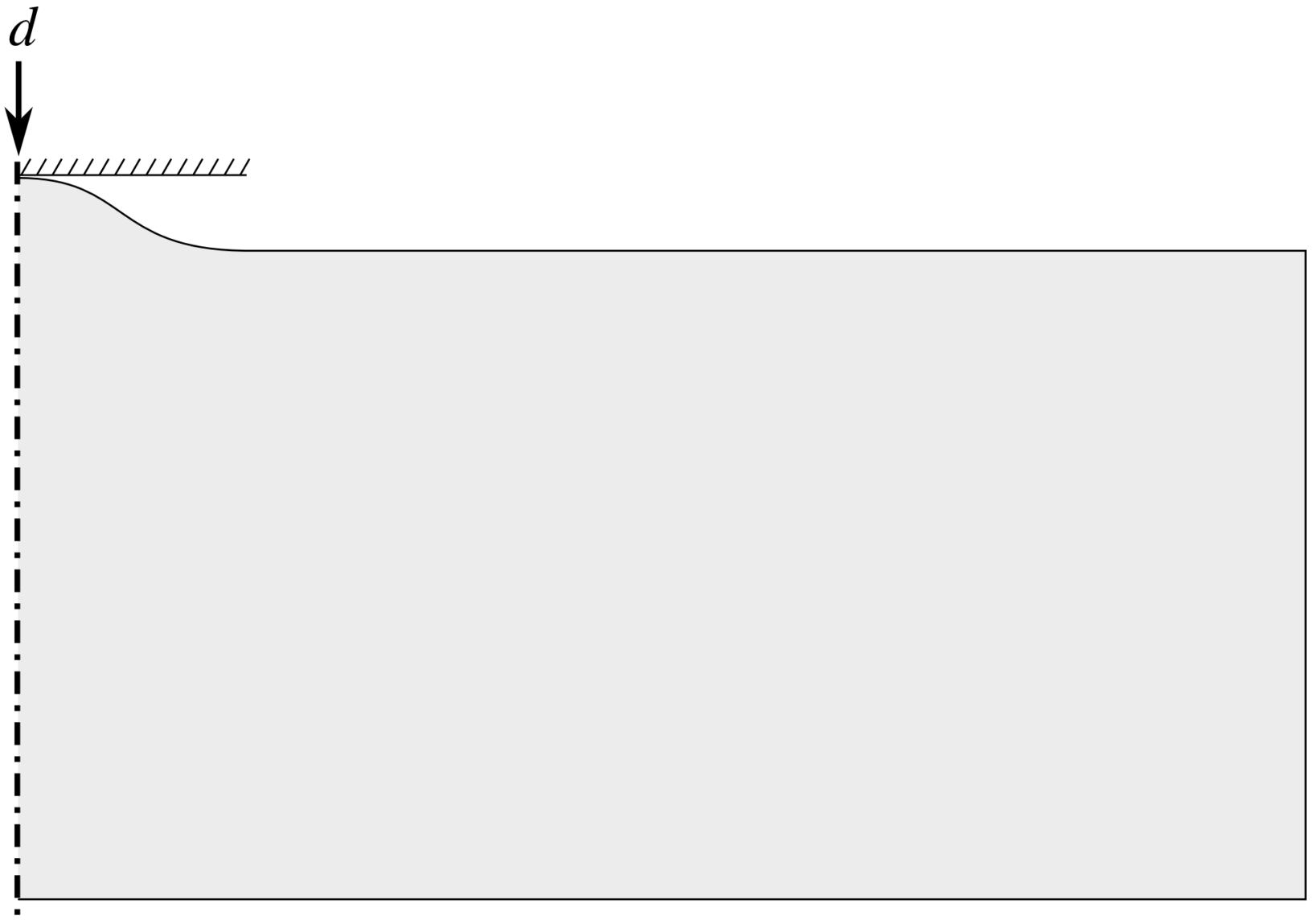
$$a_Y = \frac{1.6\pi R}{2} \frac{\sigma_Y}{E^*}$$

- Plastic flow starts at the depth

$$z_Y \approx 1.21 R \frac{\sigma_Y}{E^*}$$



Elasto-plastic transition in contact



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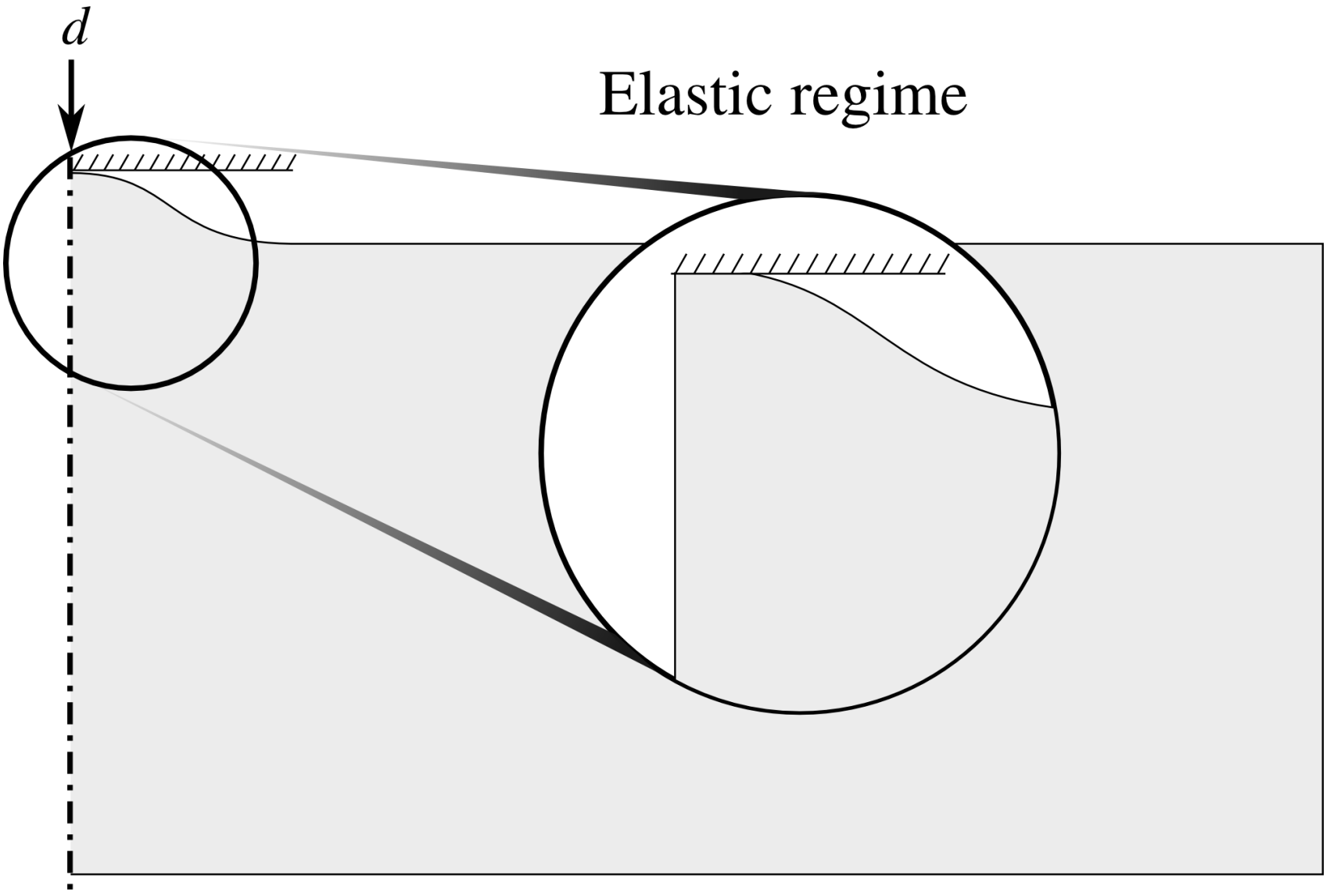

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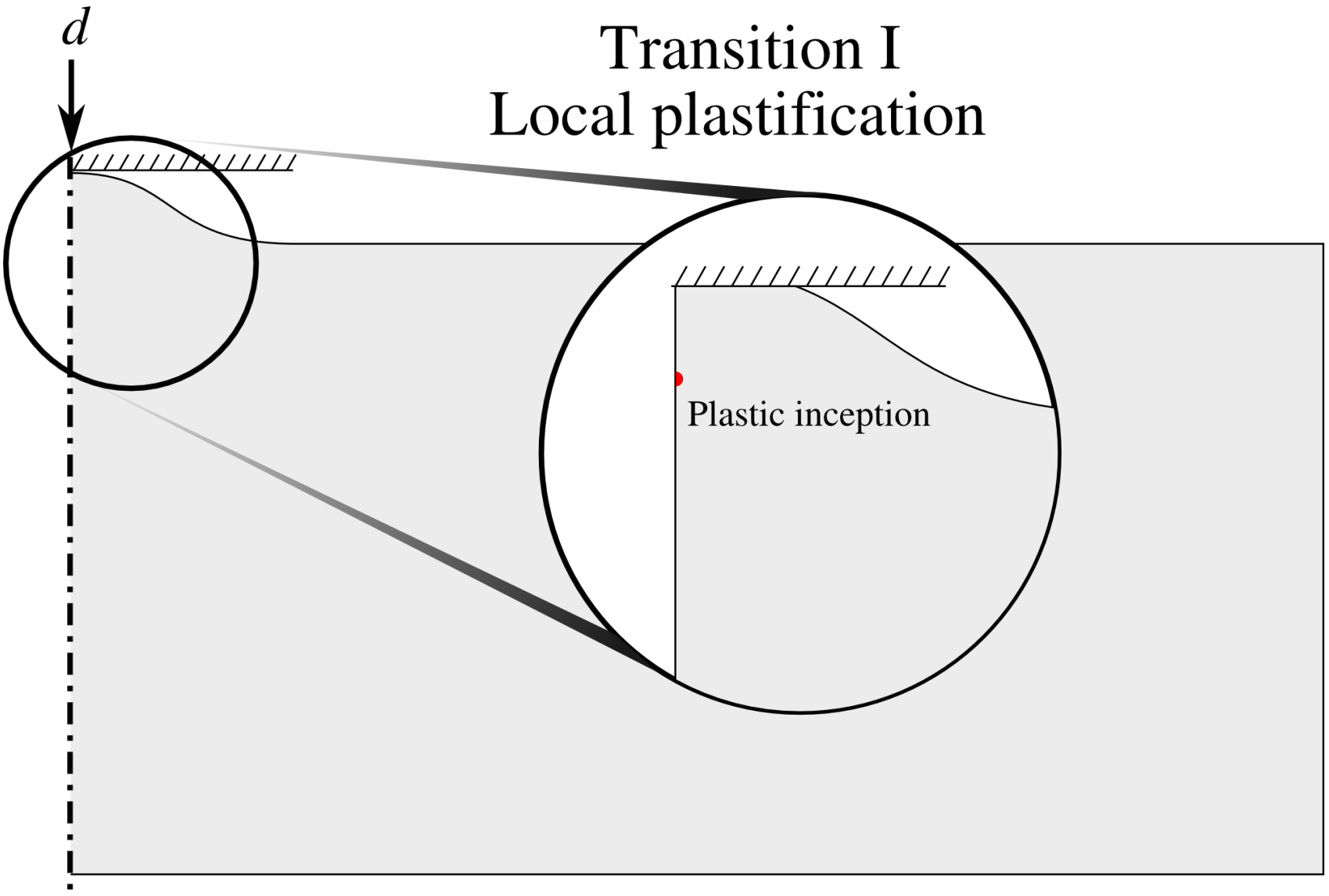

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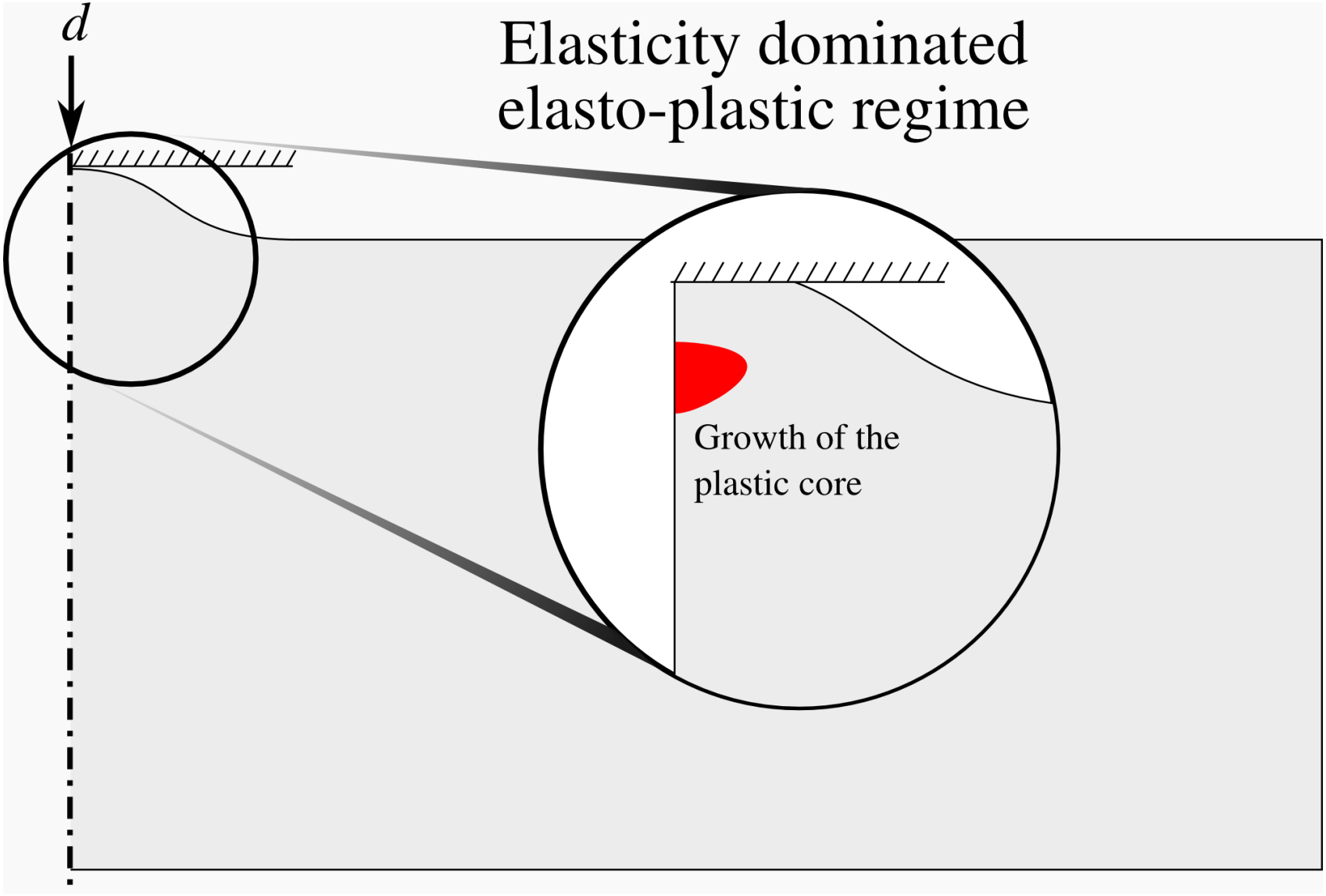

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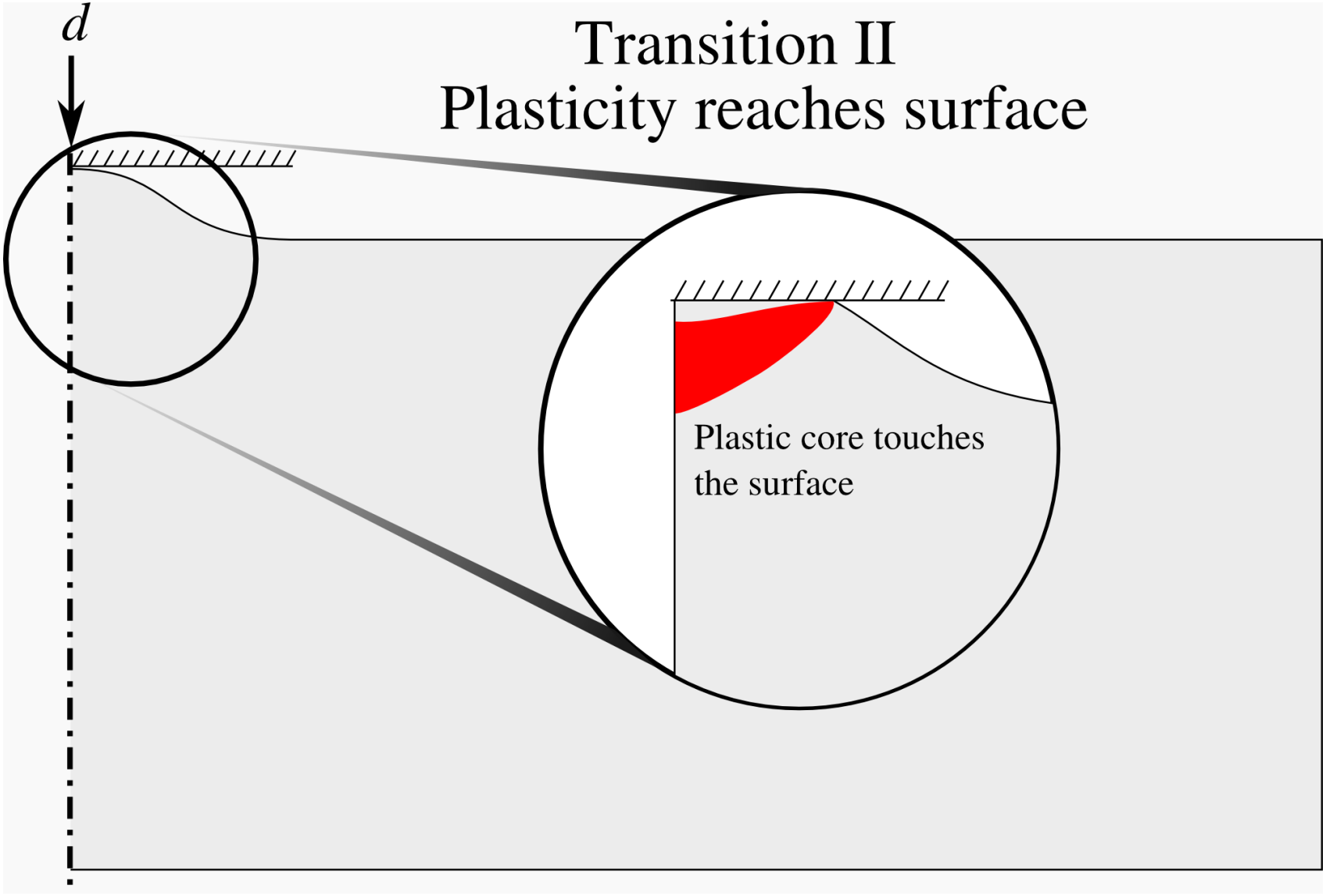

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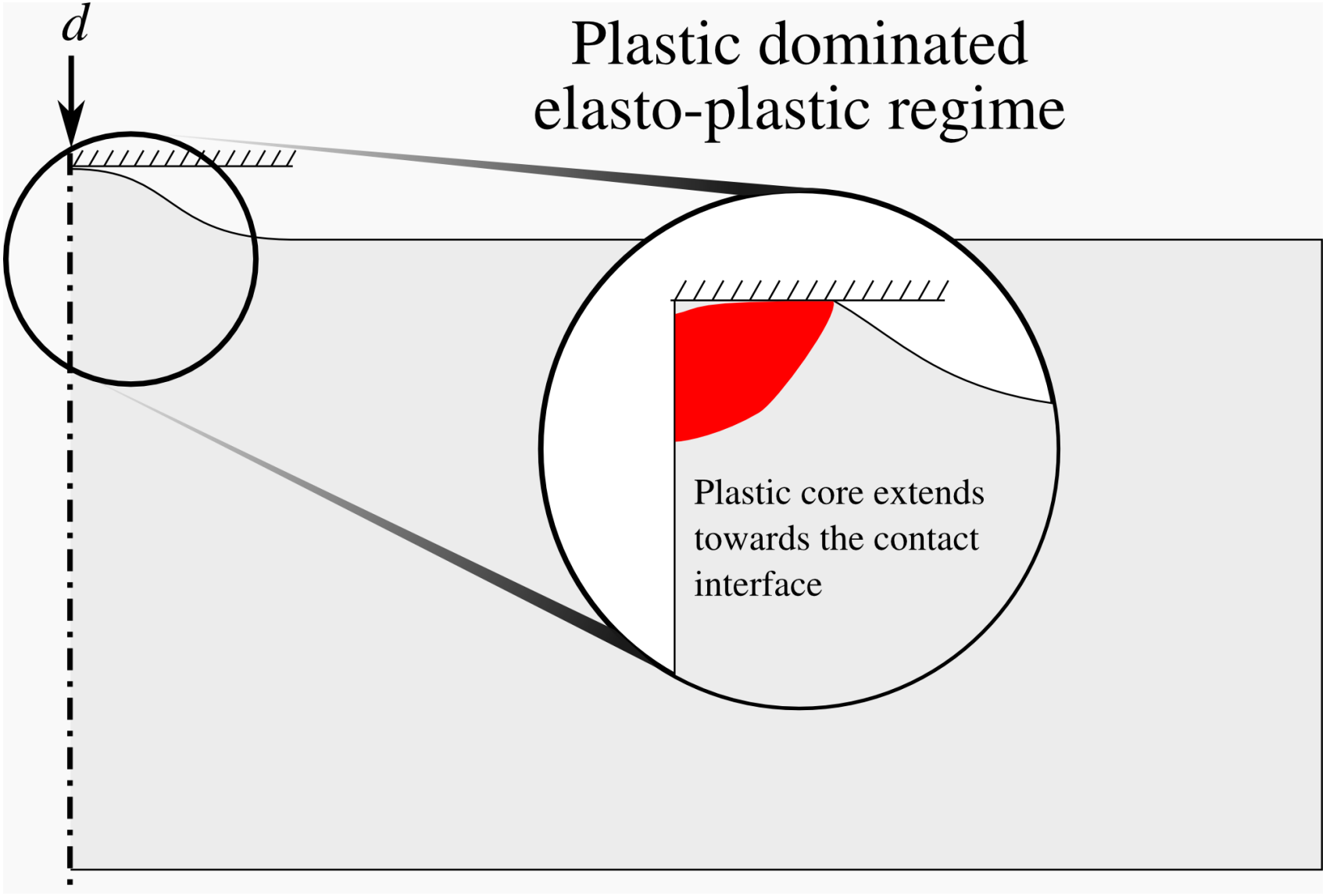

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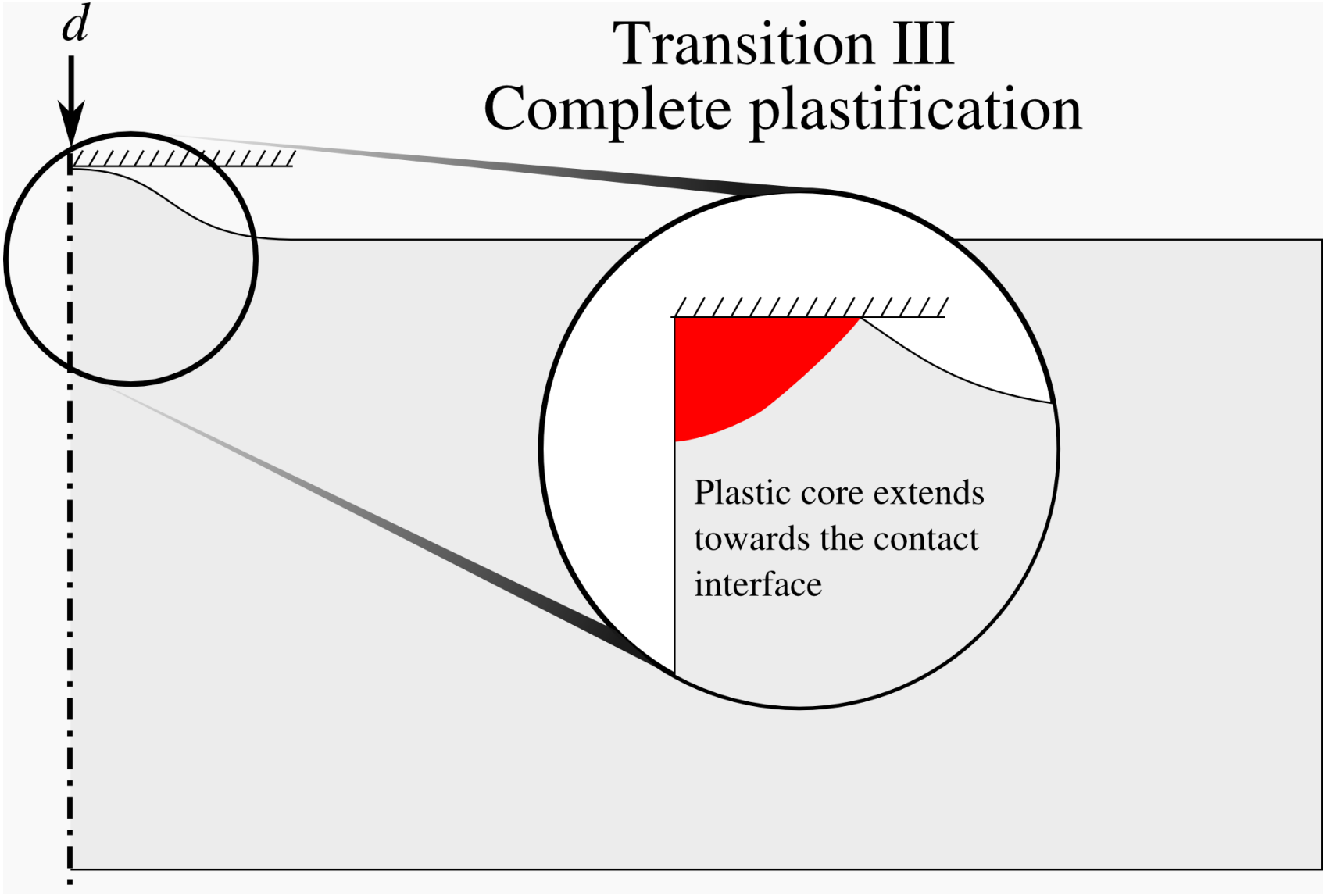

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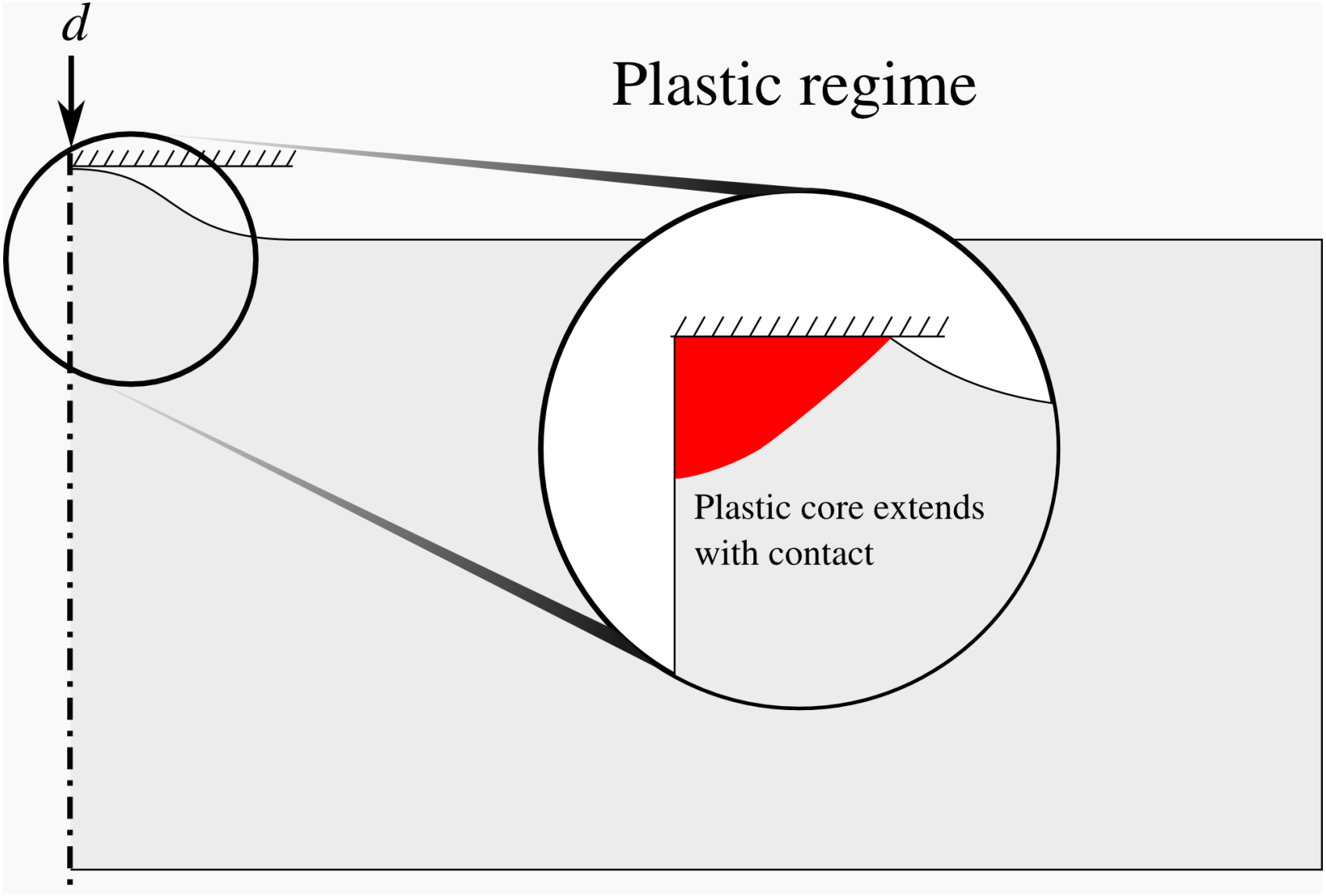

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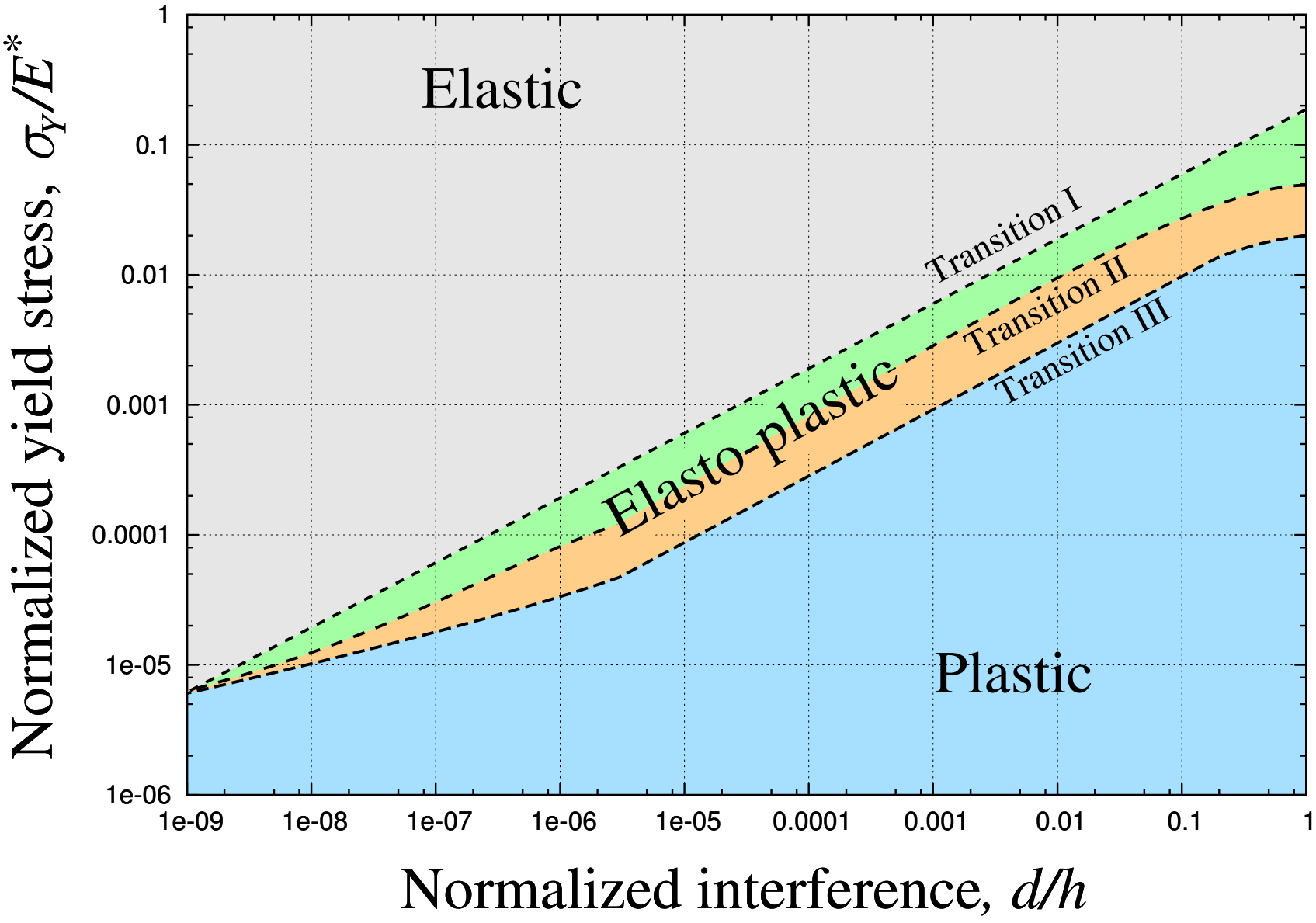
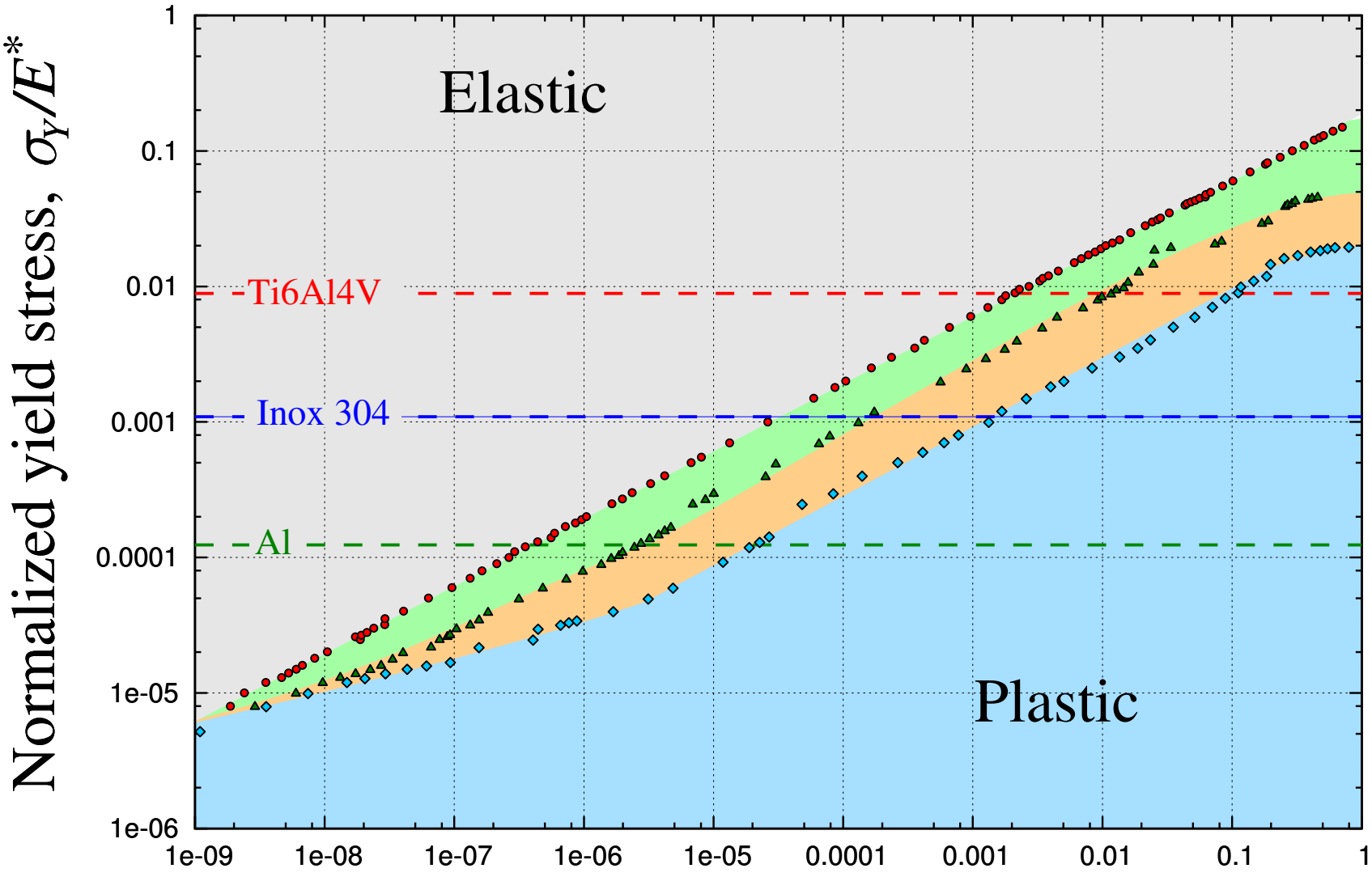


Deformation map for a sinusoidal asperity constructed with M. Liu, H. Proudhon

$h/\lambda = 0.1$

Elasto-plastic transition in contact



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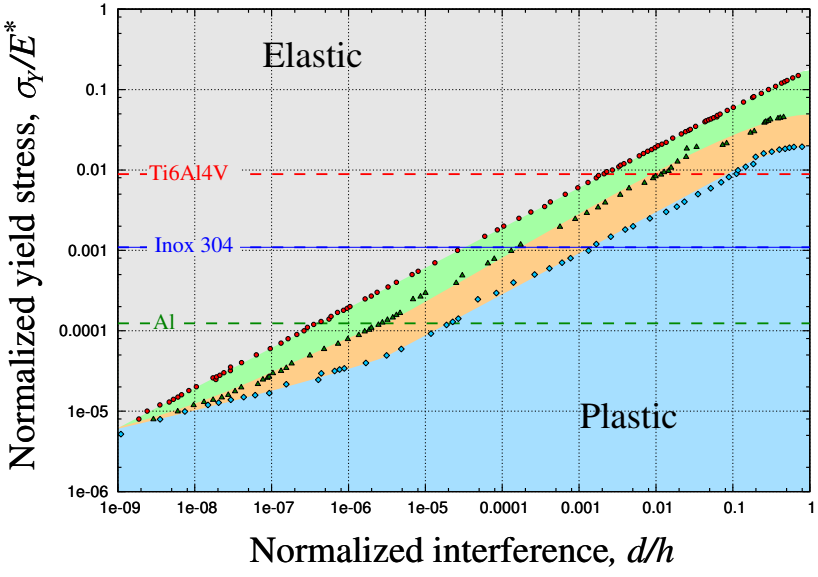
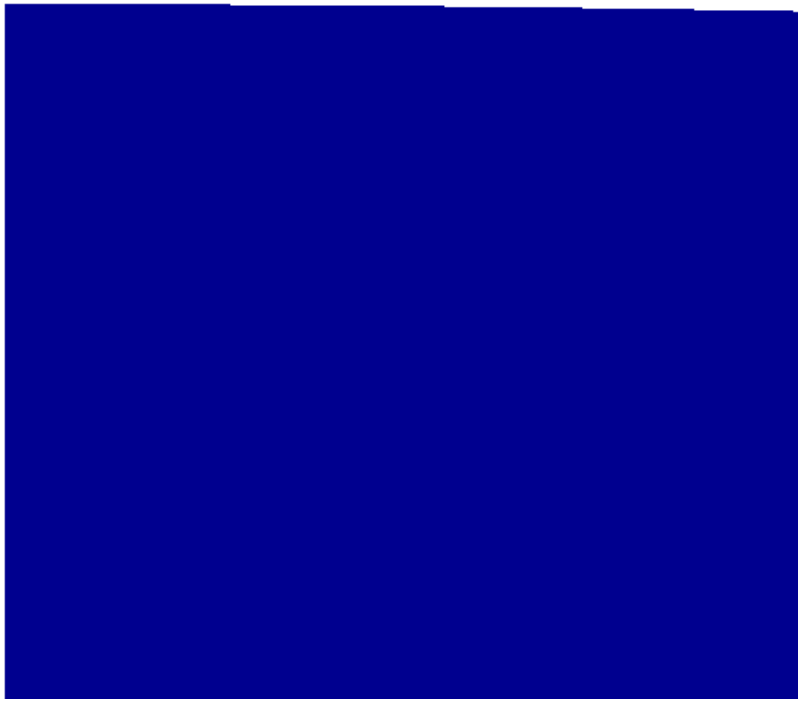
Case: $\sigma_Y/E = 0.0005$





Evolution of the plastic zone in a sinusoidal asperity in contact with a rigid flat

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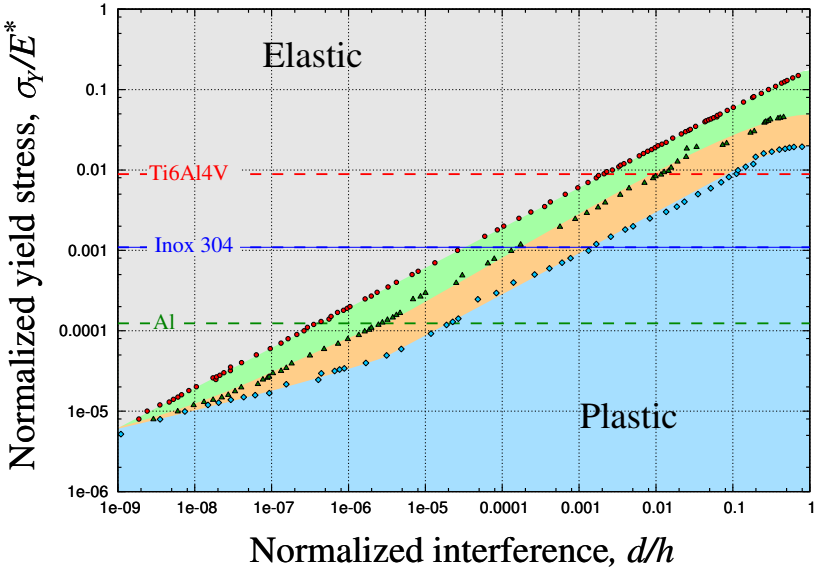

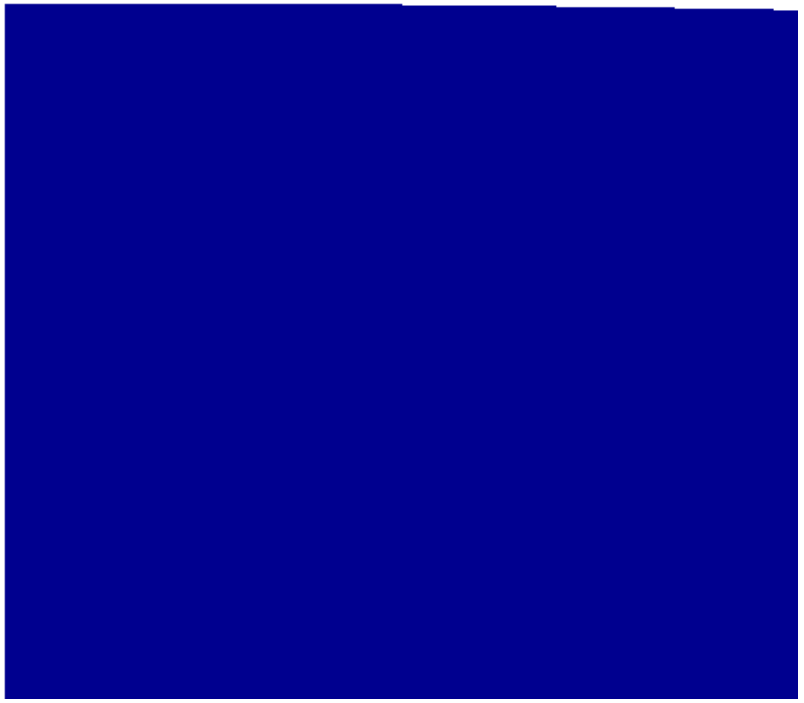

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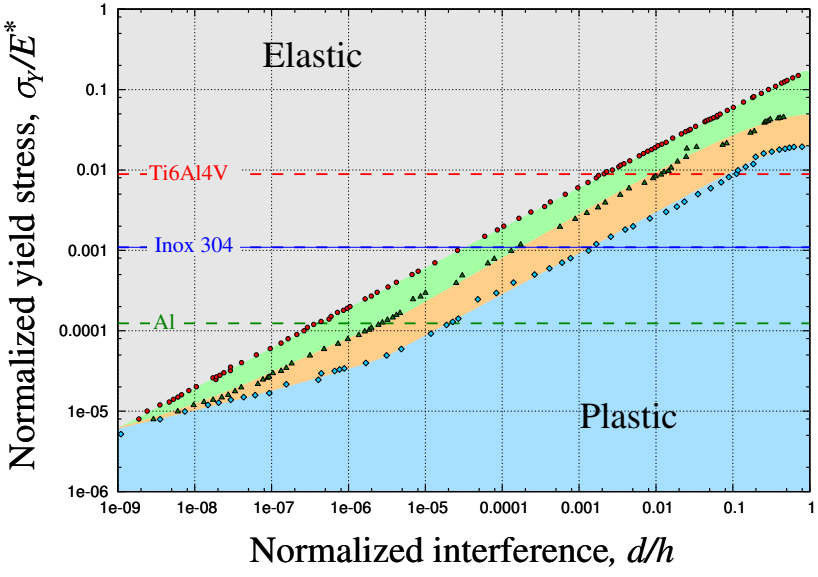

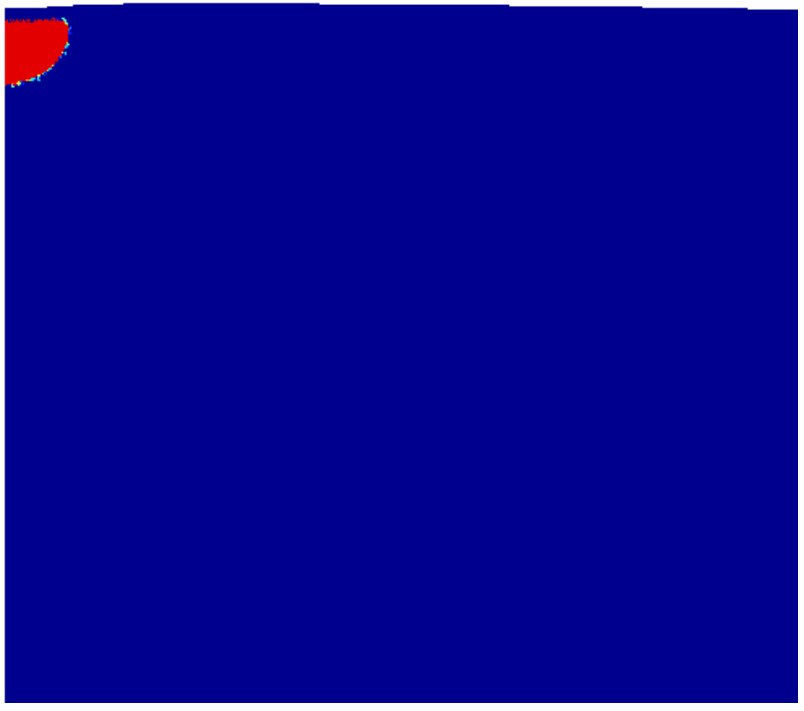

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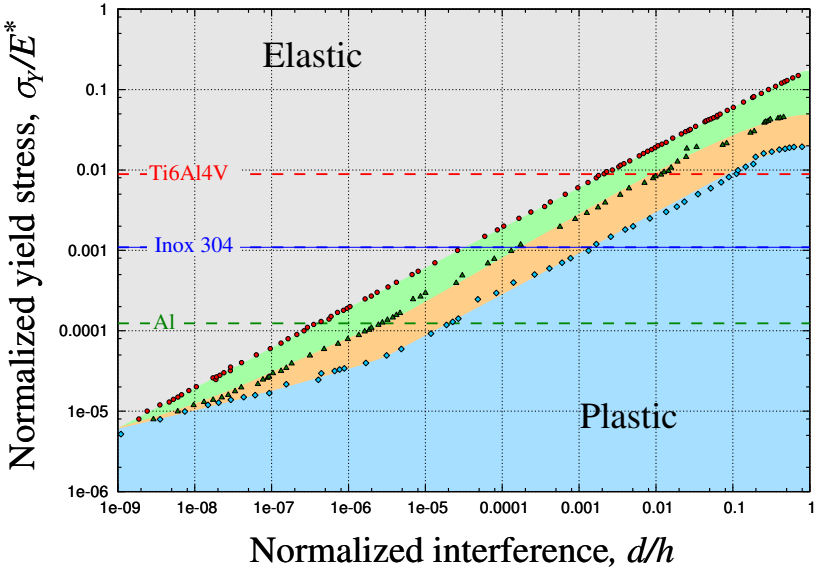

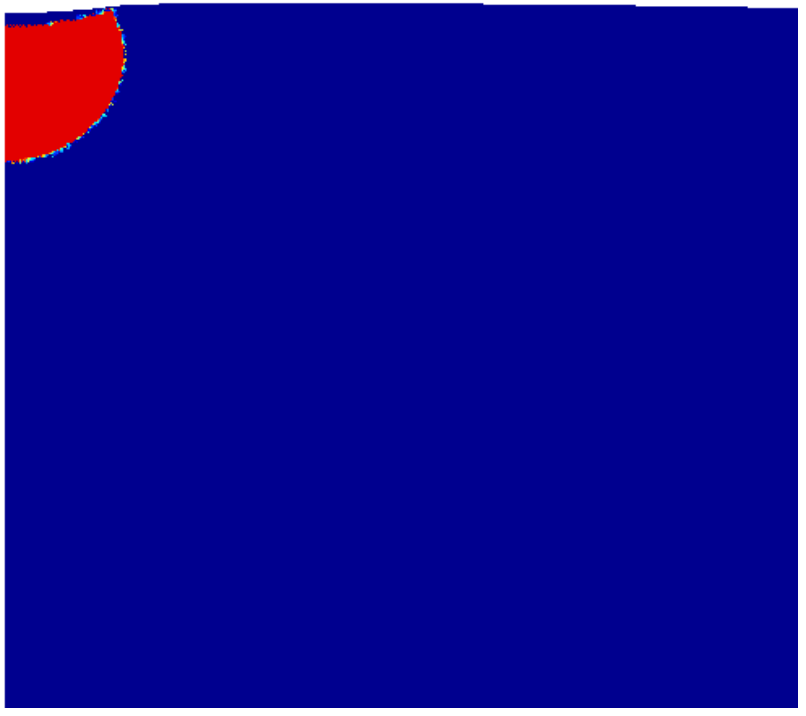

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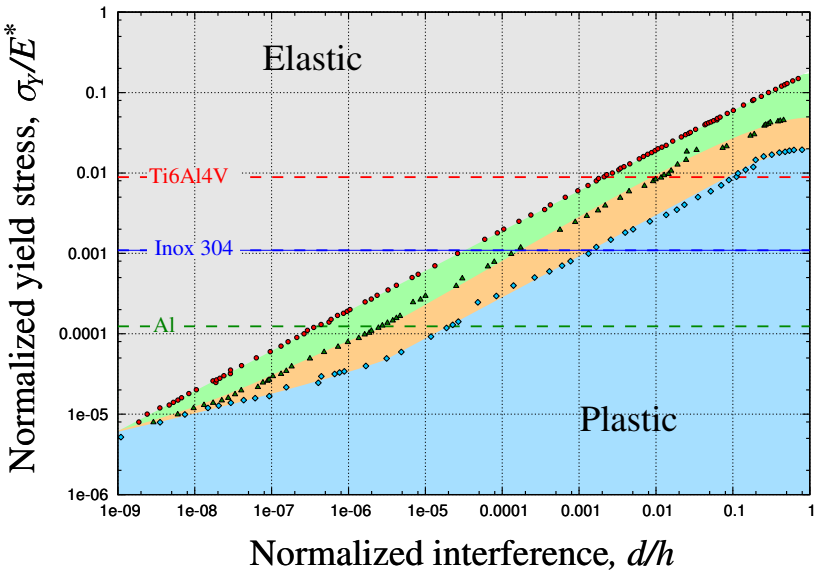

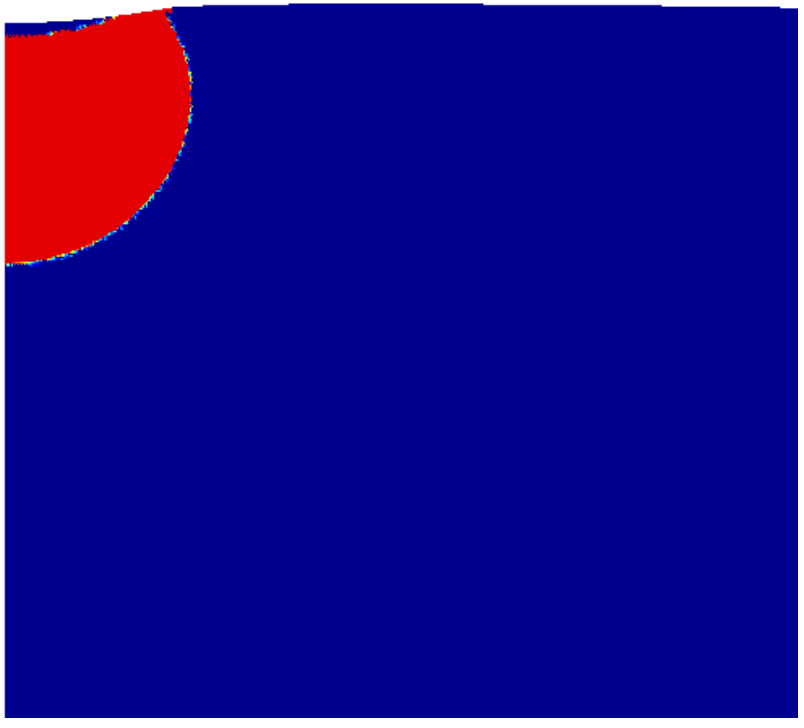

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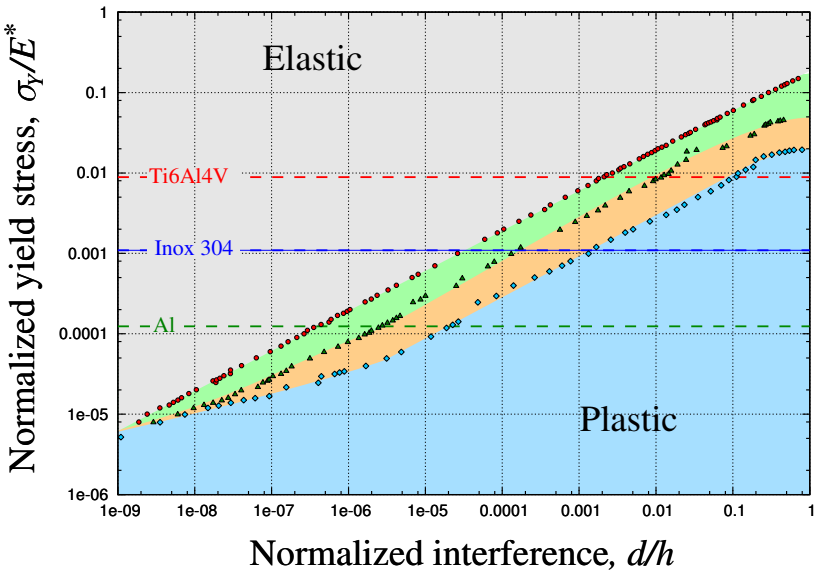

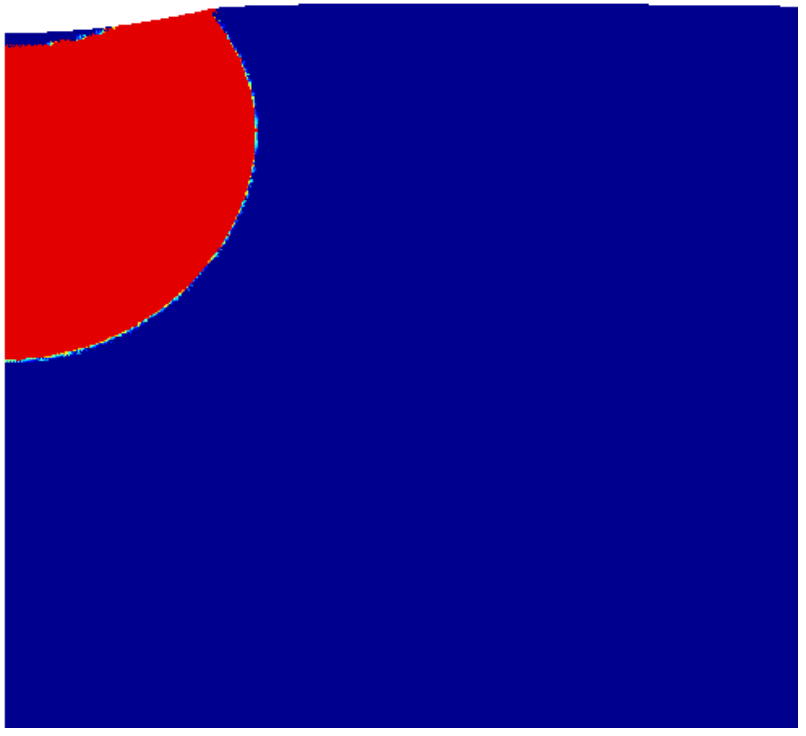

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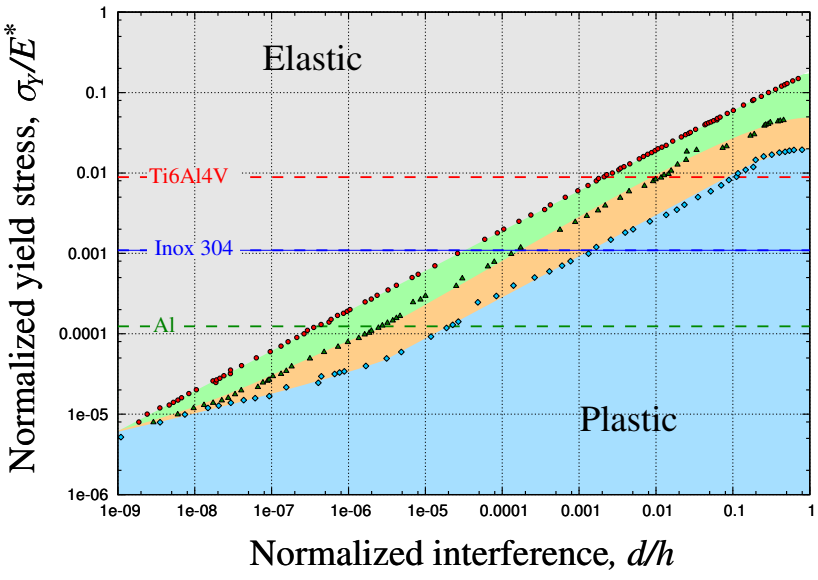

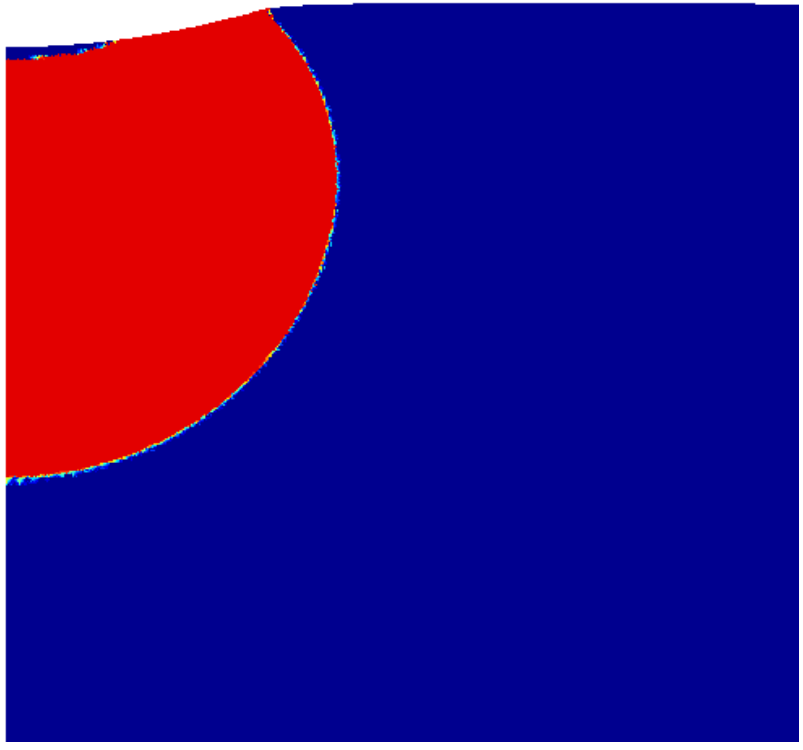

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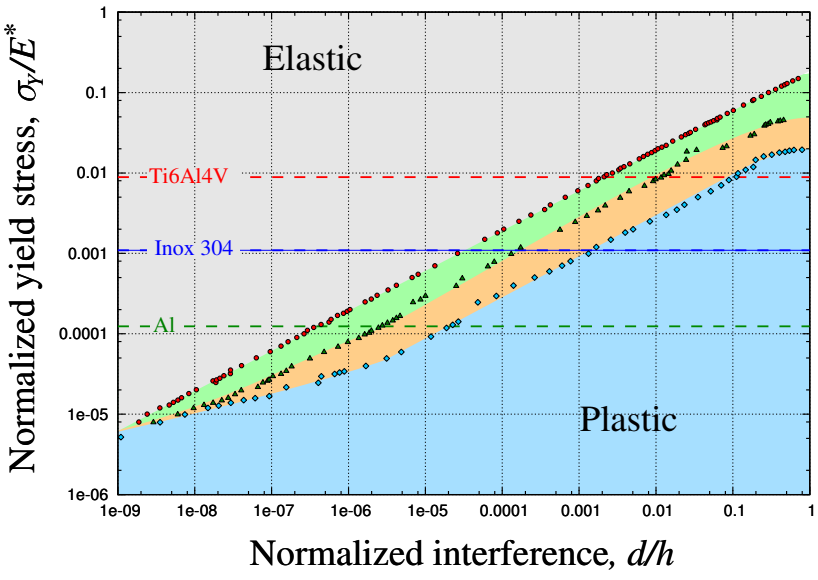
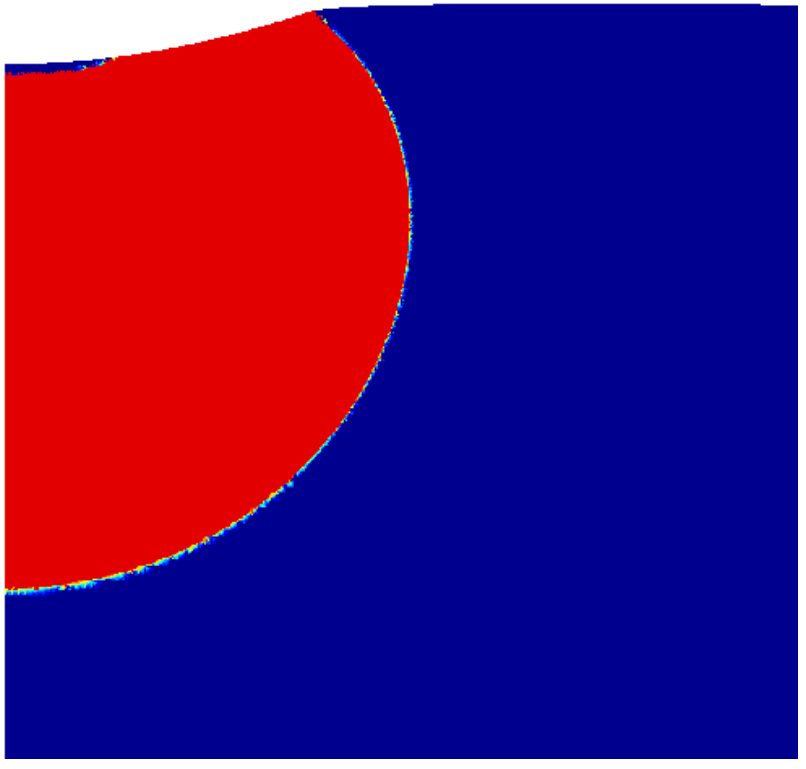




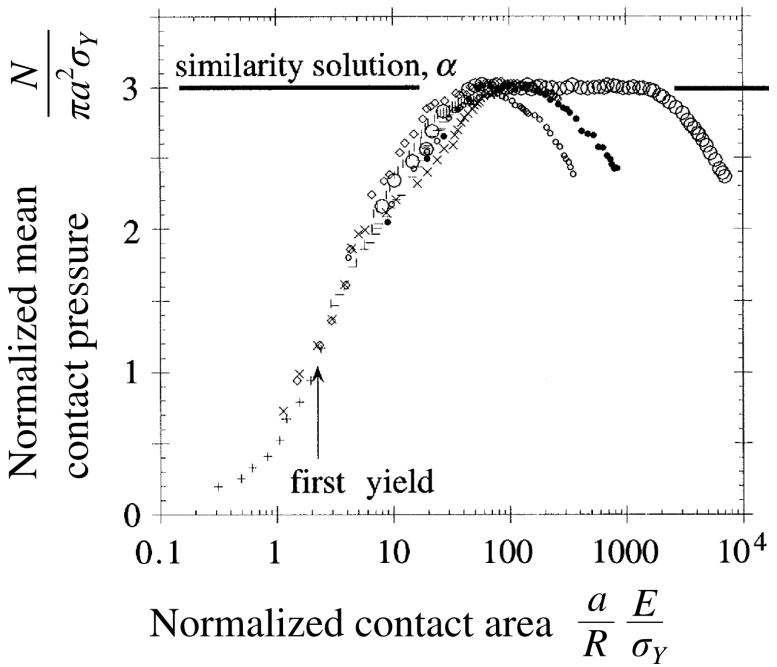
Evolution of the plastic zone in a sinusoidal asperity in contact with a rigid flat

Elastic-plastic normal contact: hardness

- Hardness ~ saturated plastic contact
- Recall: Vickers hardness $HV = N/A$
- Similarity solution[1]

$$\frac{N}{\pi a^2 \sigma_Y} = F\left(\frac{a}{R}\frac{E}{\sigma_y}\right)$$

- Hardness $H \approx 3\sigma_Y$



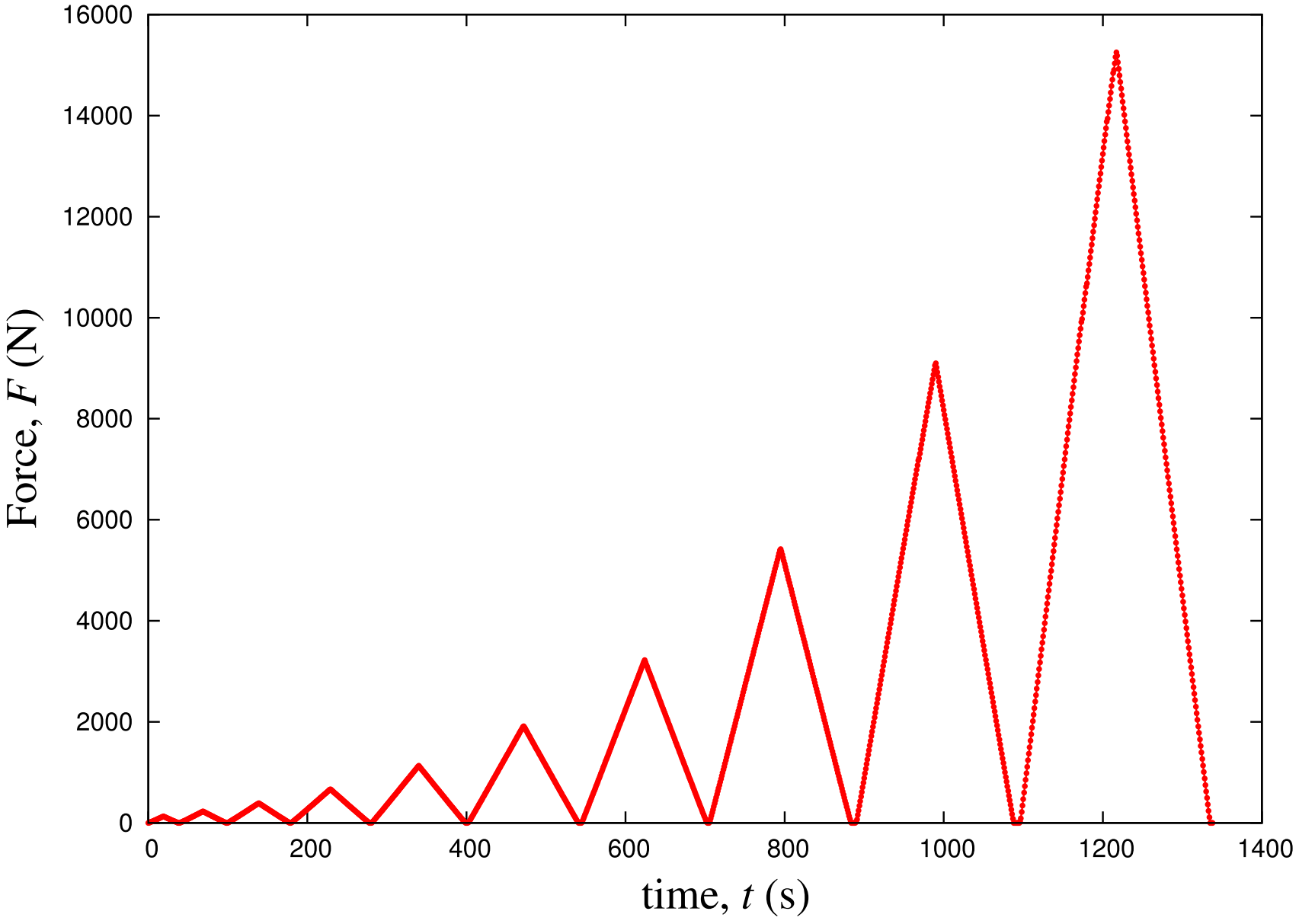
Simulation of spherical indentation of elasto-plastic solid with power-law hardening[2]

[1] Hill R., Storøakers B., Zdunek A.B. A theoretical study of the Brinell hardness test. Proc R Soc Lond A 436 (1989)

[2] Tabor, D. The Hardness of Metals. Oxford at the Clarendon Press (2000).

[3] Mesarovic S., N. Fleck, Spherical Indentation of Elastic-Plastic Solids, Proc R Soc Lond A 455 (1999)

Elasto-plastic contact under cyclic load



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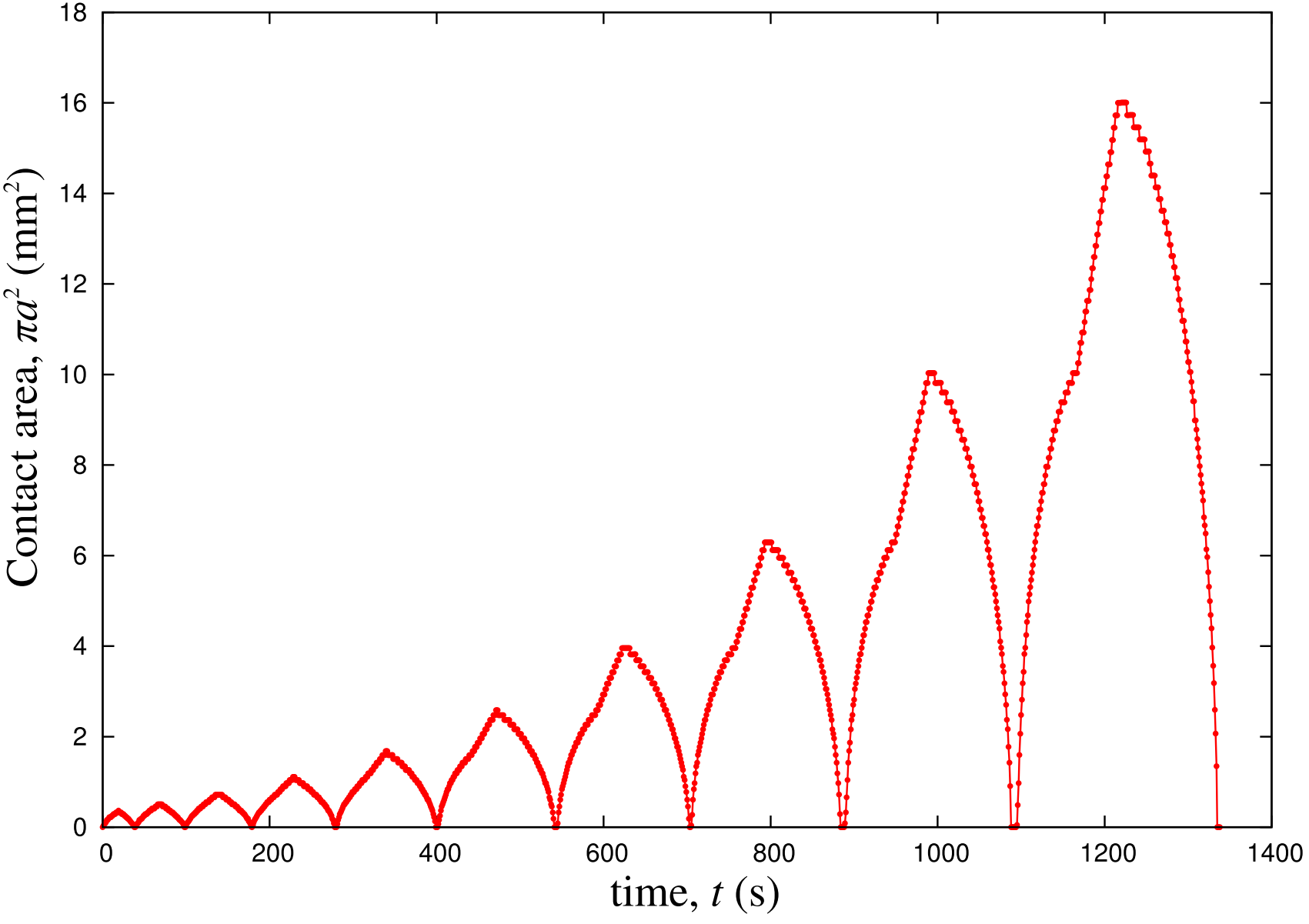

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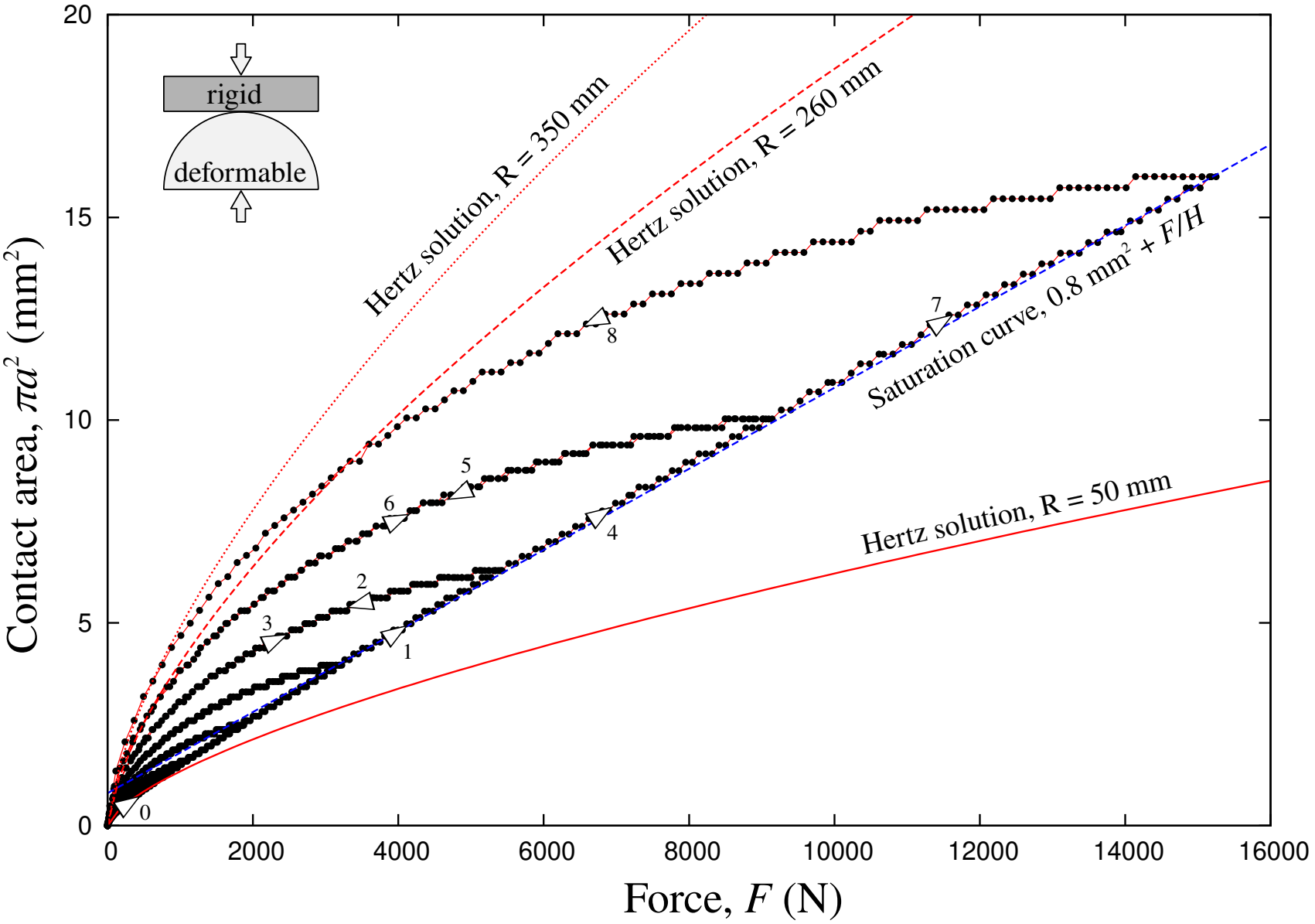

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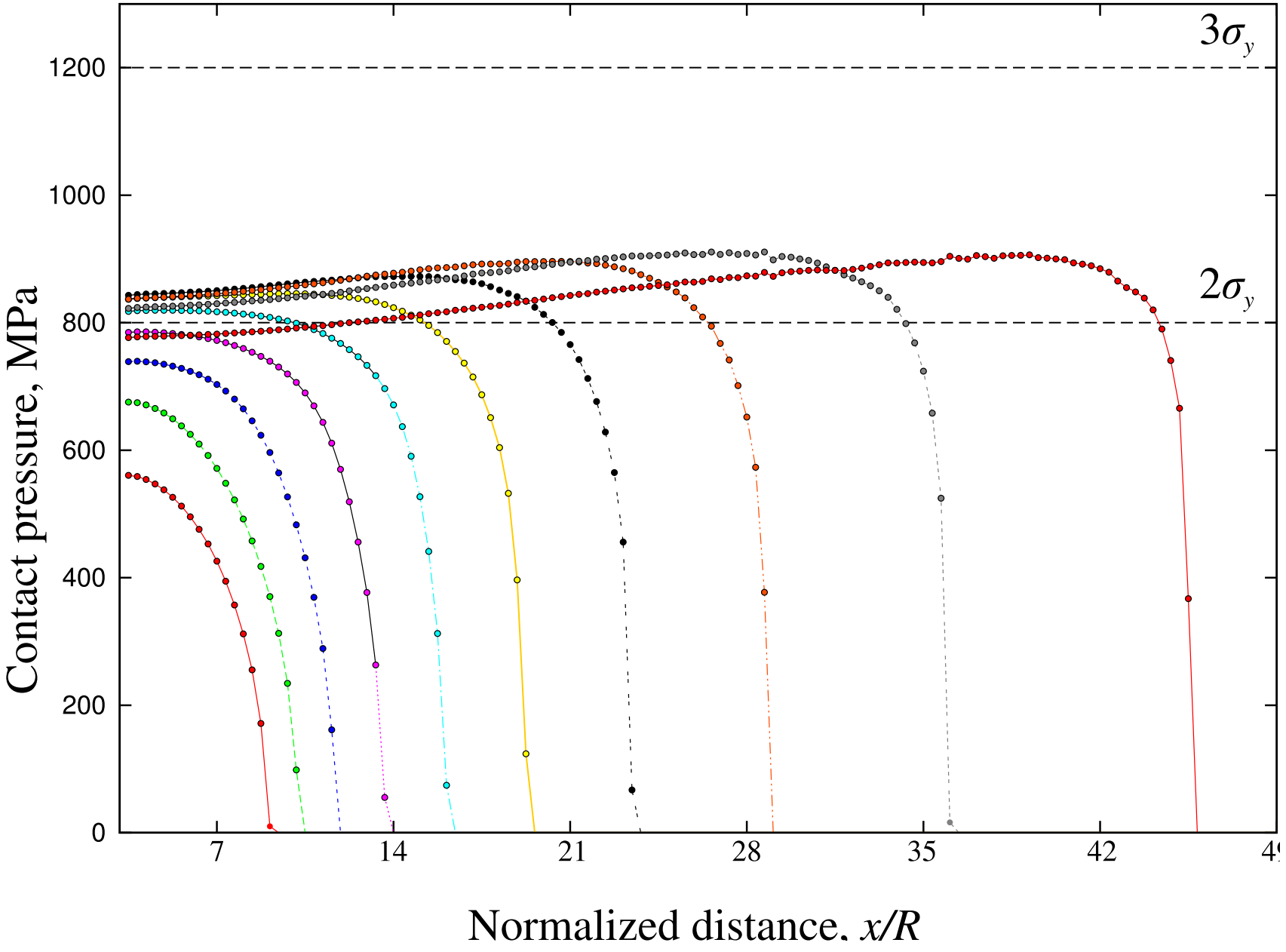

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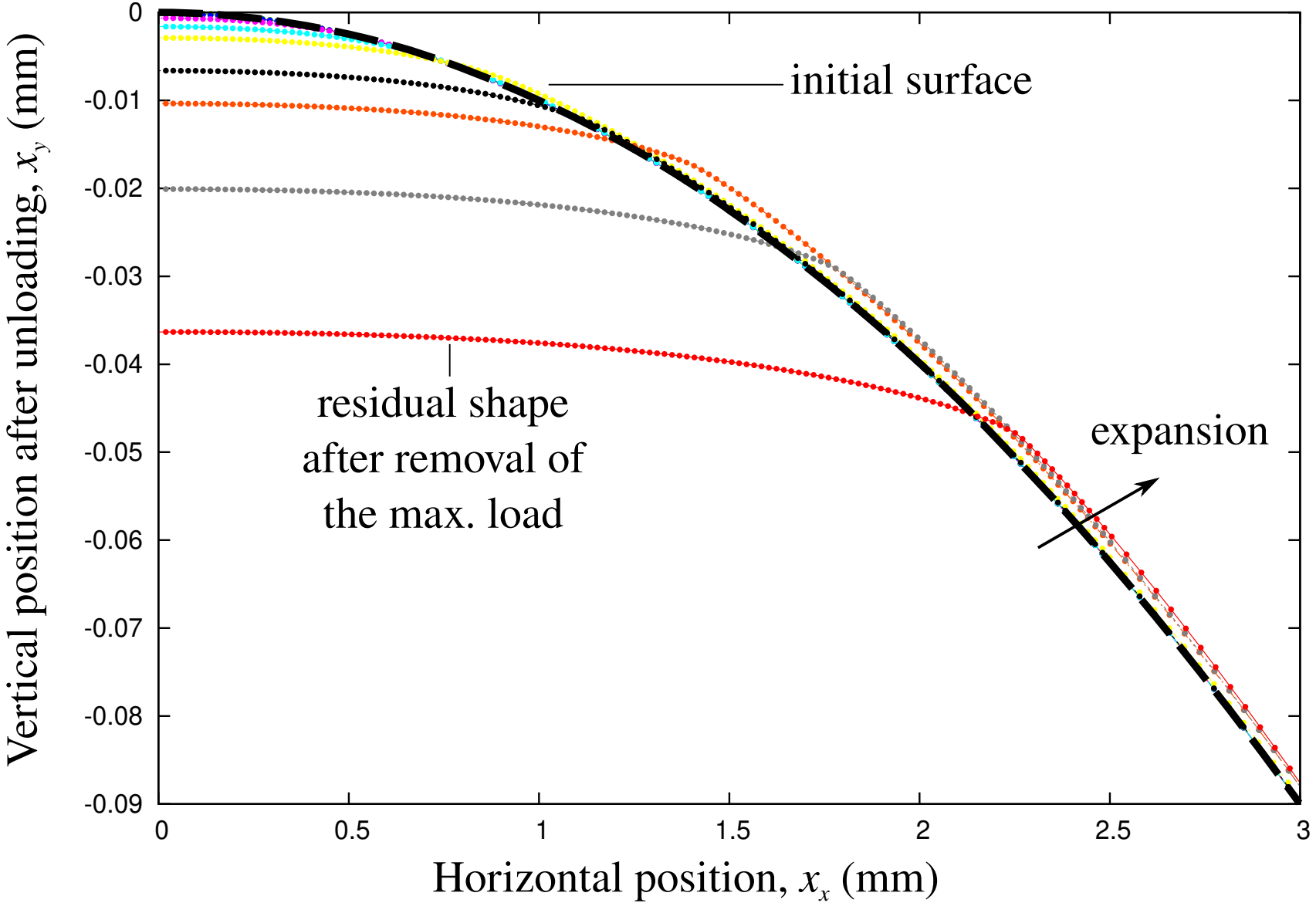
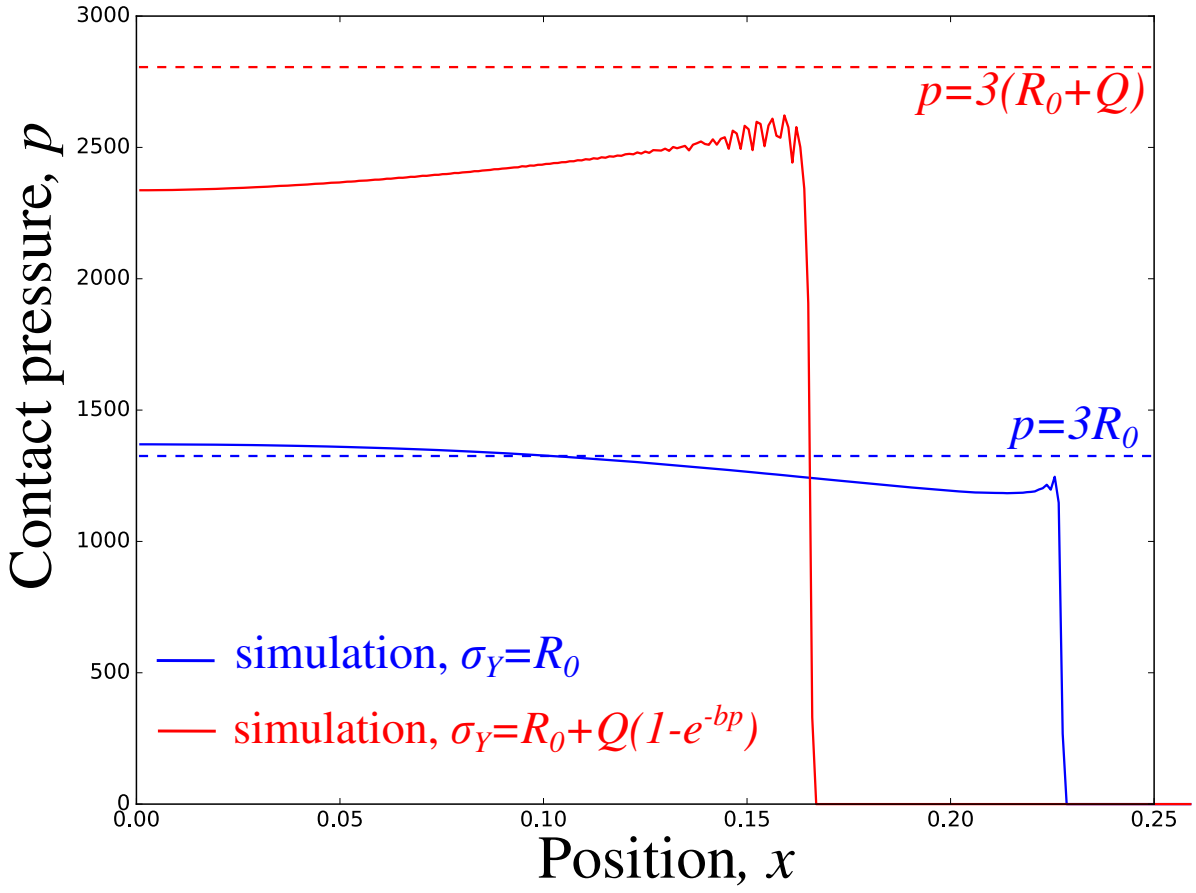


Finer-mesh and hardening effect

Two hardening models

- perfect plasticity $R = R_0$
- non-linear hardening with saturation $R(p) = R_0 + Q(1 - \exp(-bp))$



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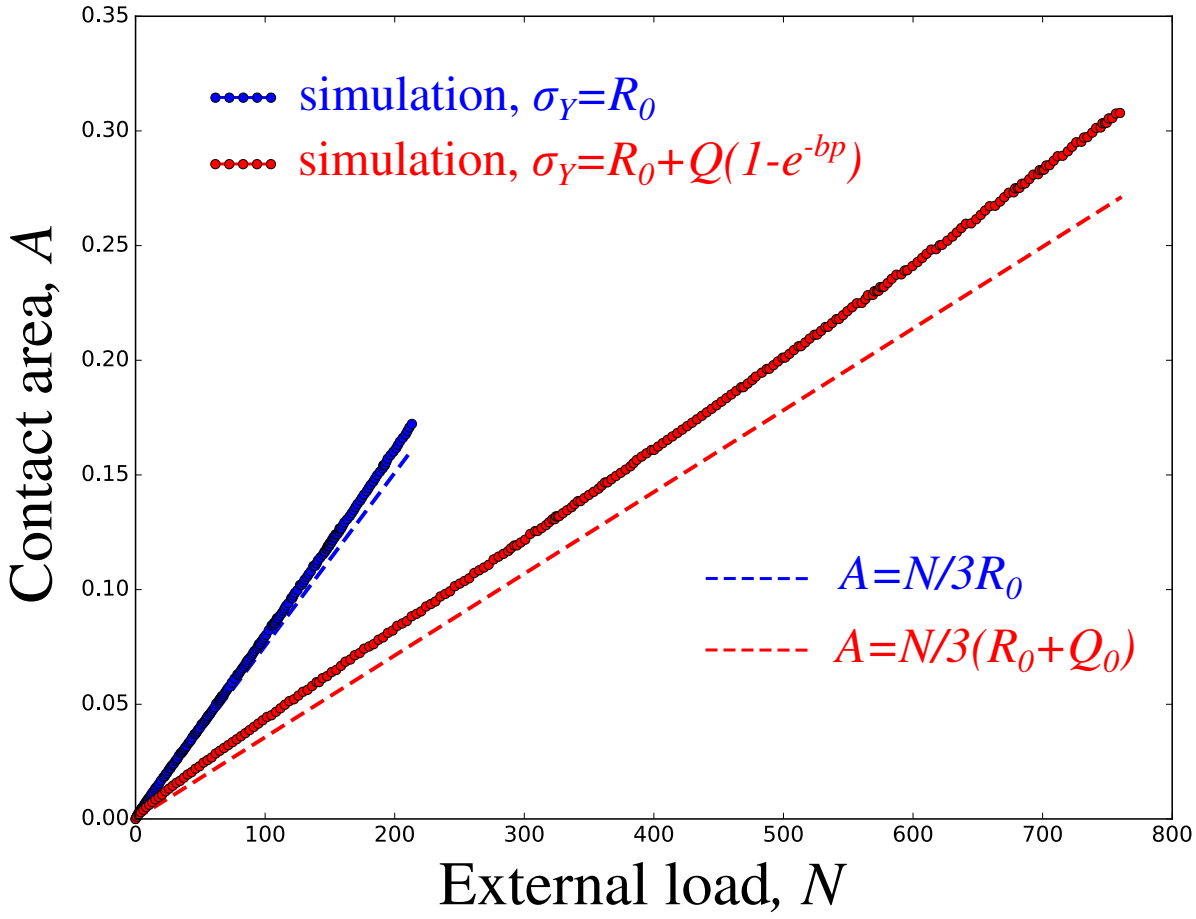
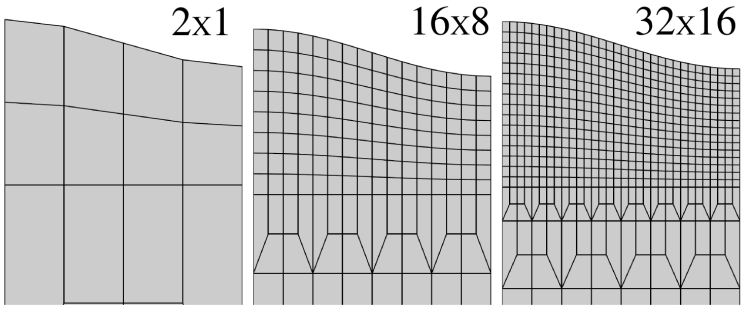
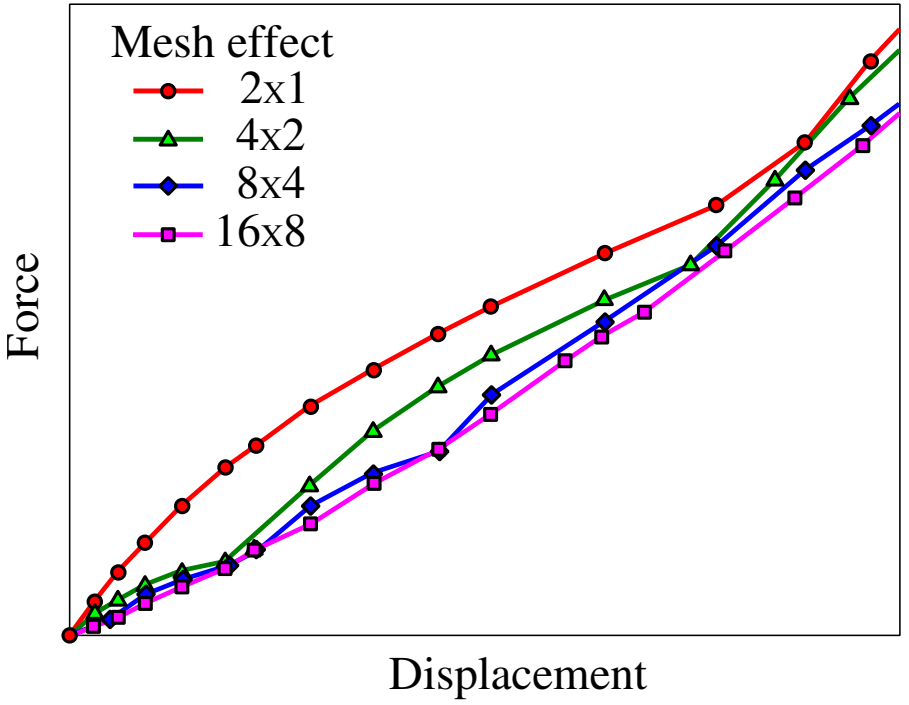


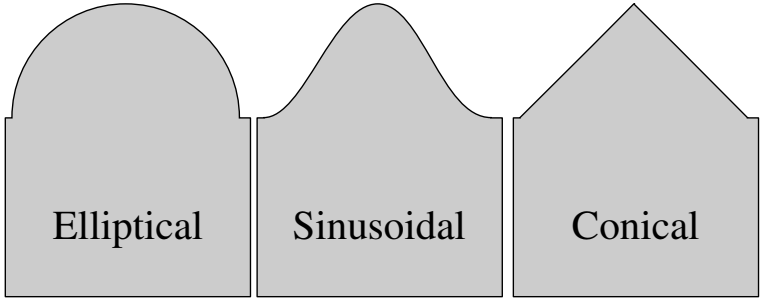
Deformation in fully plastic regime

✓ **Mesh effect**

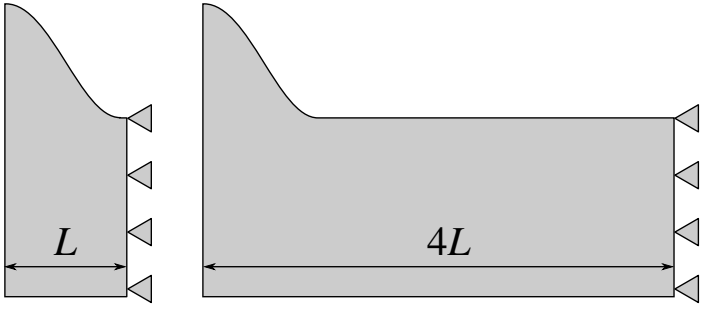




Shape effect



Edge effect



[1] V. A. Yastrebov, J. Durand, H. Proudhon, G. Cailletaud, CR Mecan, 339:473-490 (2011)

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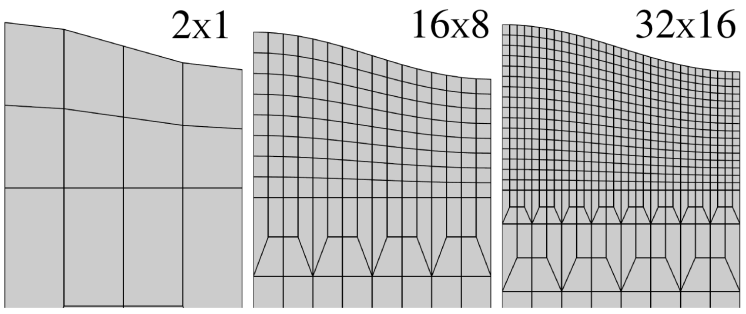

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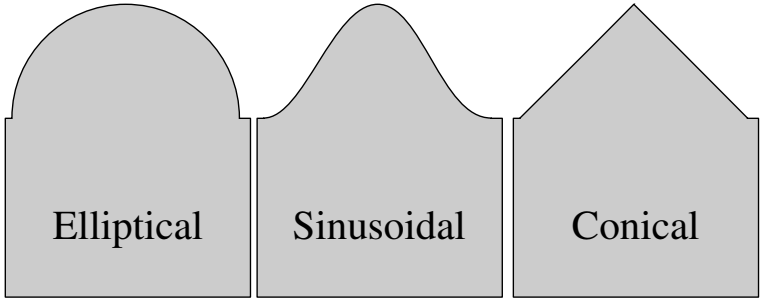

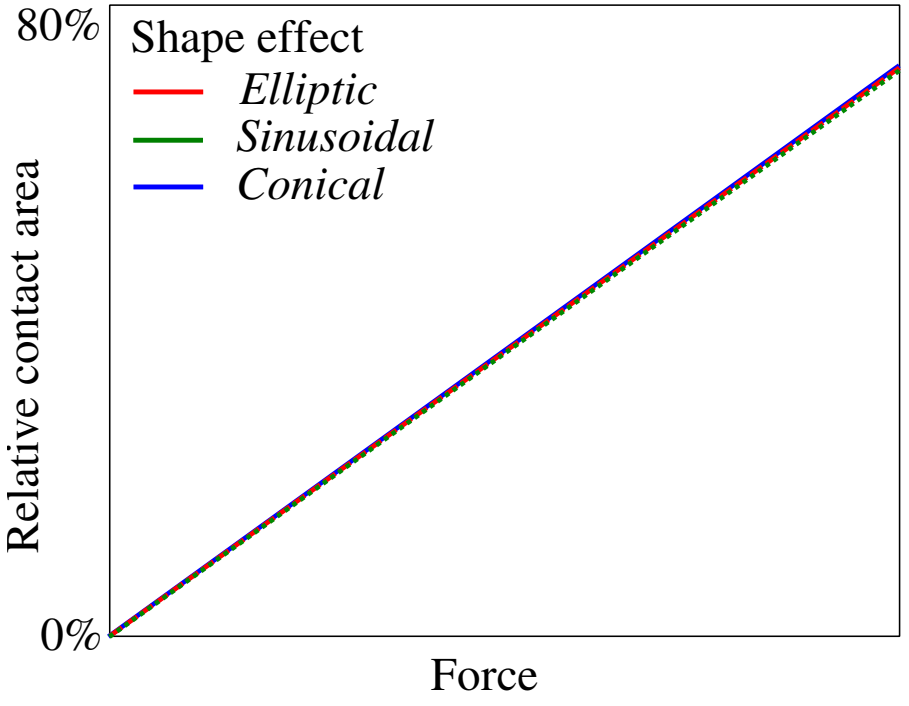

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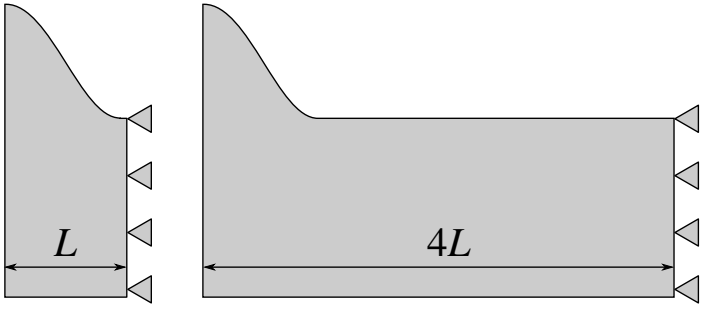

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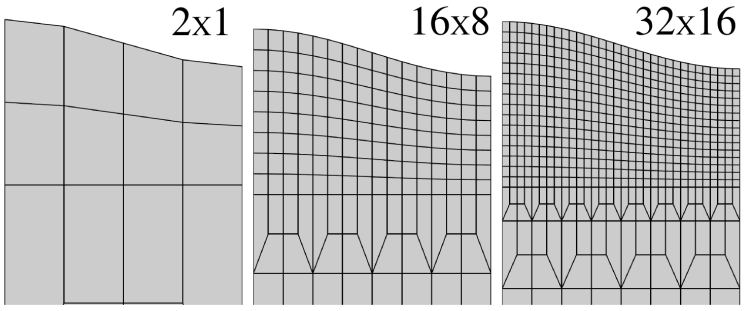

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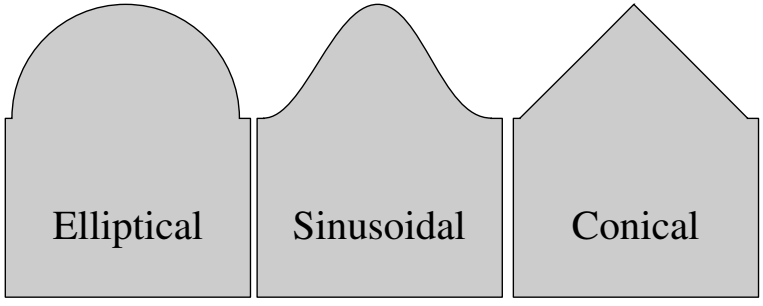

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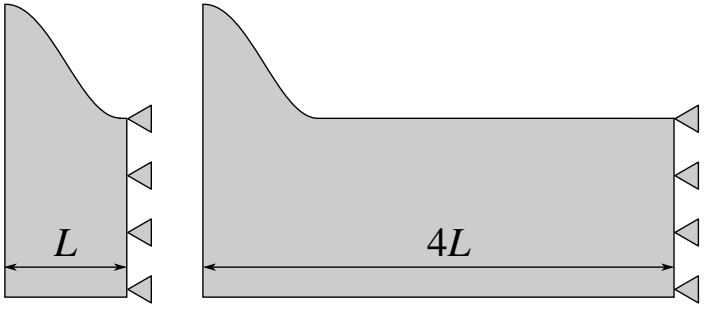

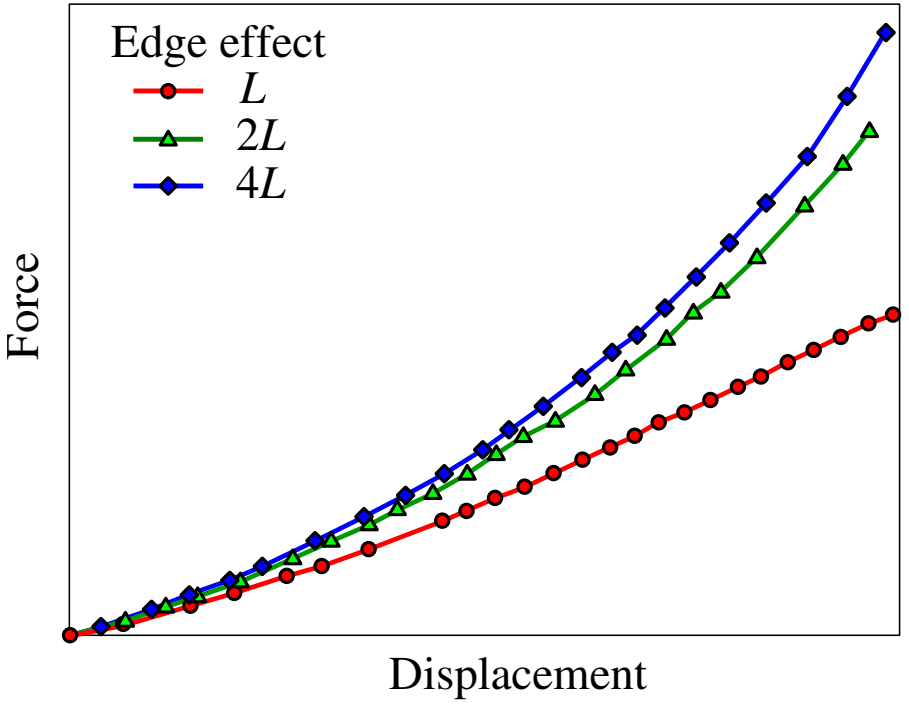

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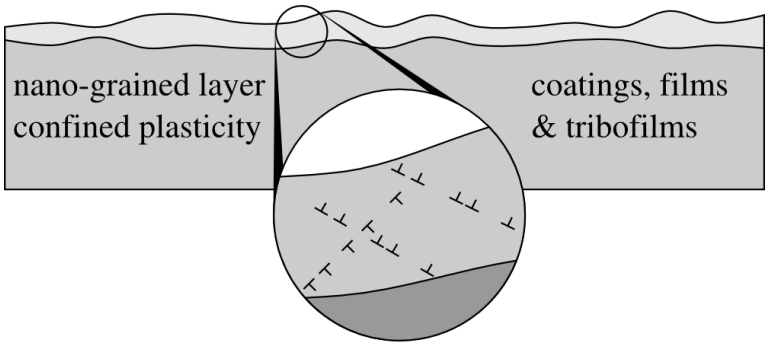
Near-surface vs bulk deformation

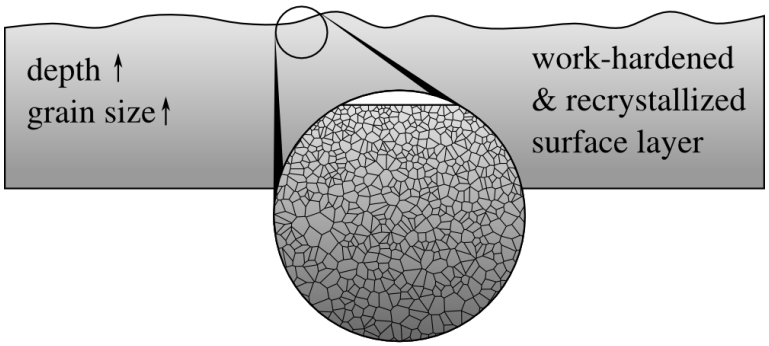
Material aspects

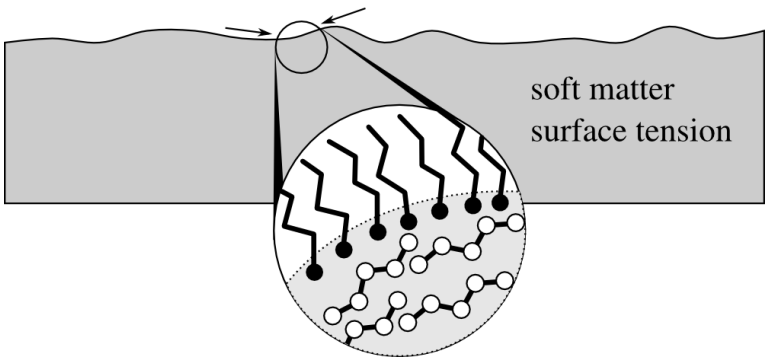
- Cold worked surface + recrystallized:
  smaller grains near the surface, Hall-Petch effect
- Thin coating films:
  nanograined, confined plasticity, Hall-Petch effect
- Oxides:
  brittle hard films

Geometrical aspects

- Roughness of all nature
- Indentation by asperities:
  confined plastic zone, high plastic strain gradients







Onset of yielding at atomic scale

Hertz contact: body of revolution

- Onset of plasticity for pressure

$$p_Y = 1.6\sigma_Y$$

- Associated force

$$F_Y = \frac{1.6^3 \pi^3 R^2}{6} \left(\frac{\sigma_Y}{E^*}\right)^2 \sigma_Y$$

- Associated contact radius

$$a_Y = \frac{1.6\pi R}{2} \frac{\sigma_Y}{E^*}$$

- Plastic flow starts at the depth

$$z_Y \approx 1.21 R \frac{\sigma_Y}{E^*}$$

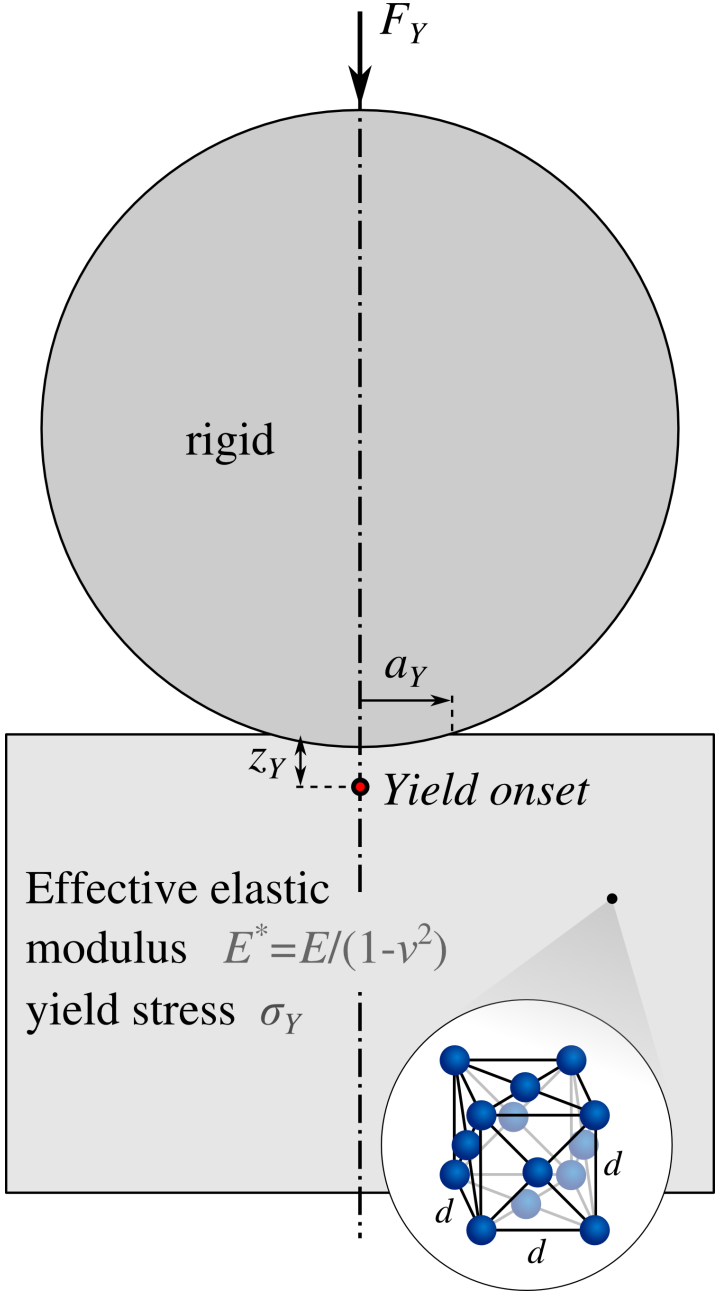
- Example: golden asperity

$E^* \approx 96$ GPa, $\sigma_y \approx 140$ MPa, $d \approx 4.1$Å

$R = 10\ \mu$m

$F_Y \approx 3.8\ \mu$N, $z_Y \approx 18$ nm, $a_Y \approx 36$ nm

$z_Y \approx 45d$, $a_Y \approx 115d$



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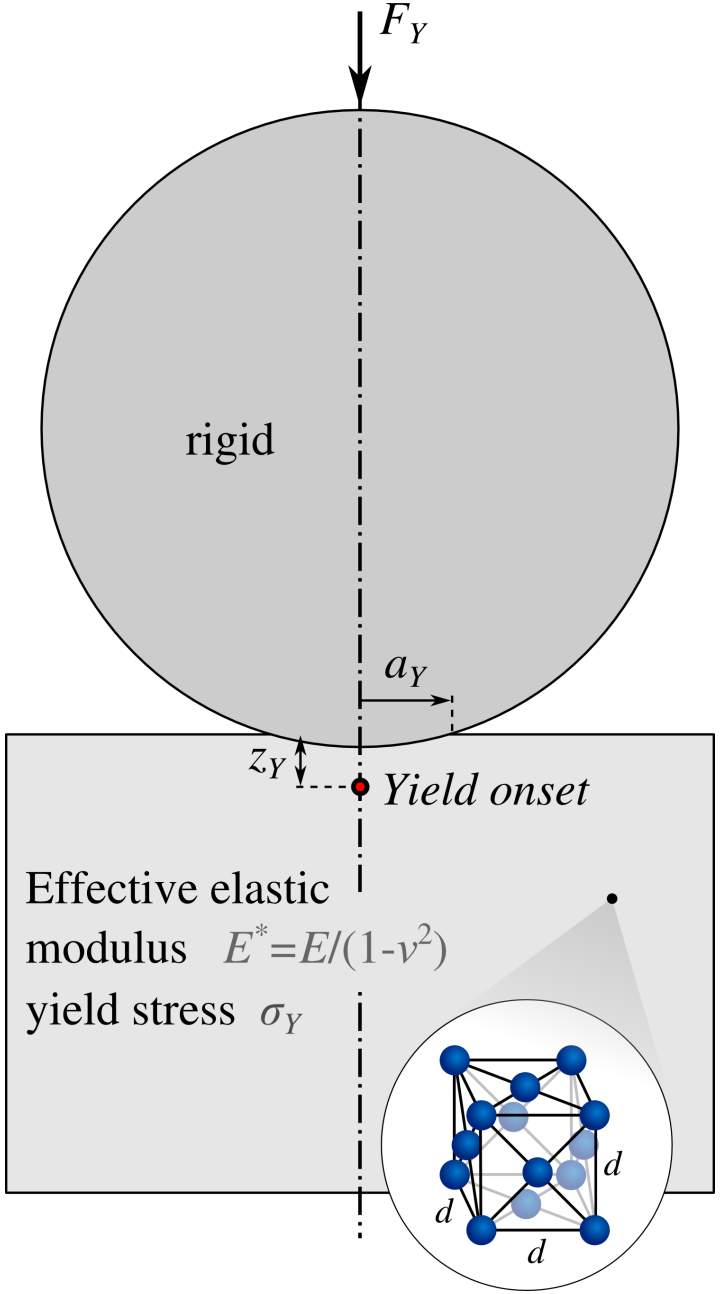
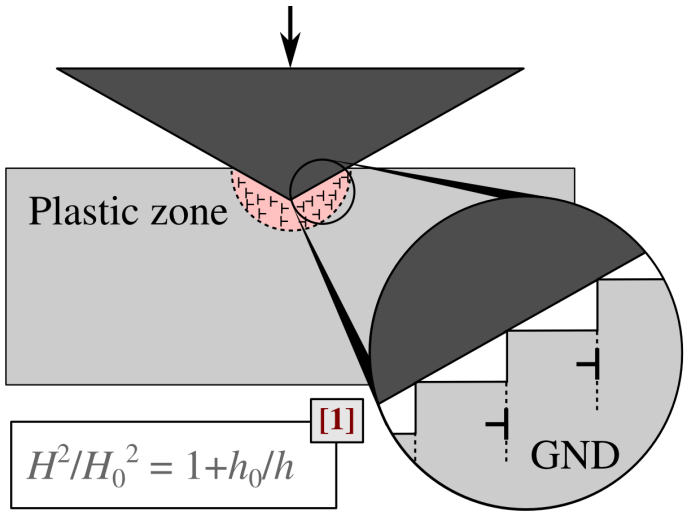
$R = 1\ \mu$m

$$F_Y \approx 38 \text{ nN}, \quad z_Y \approx 1.8 \text{ nm}, \quad a_Y \approx 3.6 \text{ nm}$$

$$z_Y \approx 4.5d, \quad a_Y \approx 11.5d$$



Indentation and hardness



[1] Nix, Gao. *J Mech Phys Solids* (1998).

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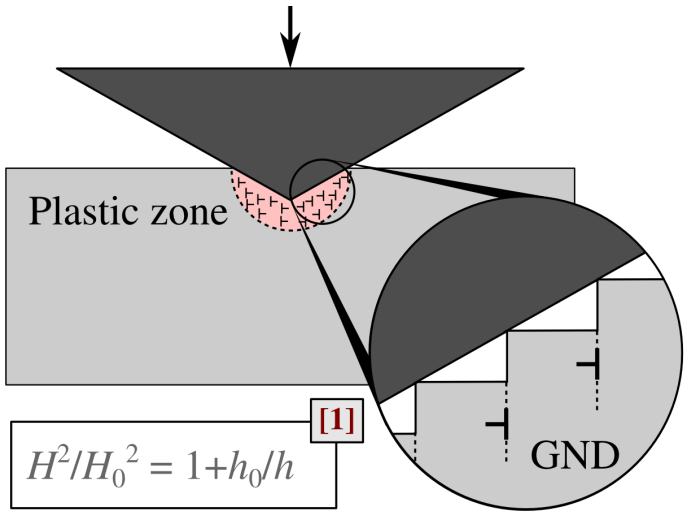

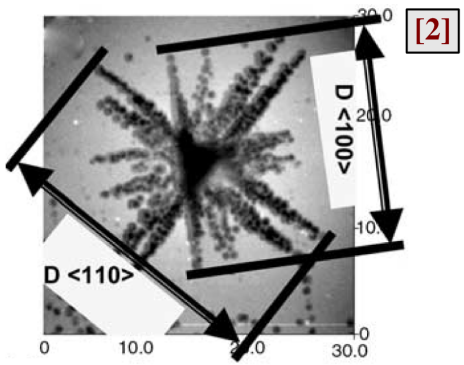

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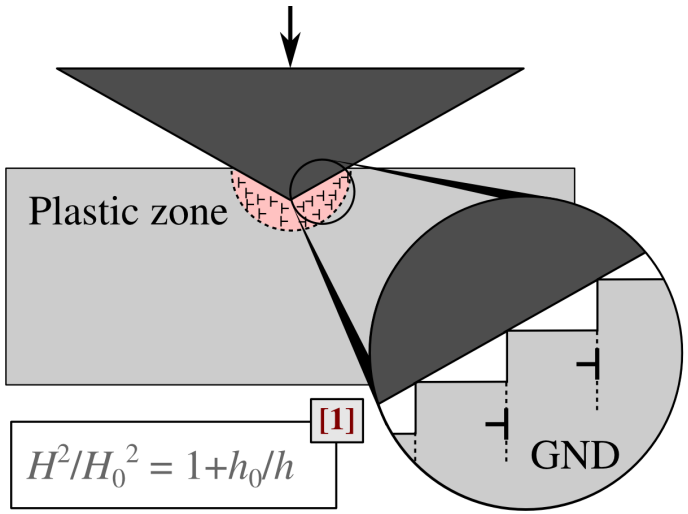

[1] $H^2/H_0^2 = 1+h_0/h$

$$\frac{\text{Plastic zone } r}{\text{Contact radius } a} \sim 1 + b\exp(-h/h_1)$$

[2] $H^2/H_0^2 = 1+[1+b\exp(-h/h_1)]^{-3}h_0/h$

[2,3] $(H-H_P)^2/(H_0-H_P)^2 = 1+[1+b\exp(-h/h_1)]^{-3}h_0/h$

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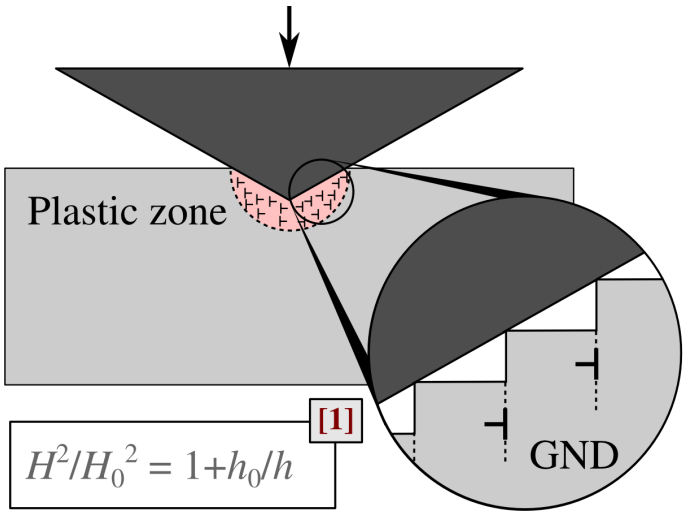

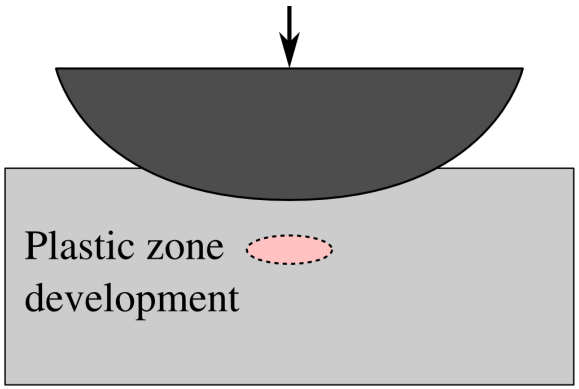

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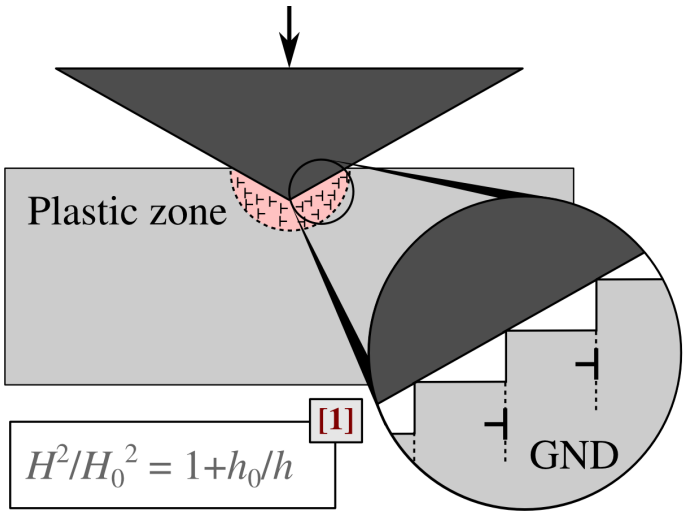

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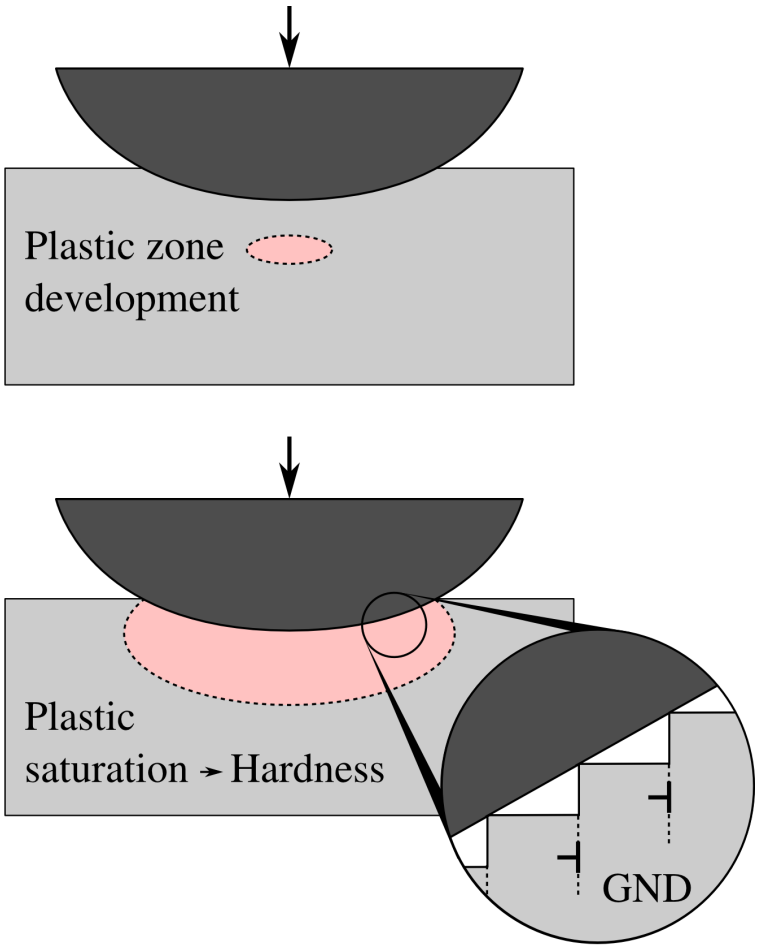

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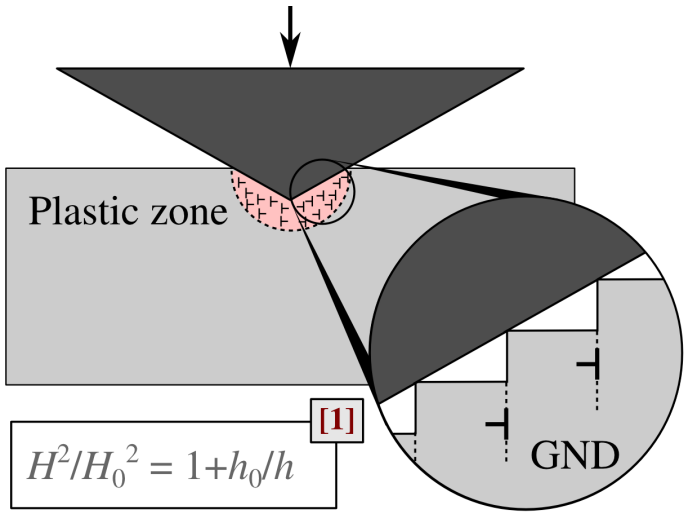

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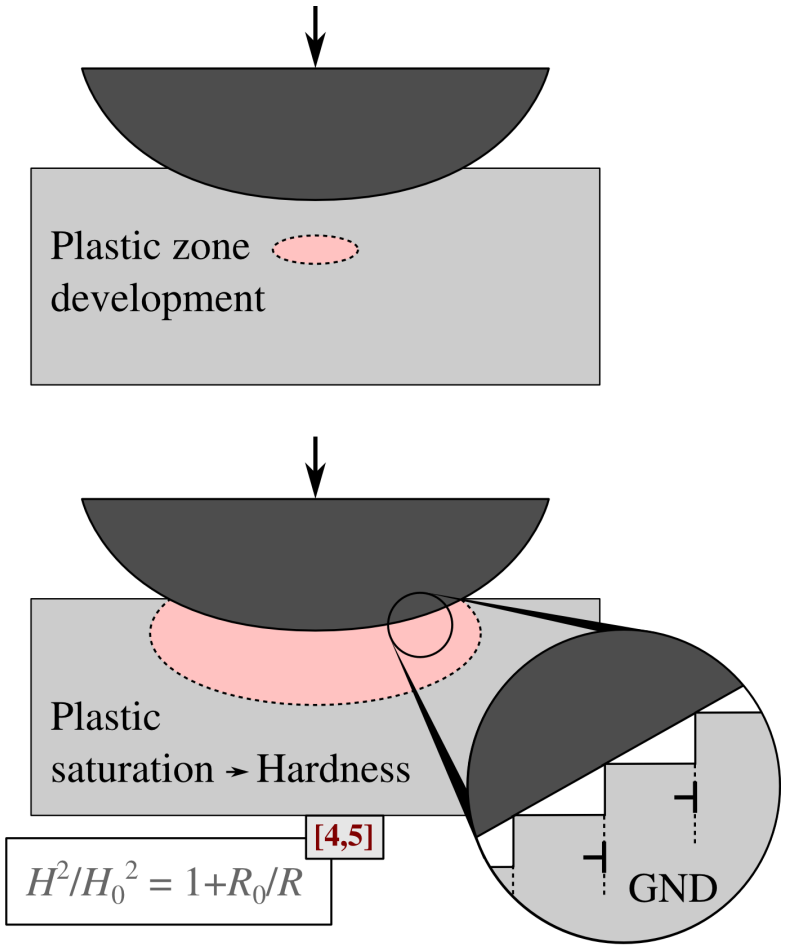

[4,5] $H^2/H_0^2 = 1+R_0/R$

[1] Nix, Gao. *J Mech Phys Solids* (1998).
[2] Feng, Nix. *Scripta Mater* (2004).
[3] Qui, Huang, Nix, Hwang, Gao. *Acta Mater* (2001).
[4] Swadener, George, Pharr. *J Mech Phys Solids* (2002).
[5] Gao, Larson, Lee, Nicola, Tischler, Pharr. *J Appl Mech* (2015).

Enhanced material behavior: beyond *J2* plasticity

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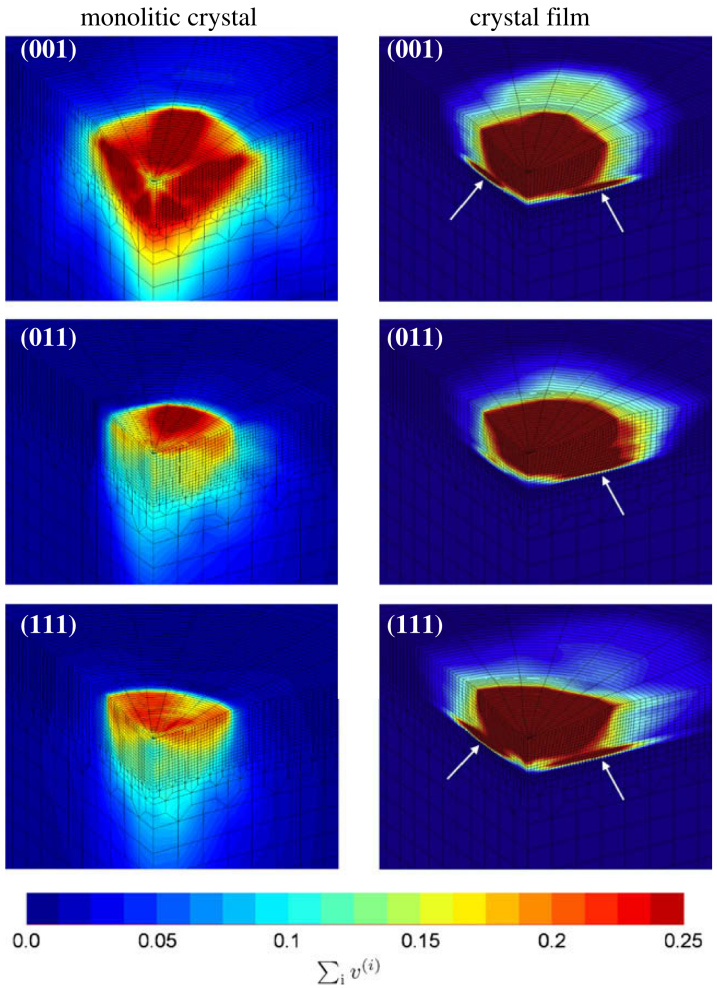

Spherical indentation of FCC copper crystal using crystal plasticity model in Zset

[3] Casals O., Forest S. Finite element crystal plasticity analysis of spherical indentation in bulk single crystals and coatings. Comp Mater Sci 45 (2009)

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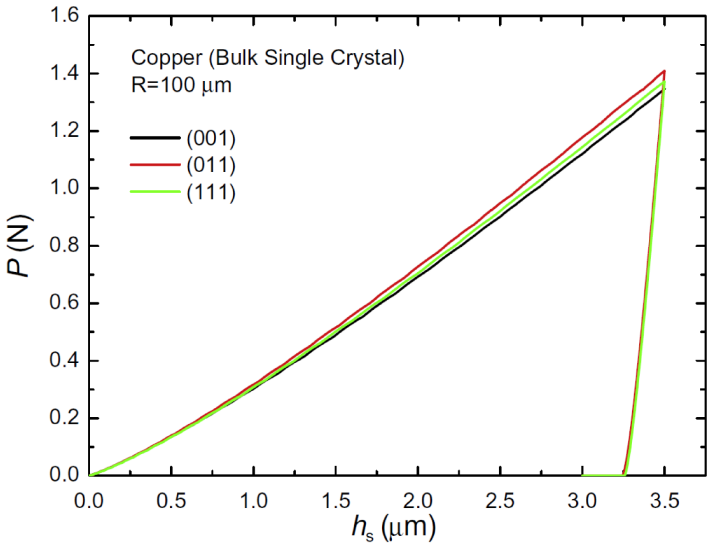

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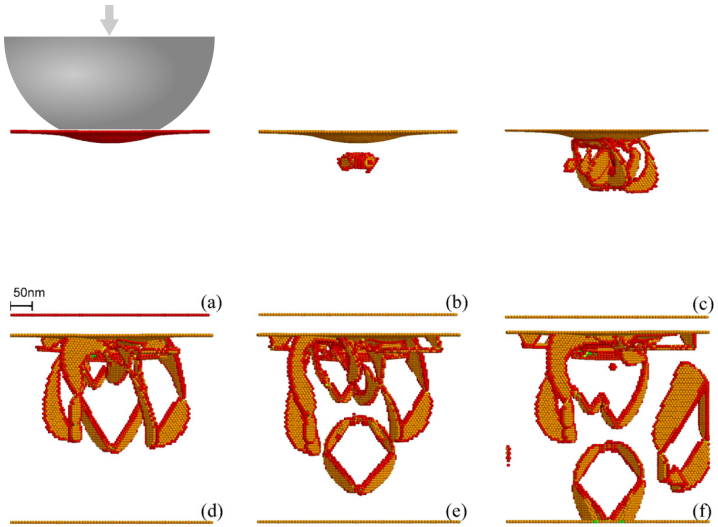

MD simulation of a spherical indentation on (111) FCC cube

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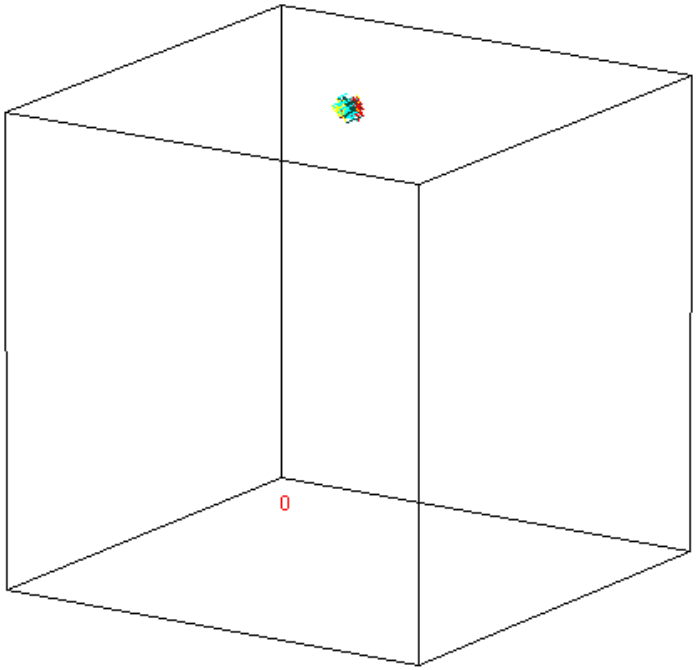

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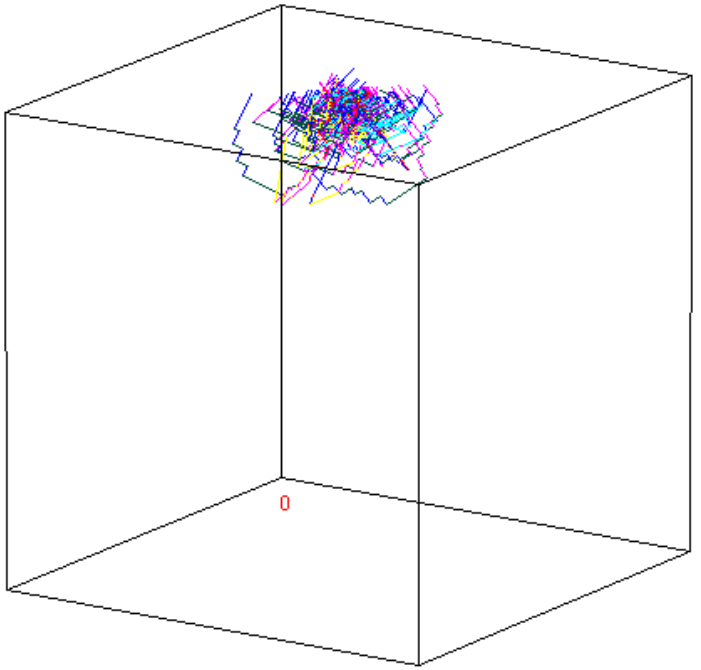

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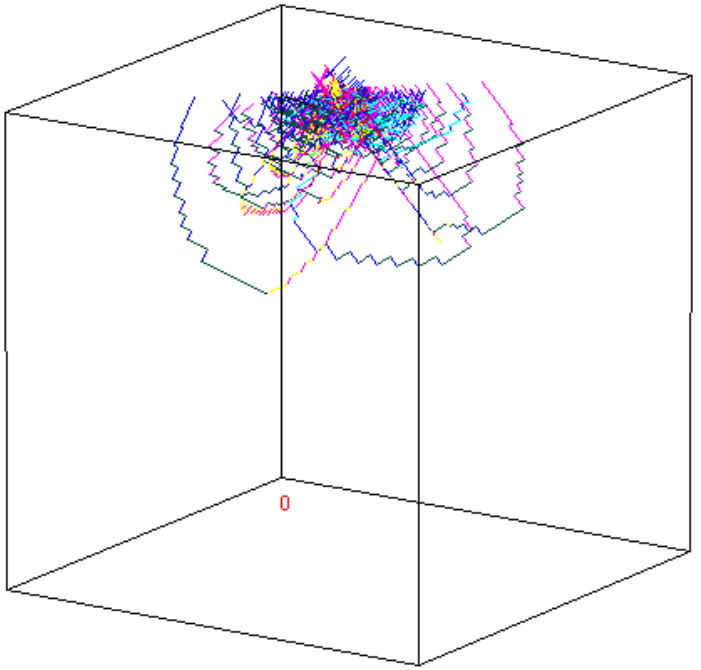

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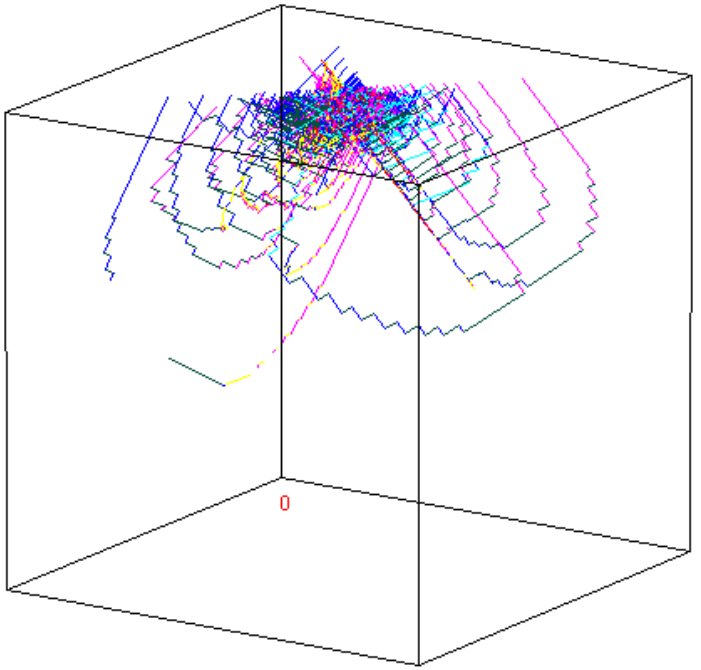

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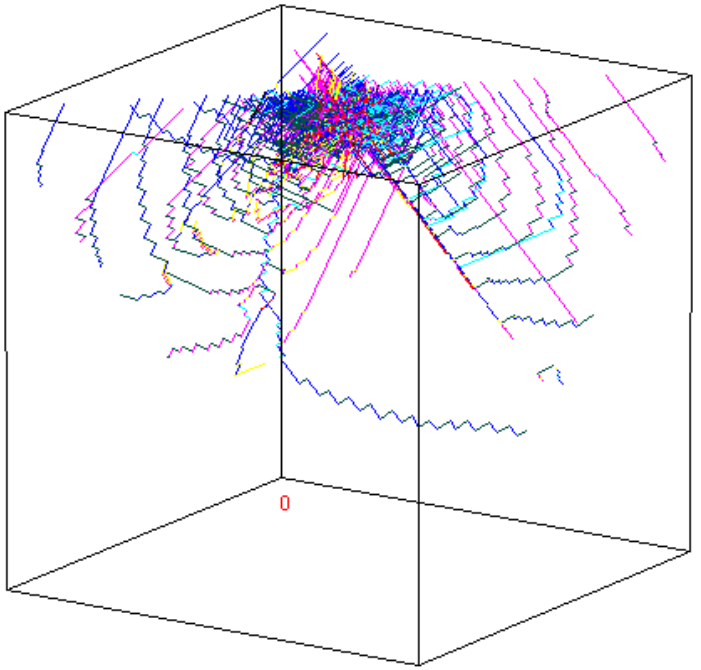

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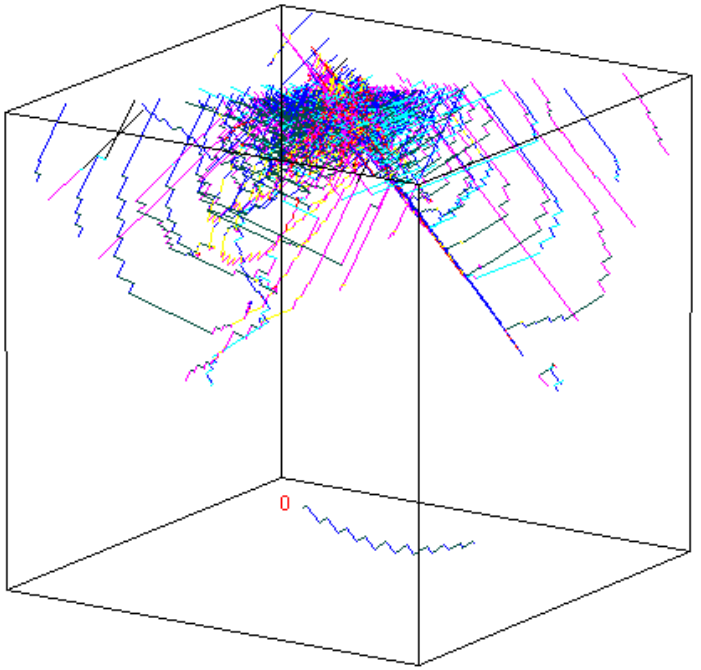

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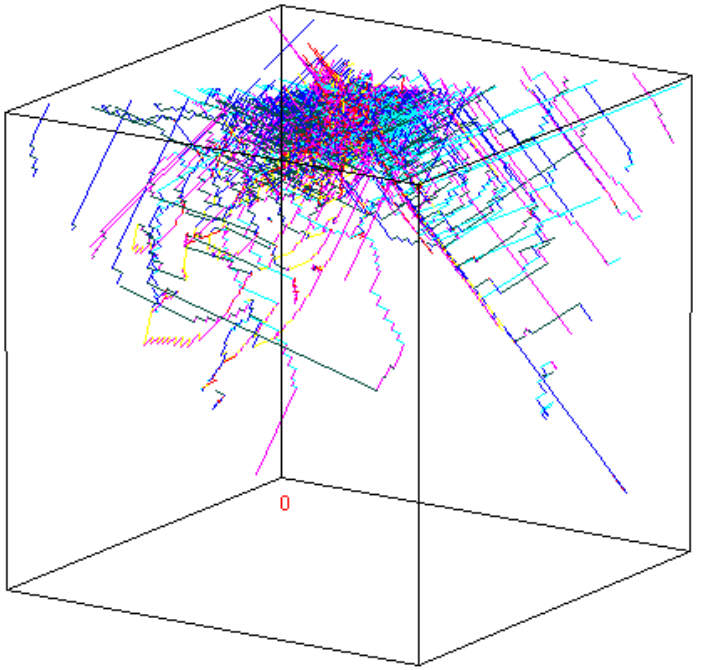

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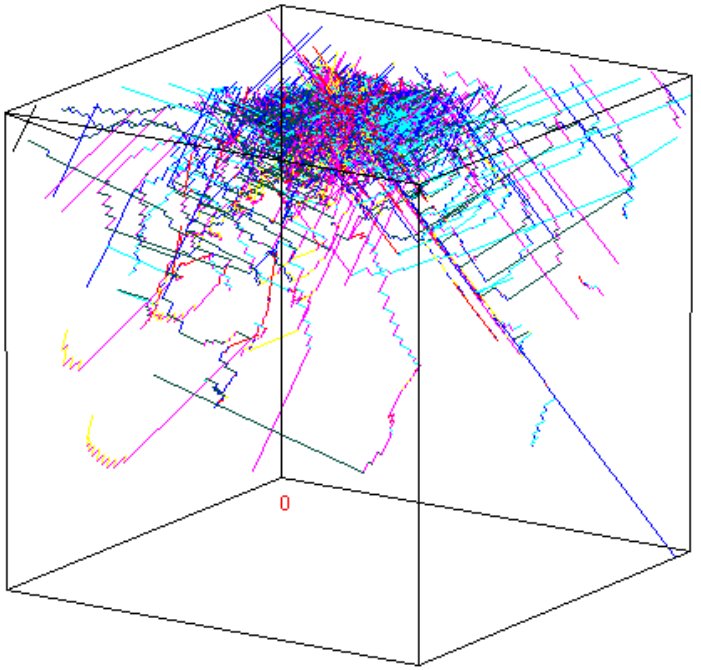

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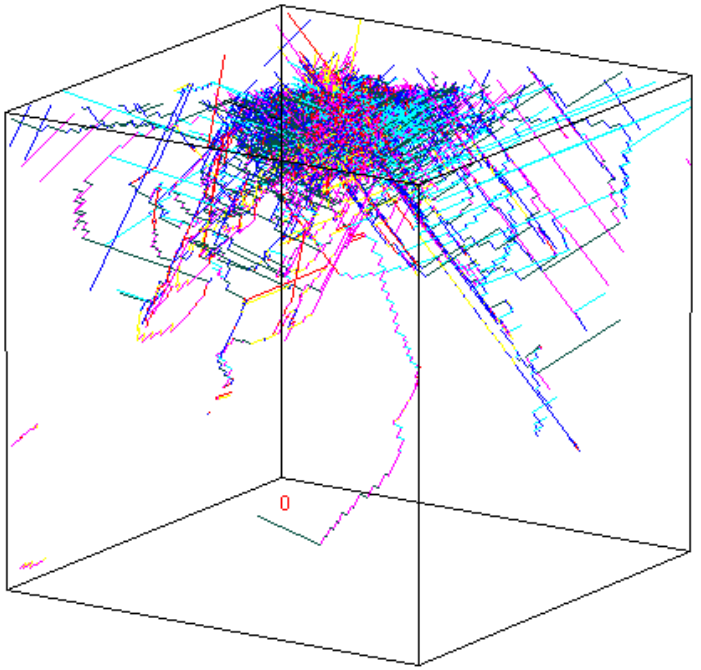

Resulting DD simulation of a spherical indentation on (111) FCC cube

[4] H.J. Chang, M. Fivel, D. Rodney, M. Verdier. Multiscale modelling of indentation in FCC metals: From atomic to continuum, CR Physique 11 (2010)
www.numodis.com

[1] Forest S., Sievert R. Elastoviscoplastic constitutive frameworks for generalized continua. Acta Mech 160 (2003)

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- Crystal plasticity models
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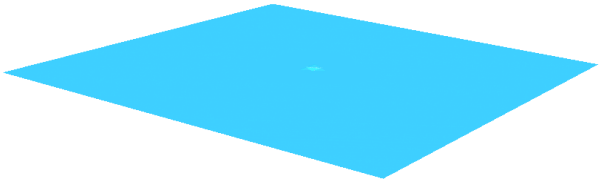

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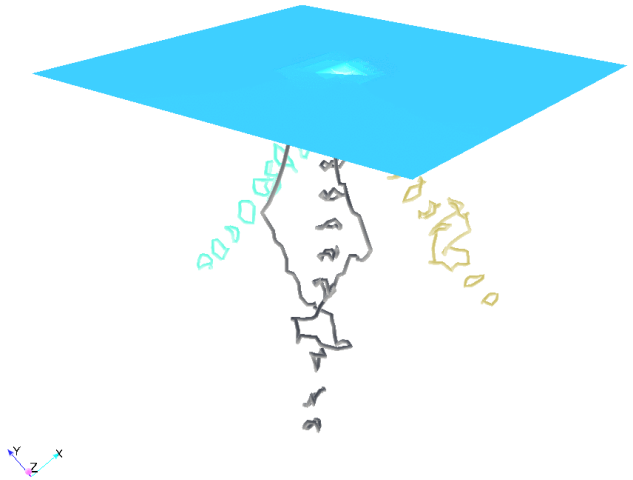

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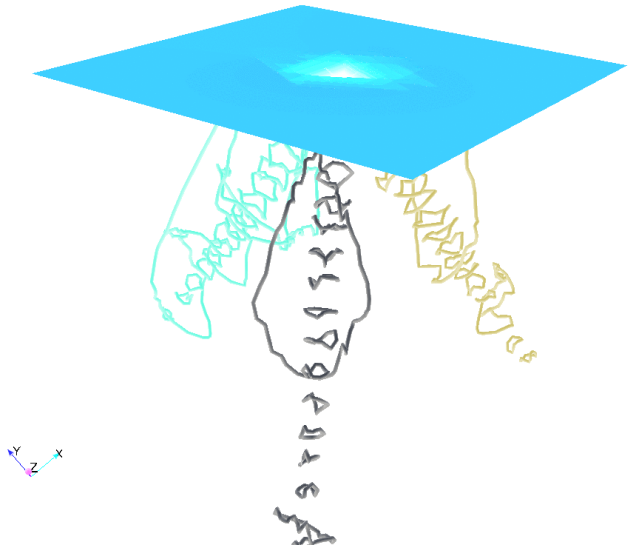

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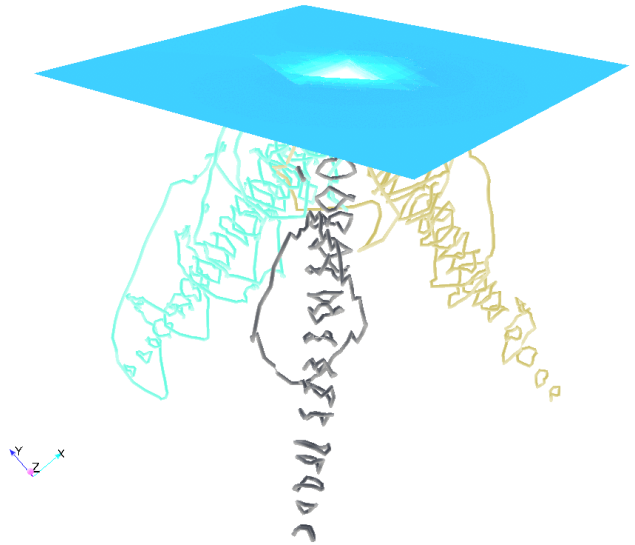

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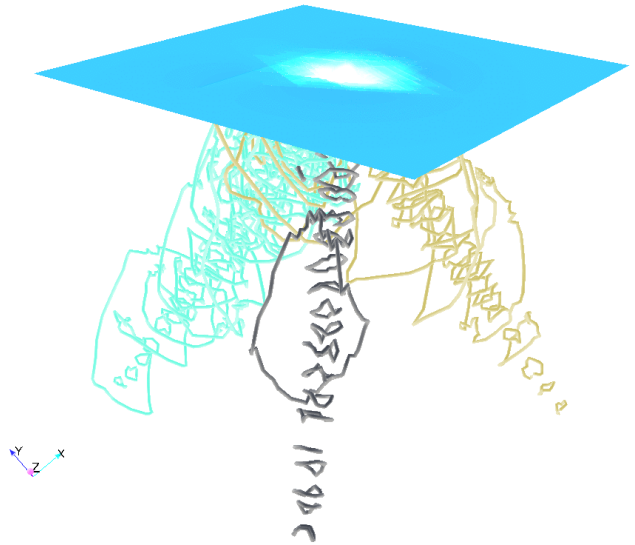

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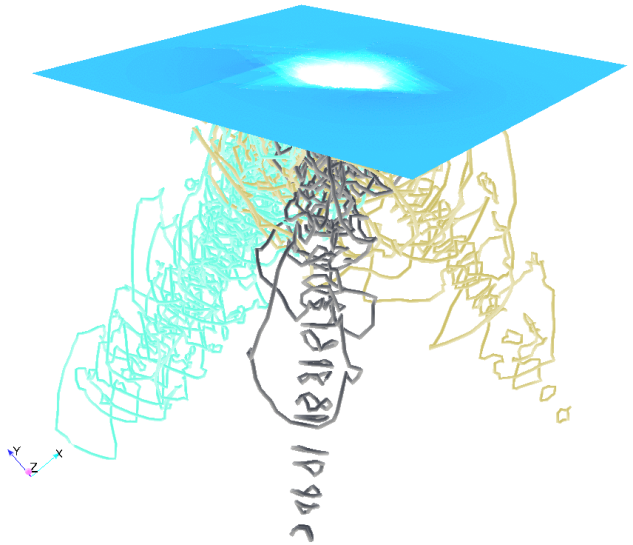

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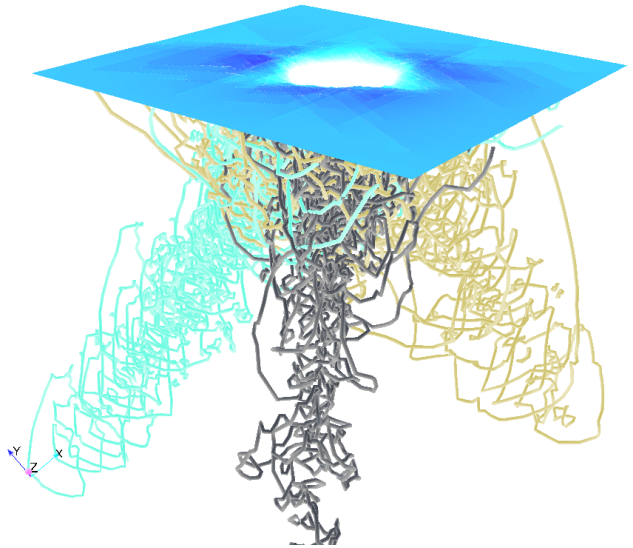

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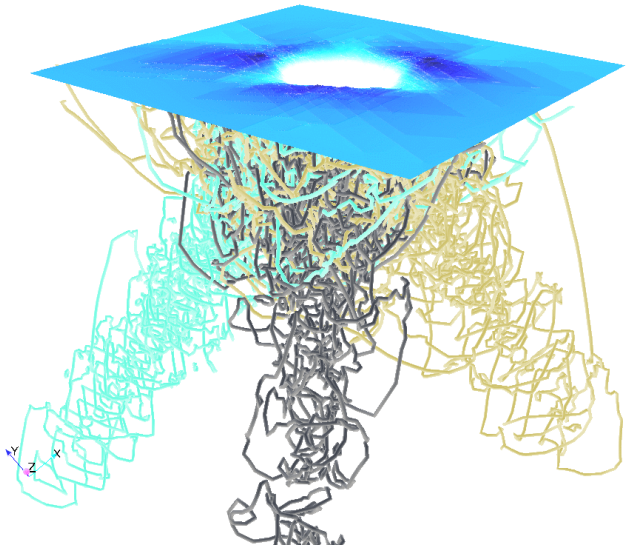

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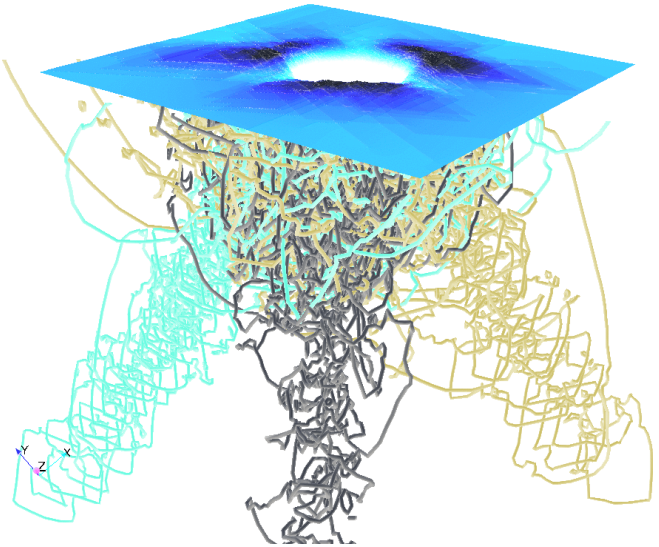

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Cosserat continuum

- Field variables (displacement & rotation): $\underline{u}, \underline{\omega}$
- Small deformation tensor: $\underline{\underline{\varepsilon}} = \nabla \underline{u} + {}^3\underline{\epsilon} \cdot \underline{\omega}$
- Torsion-curvature tensor: $\underline{\kappa} = \nabla \underline{\omega}$
- Elasticity: $\underline{\sigma} = \lambda \mathrm{tr}\left(\underline{\underline{\varepsilon}}_e\right) \underline{\underline{I}} + \mu(\underline{\underline{\varepsilon}}_e + \underline{\underline{\varepsilon}}_e^\top) + \mu_c(\underline{\underline{\varepsilon}}_e - \underline{\underline{\varepsilon}}_e^\top), \quad \underline{m} = \alpha \mathrm{tr}\left(\underline{\kappa}_e\right) \underline{\underline{I}} + 2\beta \underline{\kappa}_e$

  $$\boxed{l_e = \sqrt{\beta/\mu}}$$

  Note: $\underline{\underline{\varepsilon}}^\top \neq \underline{\underline{\varepsilon}}, \underline{\kappa}^\top \neq \underline{\kappa}, \underline{\sigma}^\top \neq \underline{\sigma}, \underline{m}^\top \neq \underline{m}$
- In non-inertial problems without volume forces and couple-forces, balance of momentum and of moment of momentum:

  $$\nabla \cdot \underline{\sigma} = 0, \quad \nabla \cdot \underline{m} - {}^3\underline{\epsilon} : \underline{\sigma} = 0$$
- Plasticity: equivalent stress[1,2] $Y = \sqrt{\frac{3}{2}\left(a_1 \underline{s} : \underline{s} + a_2 \underline{s} : \underline{s}^\top + \boxed{\frac{1}{l_p^2}} \underline{m} : \underline{m}\right)}$
- Internal lengths: elastic l_e, plastic l_p

[1] R. de Borst, L.J. Sluys, *Comp Meth Appl Mech Engin* (1991)
[2] S. Forest, R. Sievert, *Acta Mech* (2003)

---

where permutation tensor ${}^3\underline{\epsilon} \sim \epsilon_{ijk} = \begin{cases} 1, & \text{if } \{ijk\} = \{123\} \text{ or } \{231\} \text{ or } \{312\} \\ -1, & \text{if } \{ijk\} = \{321\} \text{ or } \{213\} \text{ or } \{132\} \\ 0, & \text{otherwise} \end{cases}$

Single asperity analysis

Assumptions

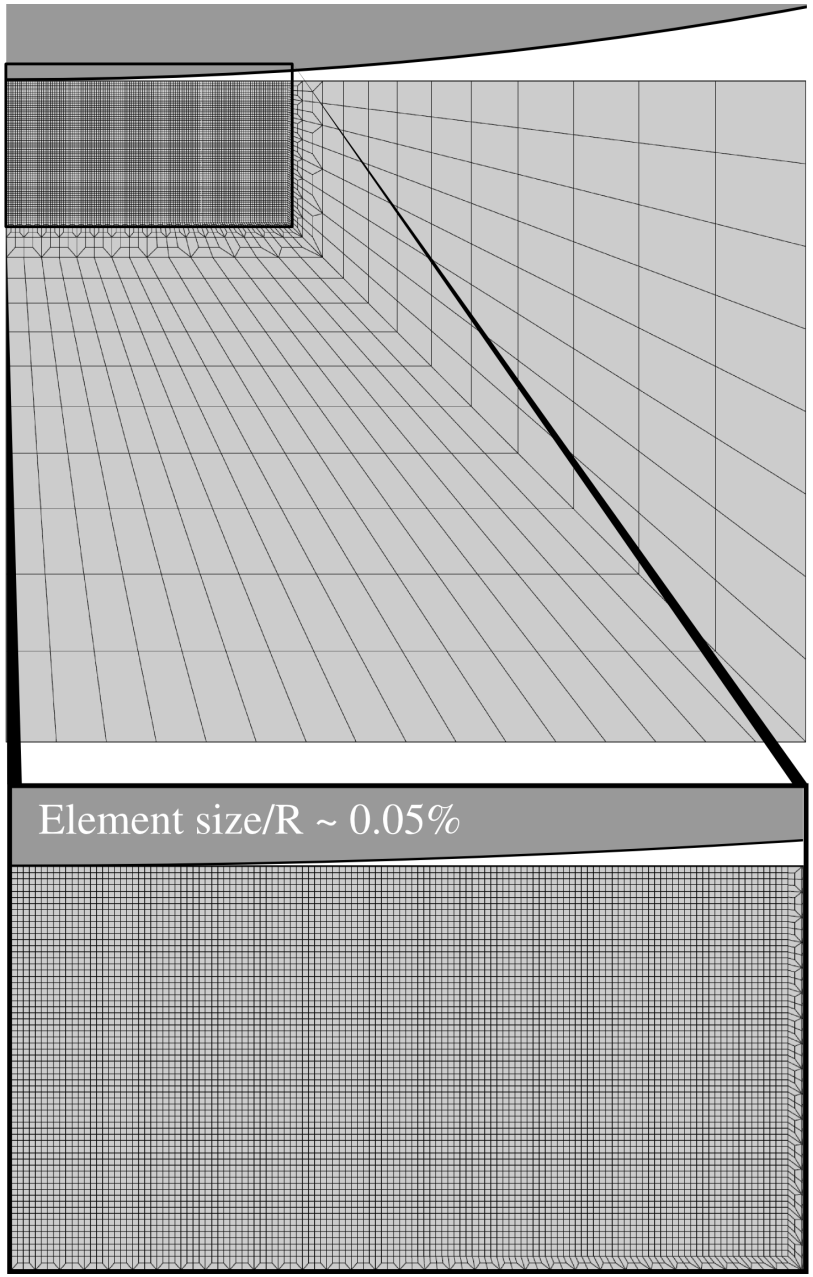
- Rigid spherical asperity
- Axisymmetric FE problem
- Generalized Cosserat continuum

Parameters

- Au: $E = 96$ GPa, $\nu = 0.42$, $\sigma_y = 140$ MPa
- $\mu_c = 10\mu, l_e = 100$ nm, $a_1 = 1$
- Indenter radius $R \in [0.002, 2000]$ μm

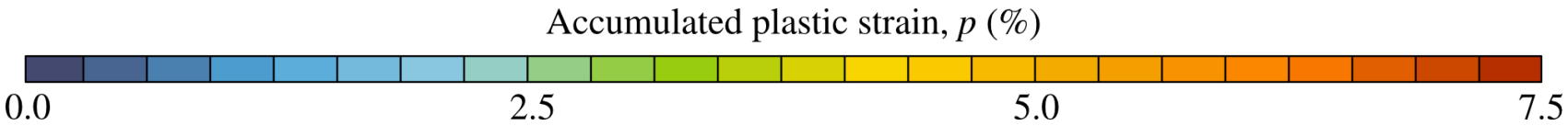
Objectives

- Study size effect
- Enhance asperity based models for rough contact



Accumulated plasticity

- Different plastic distribution



Displacement × 5

Indenter radius $R = 20\mu$m
Max plastic strain $p_{\max} \approx 7.5\%$

Displacement × 5

Indenter radius $R = 2\mu$m
Max plastic strain $p_{\max} \approx 11\%$

Accumulated plasticity

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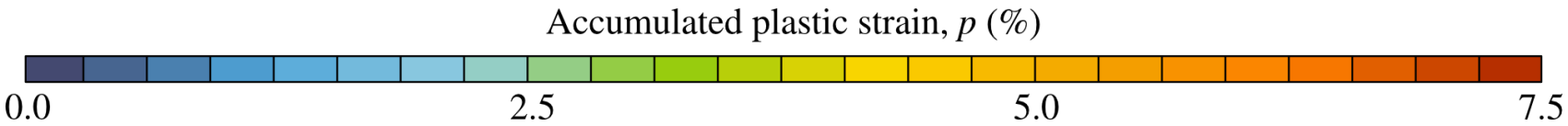

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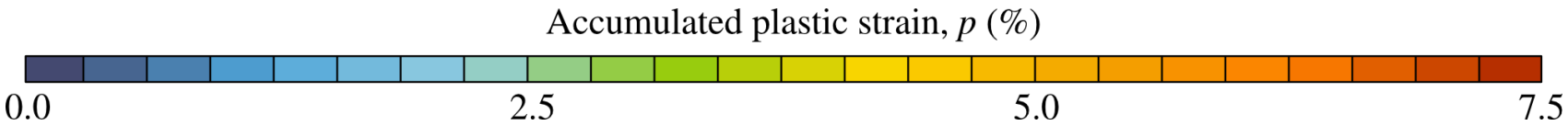

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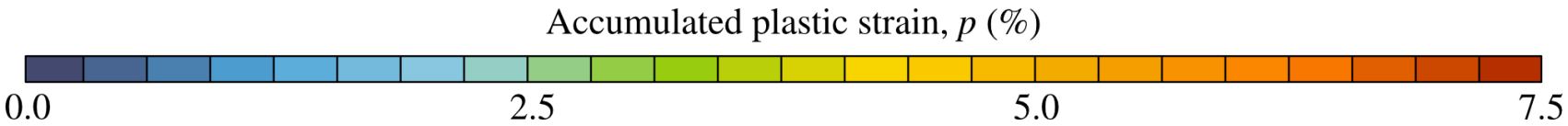

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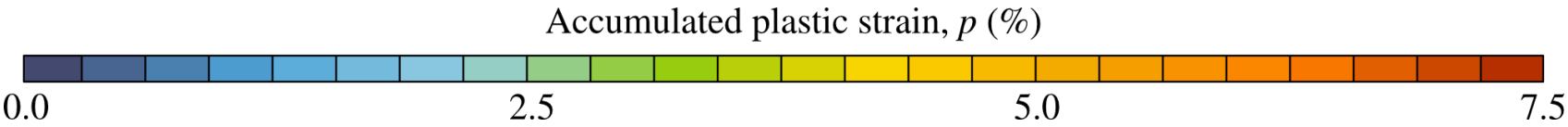

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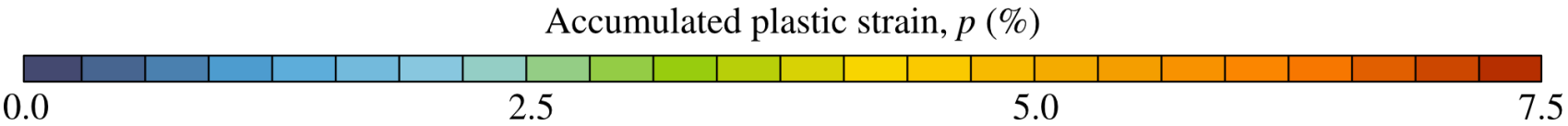

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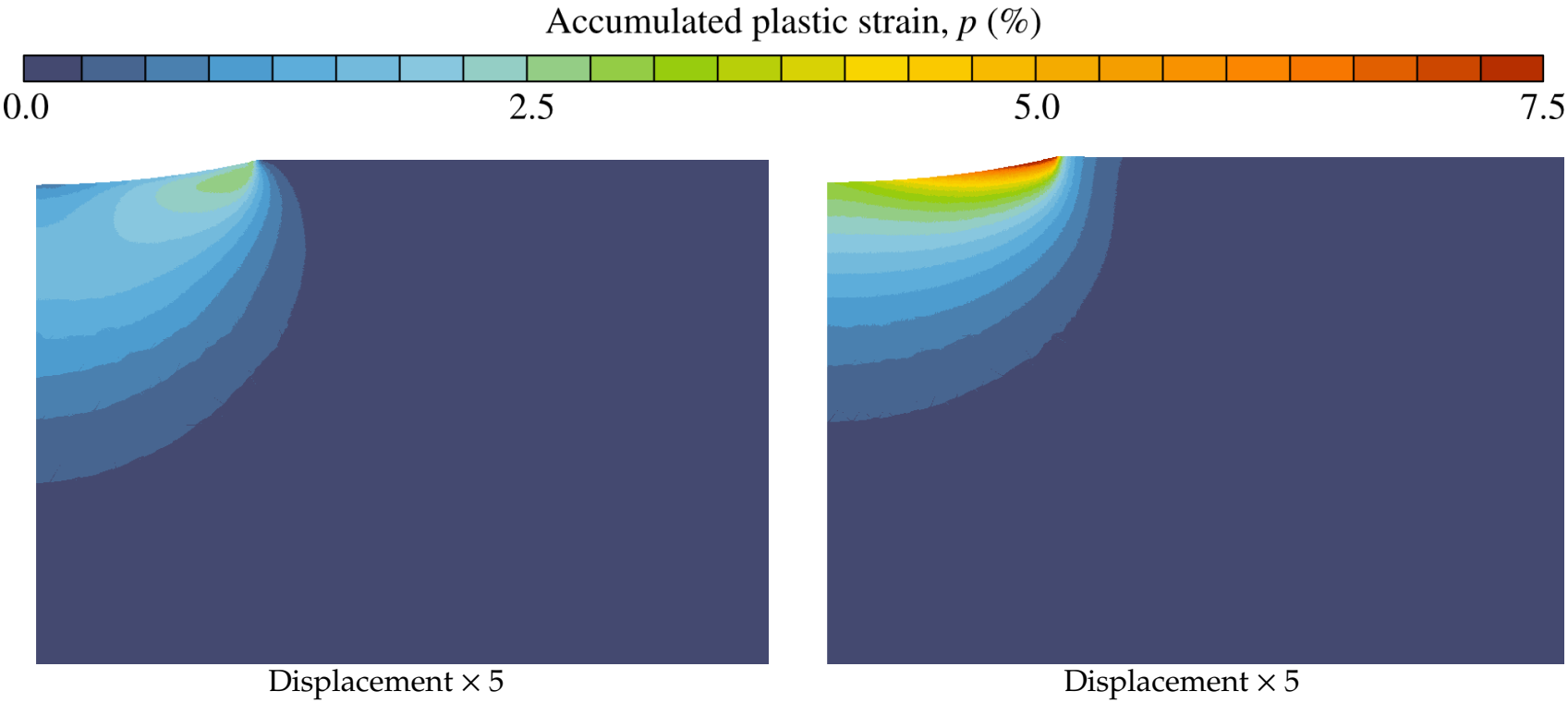

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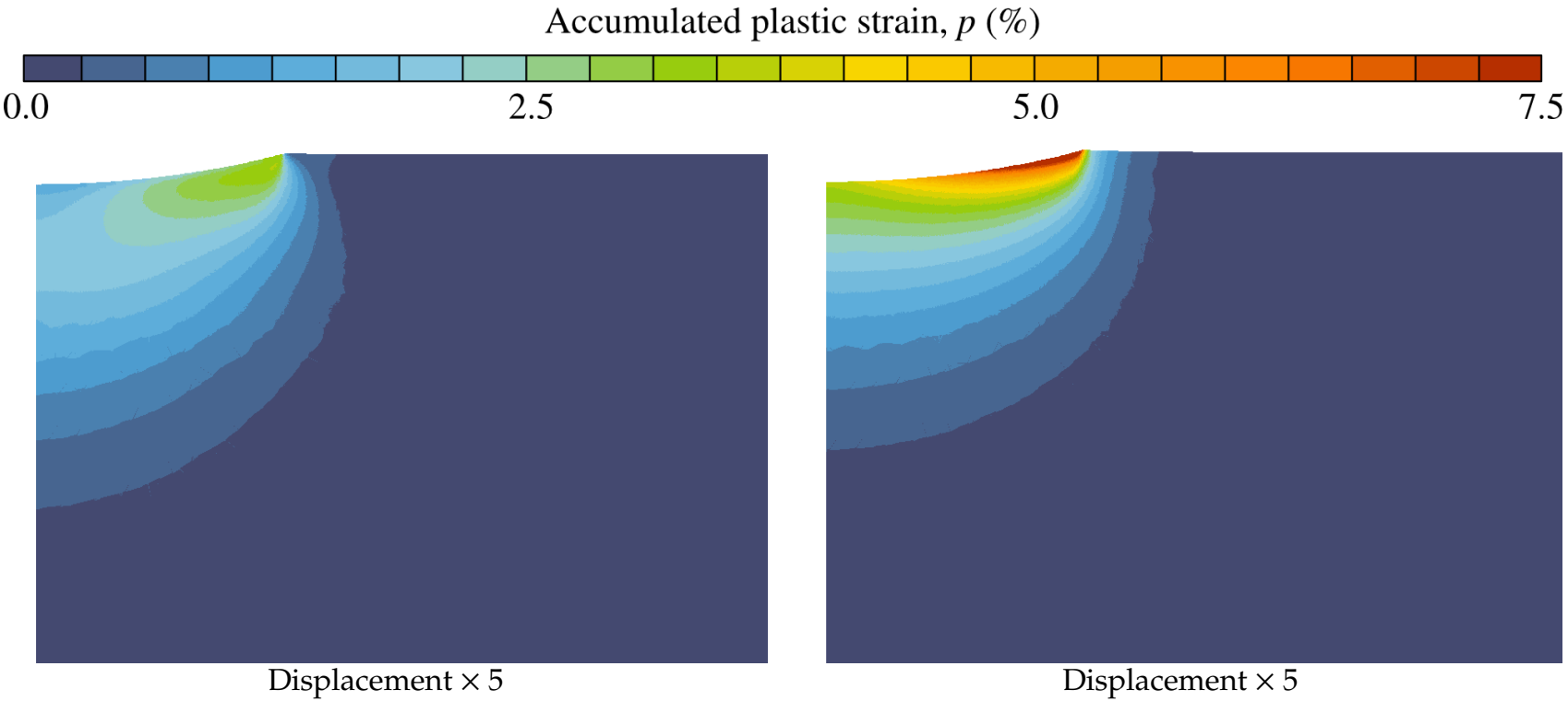

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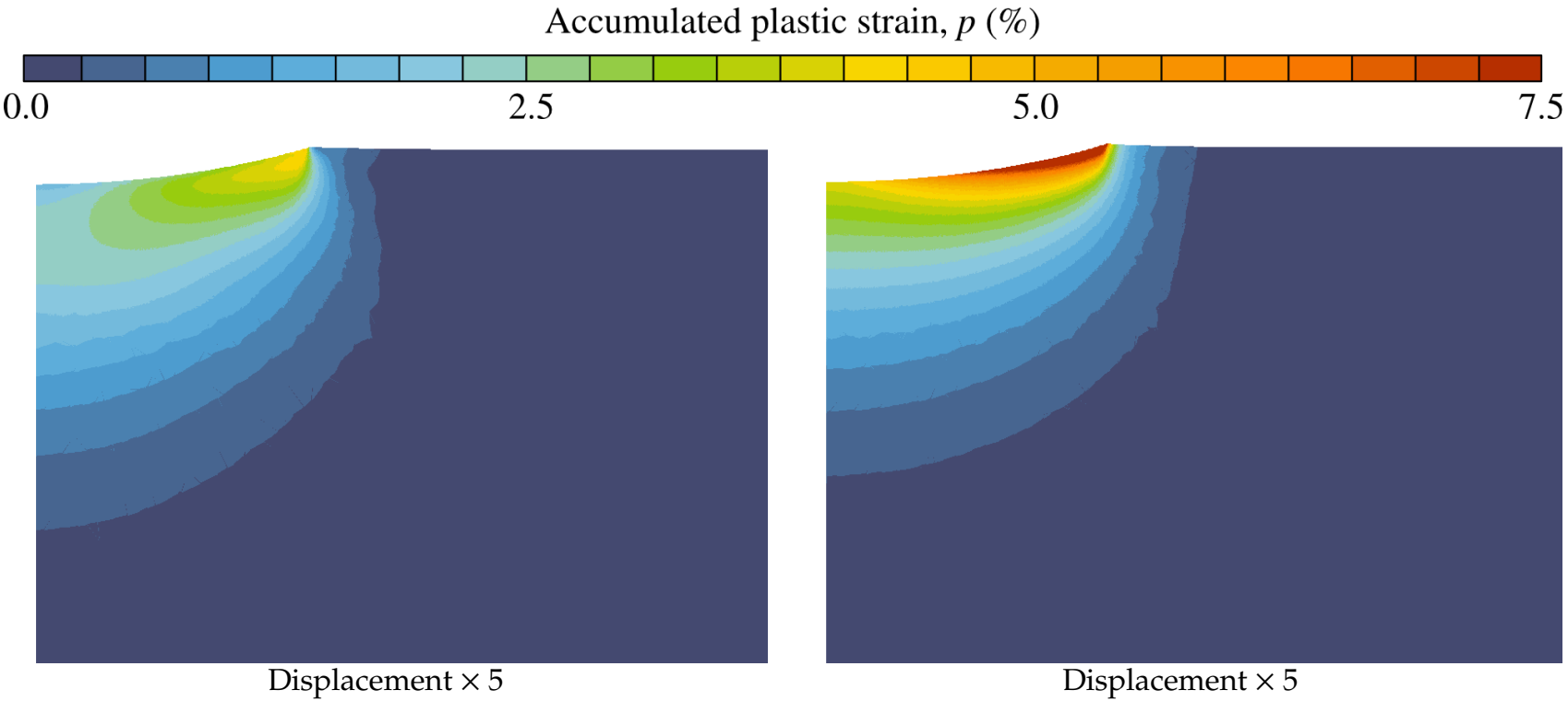

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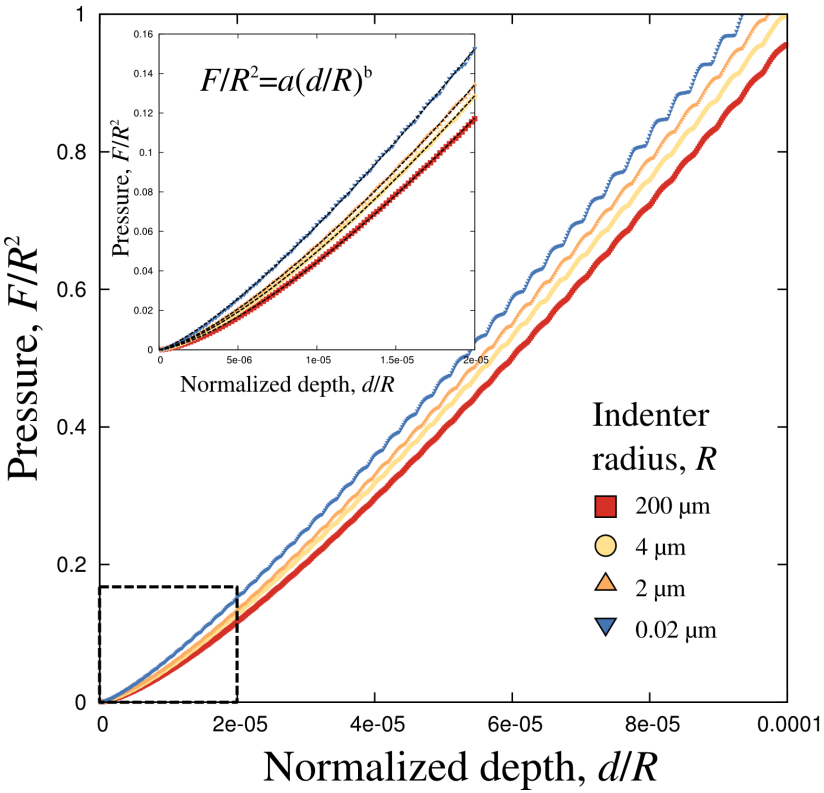
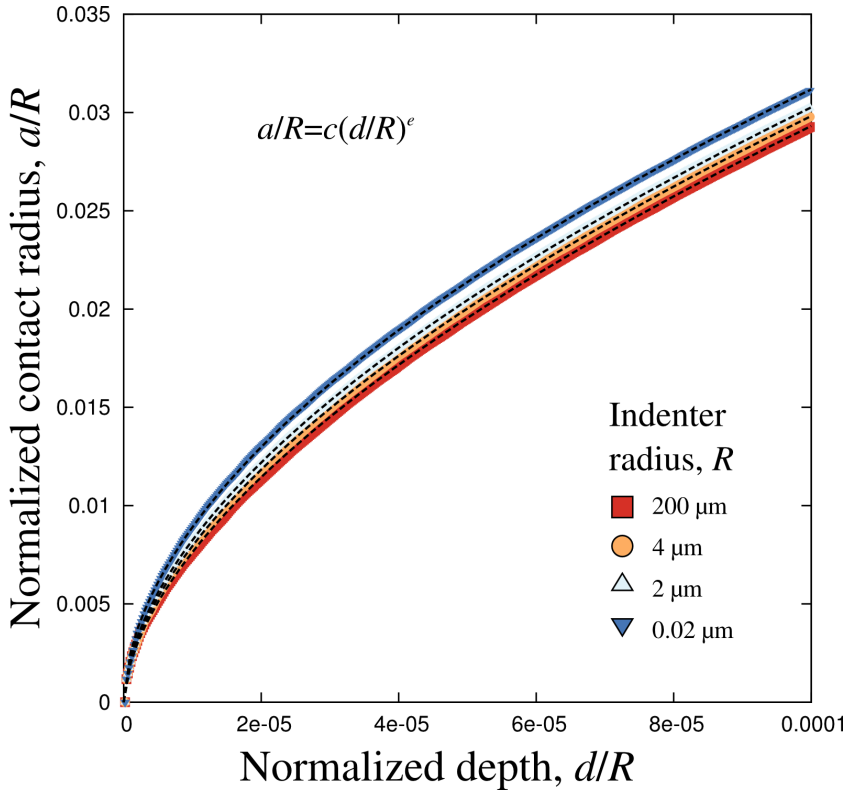
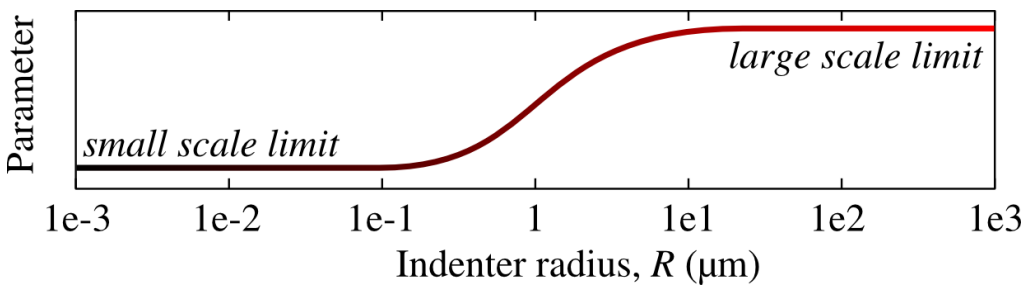
Displacement–force–contact radius



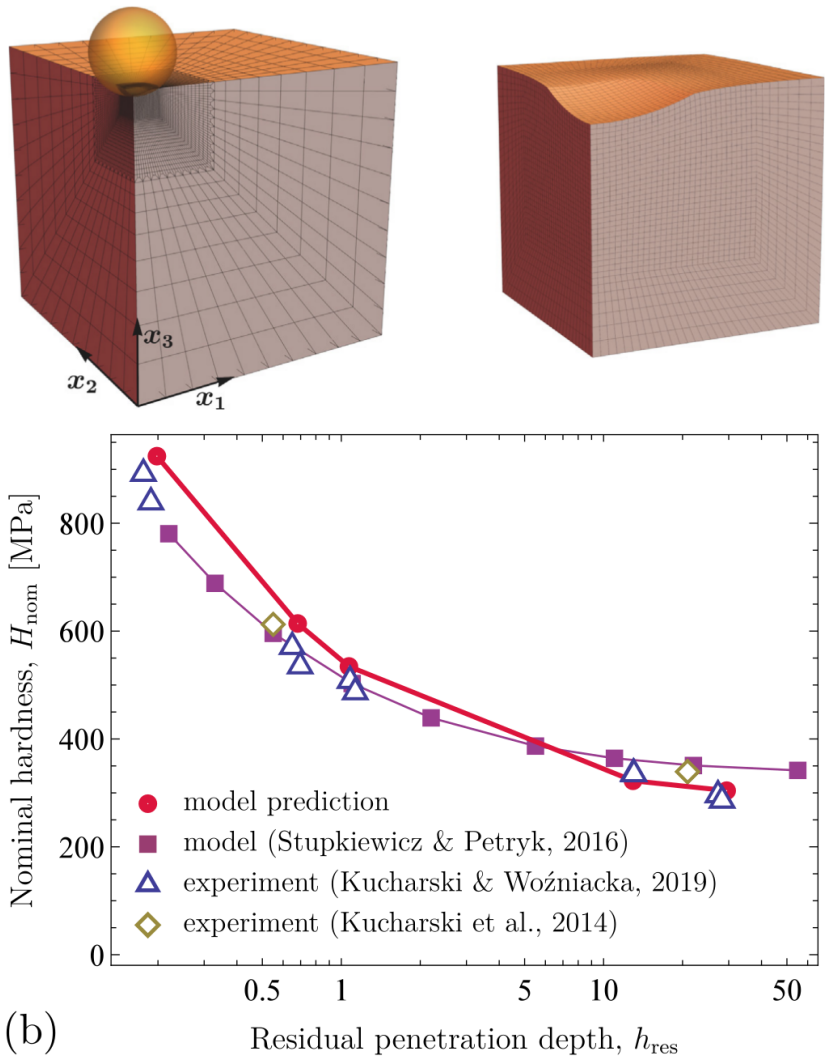


Parameters:
$a(R), b(R), c(R), e(R)$
`else:` using NURBS curve fit.



Working model

- Gradient-enhanced hardening with the Cosserat crystal plasticity theory[1]
- Cubic elasticity with C_{11}, C_{12}, C_{44} is to be taken into account



[1] Ryś, M., Stupkiewicz, S. and Petryk, H., 2022. Micropolar regularization of crystal plasticity with the gradient-enhanced incremental hardening law. International Journal of Plasticity, 156, p.103355.

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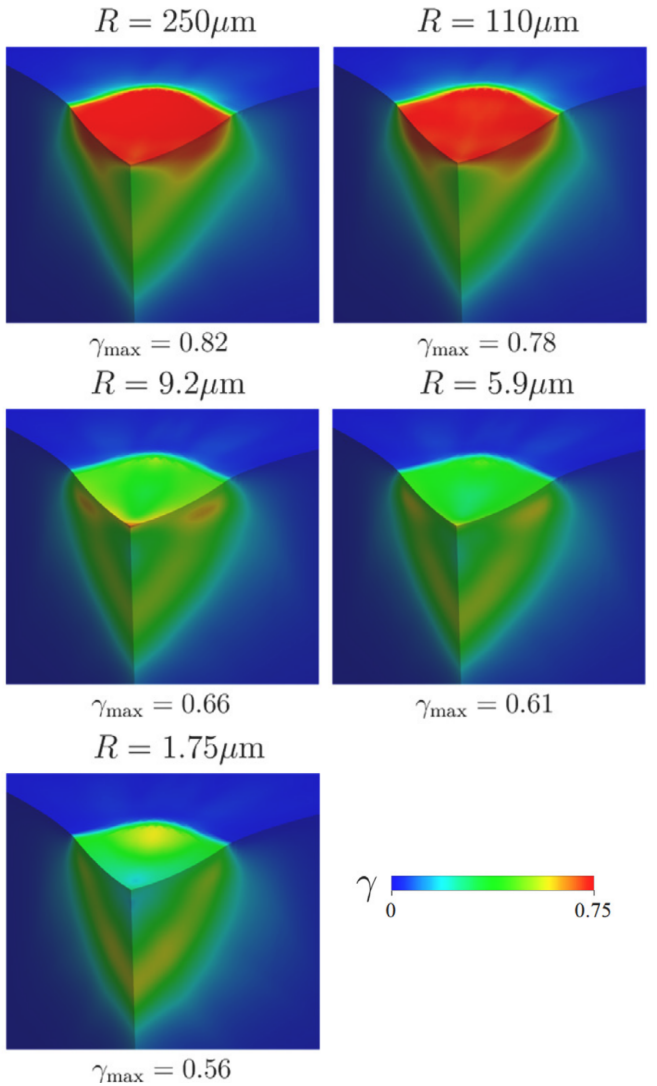


Residual plastic slip for different indenter radii R

[1] Ryś, M., Stupkiewicz, S. and Petryk, H., 2022. Micropolar regularization of crystal plasticity with the gradient-enhanced incremental hardening law. International Journal of Plasticity, 156, p.103355.

Viscoelasticity

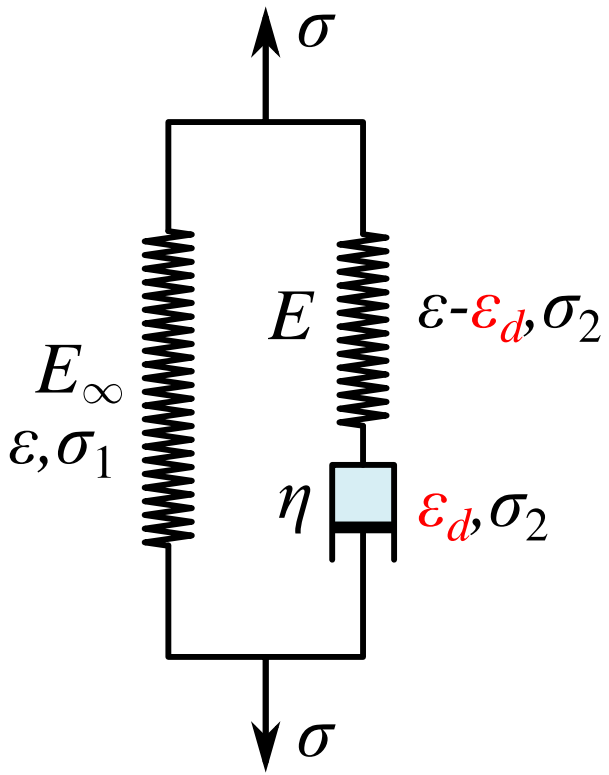
One-dimensional constitutive equations

- Applied stress σ
- In the left branch $\sigma_1 = E_\infty \varepsilon$
- In the dashpot $\sigma_2 = \eta \dot{\varepsilon}_d \quad (*)$
- In the right spring $\sigma_2 = E(\varepsilon - \varepsilon_d) \quad (**)$
- For the whole system $\sigma = \sigma_1 + \sigma_2$

$$\boxed{\sigma = (E_\infty + E)\varepsilon - E\varepsilon_d}$$

- From $(*)$ and $(**)$, and denoting $\tau = \eta/E$:

$$\boxed{\dot{\varepsilon}_d + \frac{\varepsilon_d}{\tau} = \frac{\varepsilon}{\tau},} \qquad \varepsilon_d \xrightarrow[t \to -\infty]{} 0$$



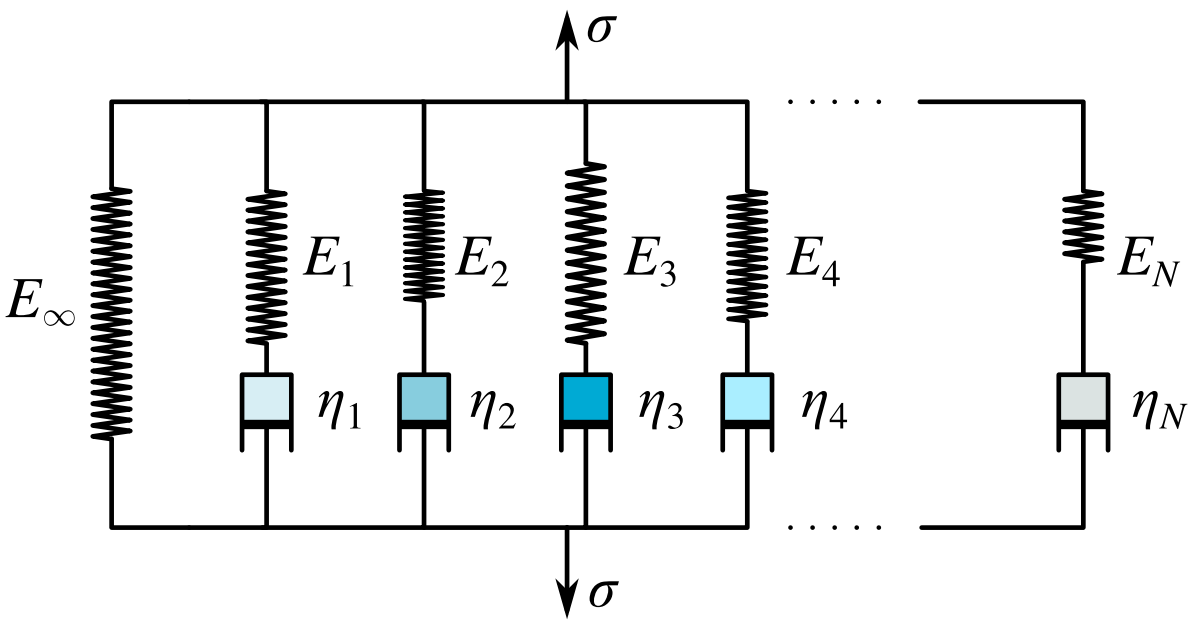
• One-dimensional viscoelastic model

- Recall: 1D model

$$\sigma = (E_\infty + E)\varepsilon - E\varepsilon_d, \quad \dot{\varepsilon}_d + \frac{\varepsilon_d}{\tau} = \frac{\varepsilon}{\tau}, \quad \varepsilon_d \xrightarrow[t\to-\infty]{} 0$$

- Multiple dashpots in parallel

$$\sigma = \underbrace{(E_\infty + \sum_i E_i)\varepsilon}_{\text{elastic stress } \sigma_0} - \sum_i E_i \varepsilon_d^i, \quad \dot{\varepsilon}_d^i + \frac{\varepsilon_d^i}{\tau_i} = \frac{\varepsilon}{\tau_i}, \quad \varepsilon_d^i \xrightarrow[t\to-\infty]{} 0, \quad \tau_i = \frac{\eta_i}{E_i}$$



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- Denote $E_0 = E_\infty + \sum_i E_i$, $\psi_i = E_i/E_0$, and $q_i = E_i\varepsilon_d^i$ we obtain

$$\sigma = E_0\varepsilon - \sum_i q_i, \quad \dot{q}_i + \frac{q_i}{\tau_i} = \frac{\psi_i}{\tau_i}\sigma_0, \quad \varepsilon_d^i \xrightarrow[t\to-\infty]{} 0$$

- By construction

$$\sum_i \psi_i + \frac{E_\infty}{E_0} = 1 \quad \Rightarrow \quad \sum_i \psi_i = 1 - \frac{E_\infty}{E_0}$$

• Three-dimensional viscoelastic model

Linear viscoelastic (generalized Maxwell model, standard solid)

- Stress-strain relation:

$$\underline{\underline{\sigma}}(t) = K\theta\underline{\underline{I}} + \int_{-\infty}^{t} G(t-\tau)\dot{\underline{\underline{e}}}(\tau)d\tau,$$

- Kernel $G(\tau)$ is given by:

$$G(\tau) = 2G_\infty + 2(G_0 - G_\infty)\Psi(\tau) \text{ with } \Psi(\tau) = \sum_{i=1}^{n} \psi_i \exp(-\tau/\tau_i)$$

  - G_∞, G_0 are the slow/fast loading shear moduli, respectively, such that $G_\infty \leq G_0$;
  - K is the bulk modulus, and for elastomers/polymers $K/G_0 \gg 1$;
  - ψ_i are the influence coefficients, such that $\sum_{i=1}^{n} \psi_i = 1$;
  - τ_i are the respective relaxation times.

• Material model: *storage* and loss moduli

- Consider a harmonic (rigid) loading: $\underline{e}(t) = \underline{e}_0 \exp(i\omega t)$
- Split the kernel: $G(t) = 2G_\infty + \tilde{G}(t)$
- Then, the storage modulus (general case):

$$G'(\omega) = 2G_\infty + \omega \int_0^\infty \tilde{G}(\tau) \sin(\omega\tau)\, d\tau$$

- The storage modulus in the framework of the generalized Maxwell model:

$$G'(\omega) = 2G_\infty + 2\omega(G_0 - G_\infty) \sum_{i=1}^{n} \psi_i \int_0^\infty \exp(-\tau/\tau_i) \sin(\omega\tau)\, d\tau$$

$$G'(\omega) = 2G_\infty + 2(G_0 - G_\infty) \sum_{i=1}^{n} \frac{\psi_i \omega^2 \tau_i^2}{1 + \omega^2 \tau_i^2}$$

- Remark:

$$\int \exp(cx) \sin(bx)\, dx = \frac{\exp(cx)}{c^2 + b^2} [c \sin(bx) - b \cos(bx)]$$

• Material model: storage and *loss* moduli

- The loss modulus (general case):

$$G''(\omega) = \omega \int_0^\infty \tilde{G}(\tau) \cos(\omega\tau)\, d\tau$$

- The loss modulus in the framework of the generalized Maxwell model:

$$G''(\omega) = 2\omega(G_0 - G_\infty) \sum_{i=1}^{n} \psi_i \int_0^\infty \exp(-\tau/\tau_i) \cos(\omega\tau)\, d\tau$$

$$G''(\omega) = 2(G_0 - G_\infty) \sum_{i=1}^{n} \frac{\psi_i \omega\tau_i}{1 + \omega^2\tau_i^2}.$$

- Remark:

$$\int \exp(cx) \cos(bx)\, dx = \frac{\exp(cx)}{c^2 + b^2} [c \cos(bx) + b \sin(bx)]$$

• Material model: example

- Material parameters: $G_0 = 1.1$ MPa, $G_\infty = 50$ kPa
- Single relaxation time: $\tau_0 = 10^{-7}$ s
- Quasi-incompressible material: $K/G_0 = 10^6 \gg 1$
- Uniaxial (rigid) loading: $\varepsilon_{xx} = A\sin(\omega t)$, $\sigma_{yy} = \sigma_{zz} = 0$, $\varepsilon_{yy} = \varepsilon_{zz} \approx -0.5\varepsilon_{xx}$
- Spherical and deviatoric parts: $\underline{\boldsymbol{\epsilon}} \approx A(1-2\nu)\sin(\omega t)\underline{\boldsymbol{I}}$, $\underline{\boldsymbol{e}} \approx \underline{\boldsymbol{\varepsilon}}$
- Stress-strain relation:

$$\underline{\boldsymbol{\sigma}}(t) = \int_{-\infty}^{t} 2(G_0 - G_\infty)\exp[-(t-\tau)/\tau_0]\underline{\dot{\boldsymbol{e}}}(\tau)d\tau + 2G_\infty\underline{\boldsymbol{e}} + K\underline{\boldsymbol{\epsilon}},$$

- Axial and radial stress components:

$$\sigma_{xx} = 2G_\infty\varepsilon_{xx} + K(\varepsilon_{xx} + 2\varepsilon_{yy}) + \int_{-\infty}^{t} 2(G_0 - G_\infty)\exp[-(t-\tau)/\tau_0]\dot{\varepsilon}_{xx}(\tau)d\tau$$

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$$\sigma_{yy} = 0$$

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$$\sigma_{yy} = 0 = 2G_\infty\varepsilon_{yy} + K(\varepsilon_{xx} + 2\varepsilon_{yy}) + \int_{-\infty}^{t} 2(G_0 - G_\infty)\exp[-(t-\tau)/\tau_0]\dot{\varepsilon}_{yy}(\tau)d\tau$$

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- Axial and radial stress components:

$$\sigma_{xx} = 2G_\infty\varepsilon_{xx} + K(\varepsilon_{xx} + 2\varepsilon_{yy}) + \int_{-\infty}^{t} 2(G_0 - G_\infty)\exp[-(t-\tau)/\tau_0]\dot{\varepsilon}_{xx}(\tau)d\tau$$

$$\sigma_{yy} = 0 = 2G_\infty\varepsilon_{yy} + K(\varepsilon_{xx} + 2\varepsilon_{yy}) + \int_{-\infty}^{t} 2(G_0 - G_\infty)\exp[-(t-\tau)/\tau_0]\dot{\varepsilon}_{yy}(\tau)d\tau$$

$$\boxed{\sigma_{xx} = 3G_\infty\varepsilon_{xx} + \int_{-\infty}^{t} 3(G_0 - G_\infty)\exp[-(t-\tau)/\tau_0]\dot{\varepsilon}_{xx}(\tau)d\tau}$$

• Material model: example II

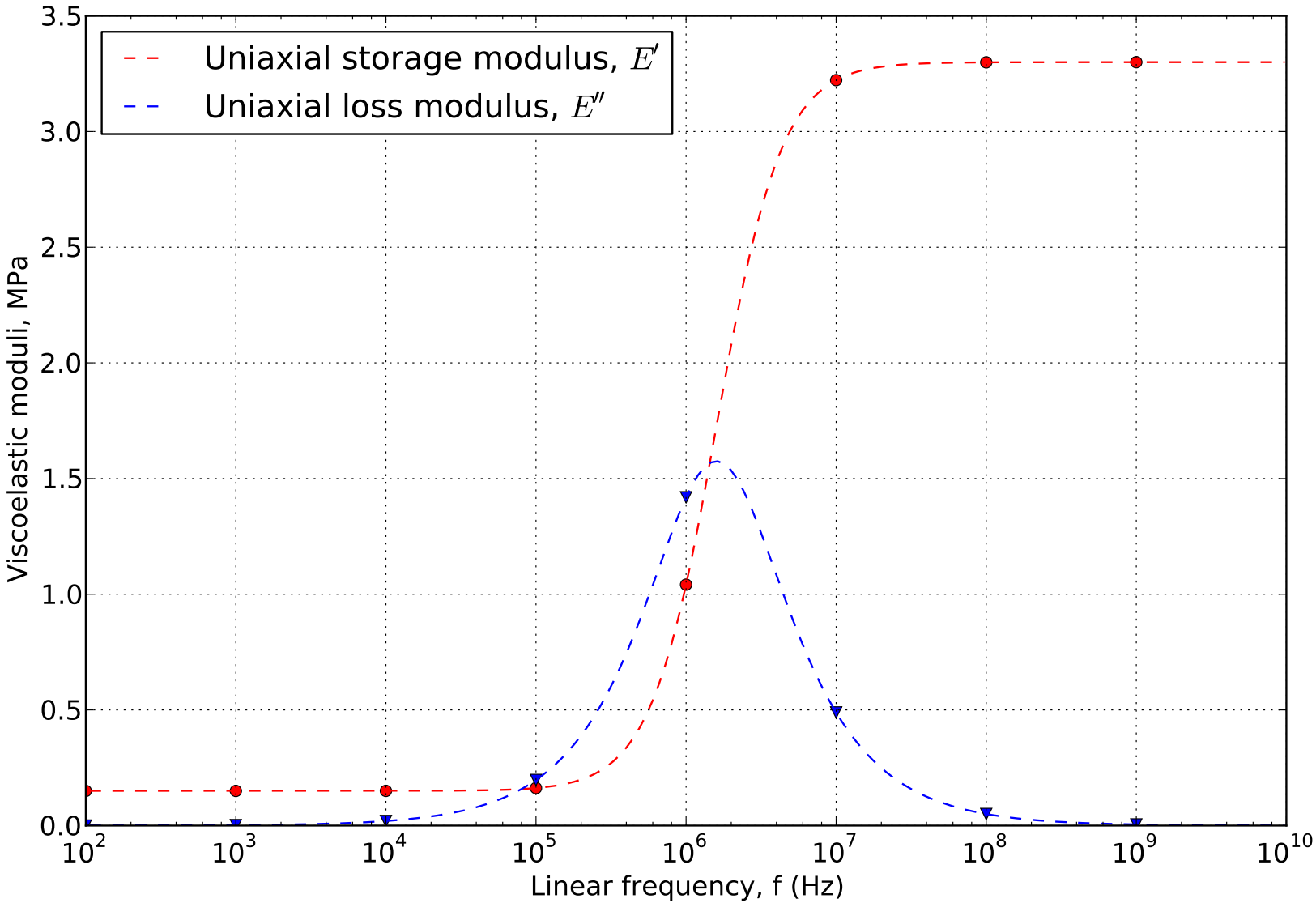
- Uniaxial storage modulus:

$$E'(\omega) = 3G_\infty + 3(G_0 - G_\infty)\frac{\omega^2\tau_0^2}{1+\omega^2\tau_0^2}$$

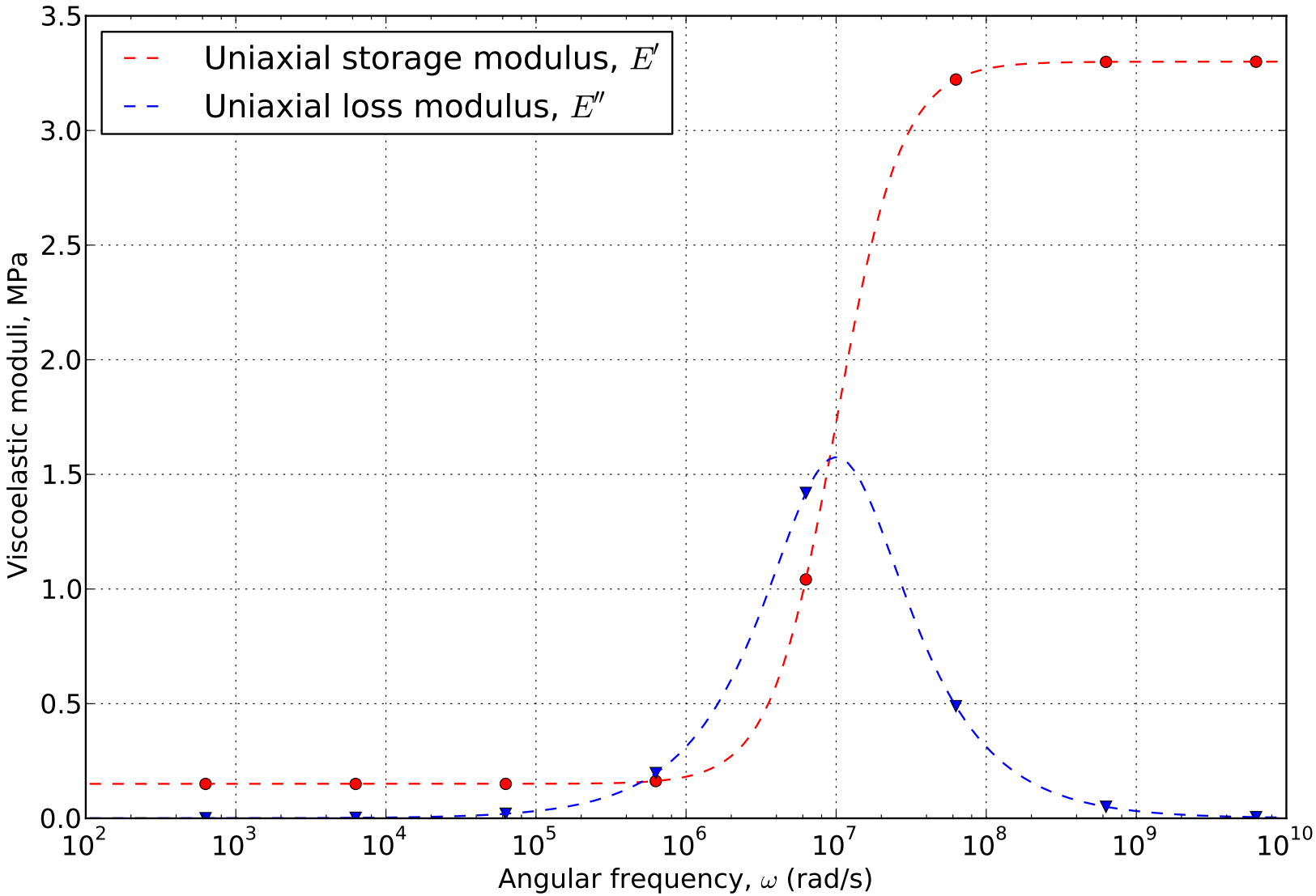
- Uniaxial loss modulus:

$$E''(\omega) = 3(G_0 - G_\infty)\frac{\omega\tau}{1+\omega^2\tau_0^2}$$

• Material model: example (FEA vs Analytics)



• Material model: example (FEA vs Analytics)



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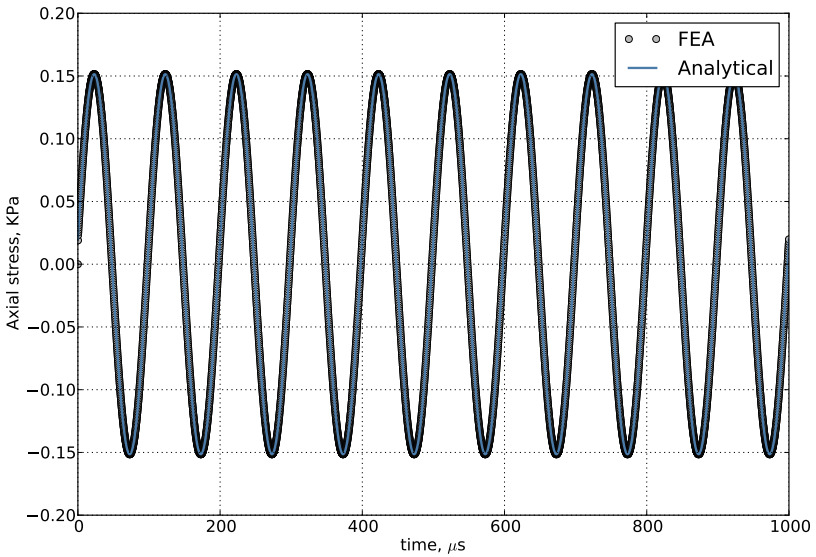
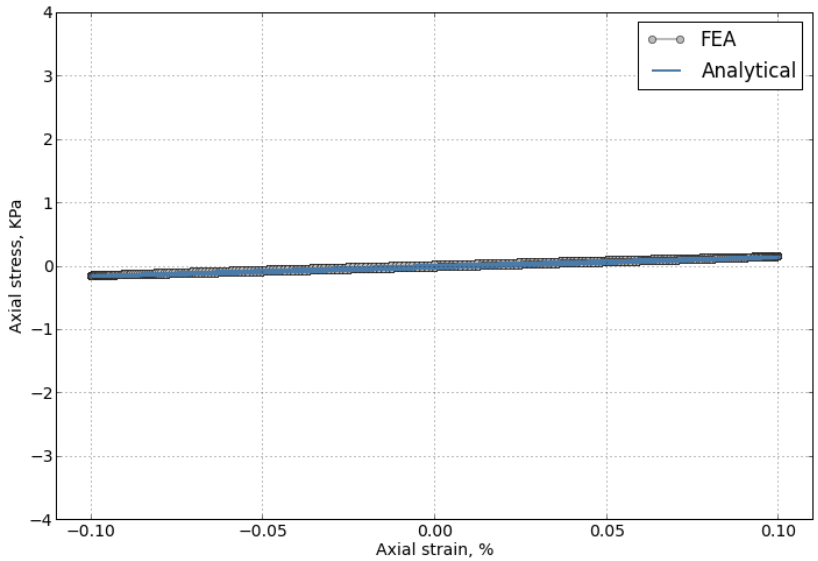
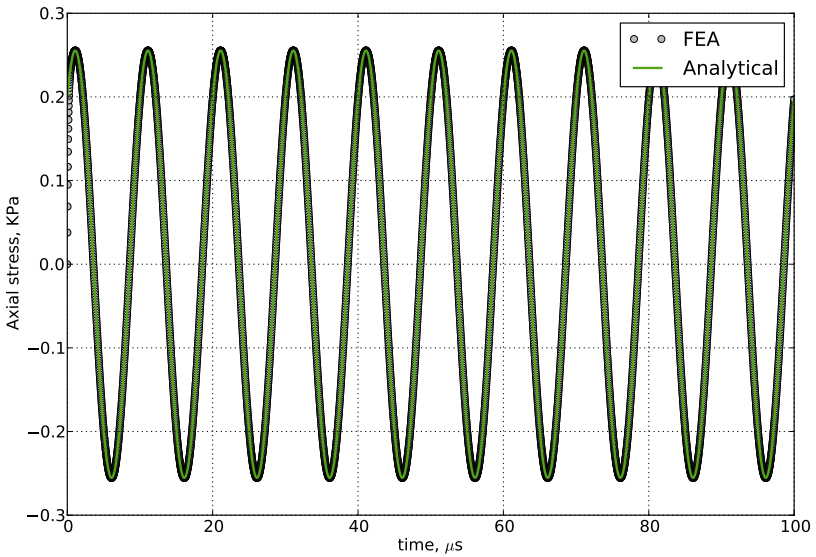
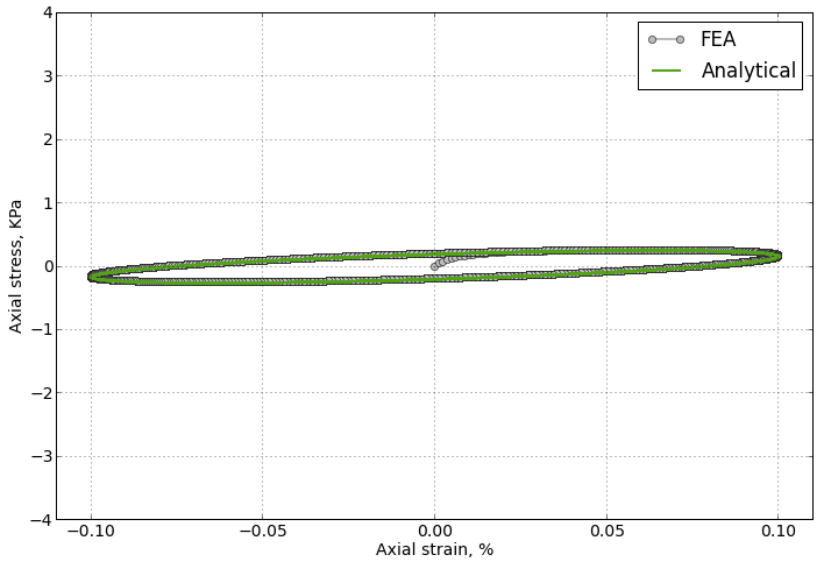




Linear frequency $f = 10^4$ Hz

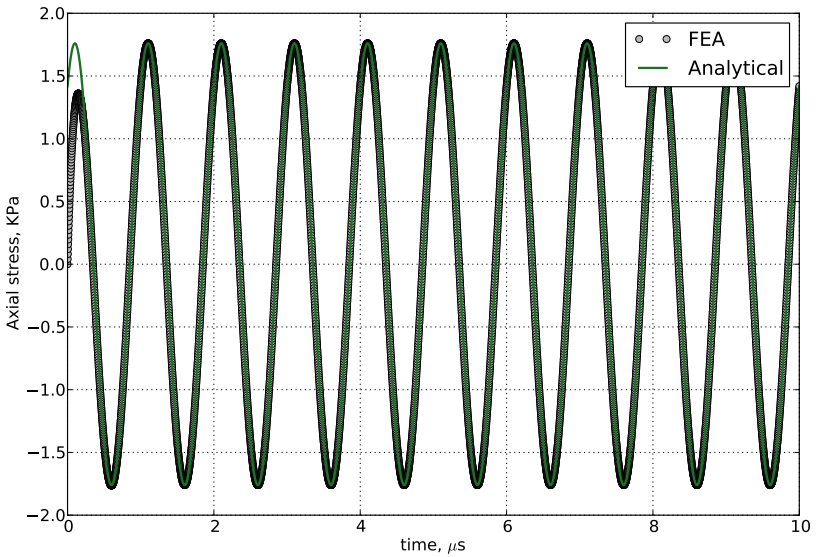
• Material model: example (FEA vs Analytics)

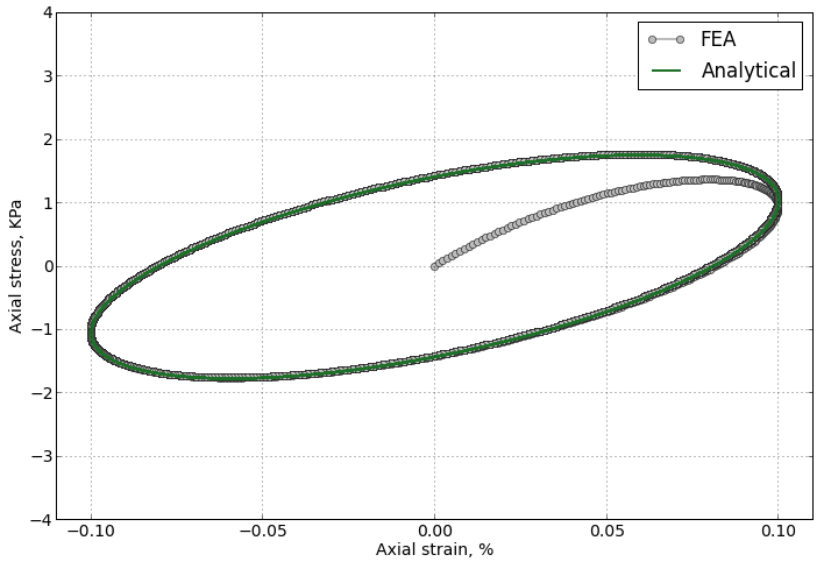




Linear frequency $f = 10^5$ Hz

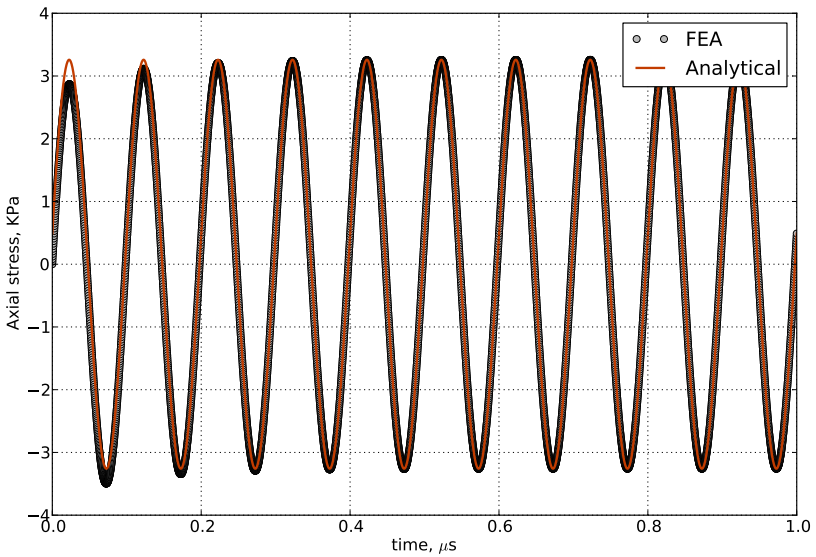
• Material model: example (FEA vs Analytics)

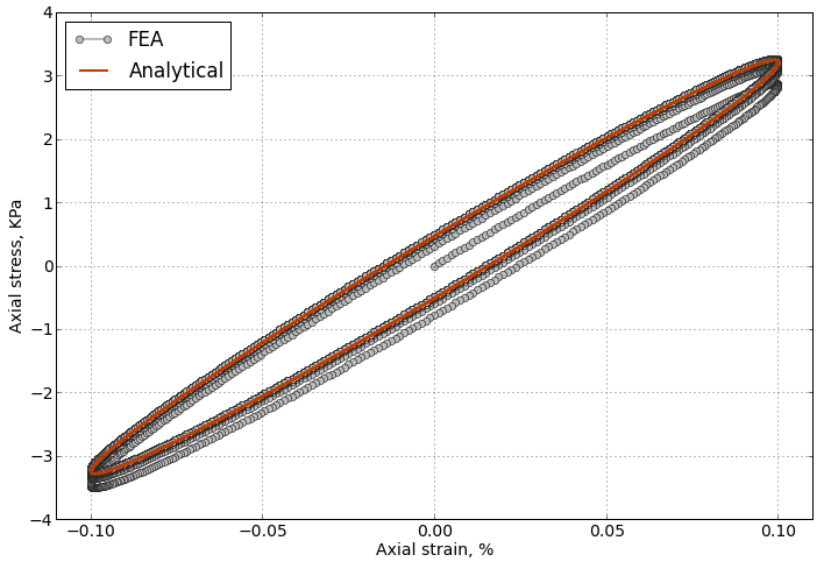




Linear frequency $f = 10^6$ Hz

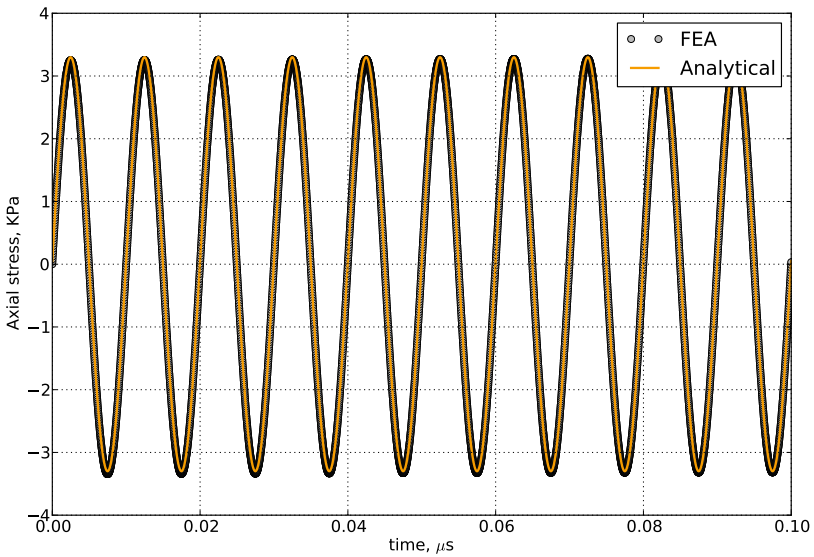
• Material model: example (FEA vs Analytics)

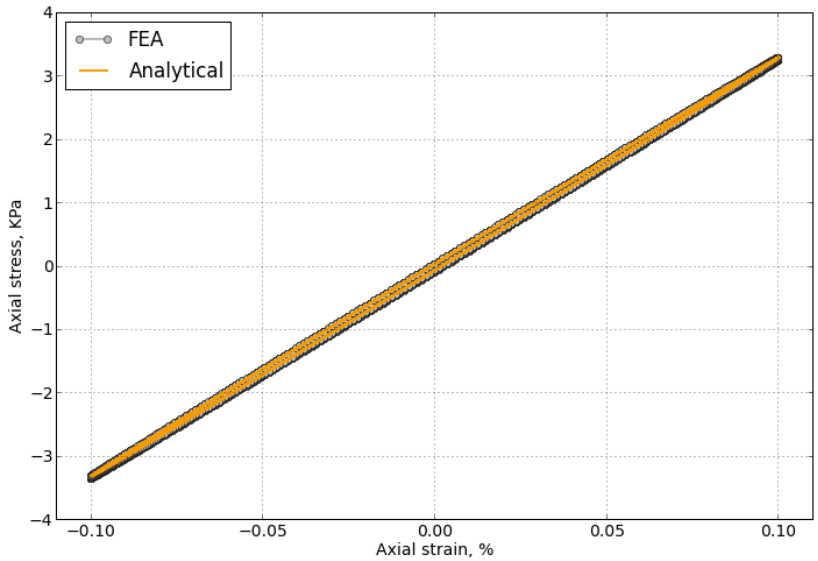




Linear frequency $f = 10^7$ Hz

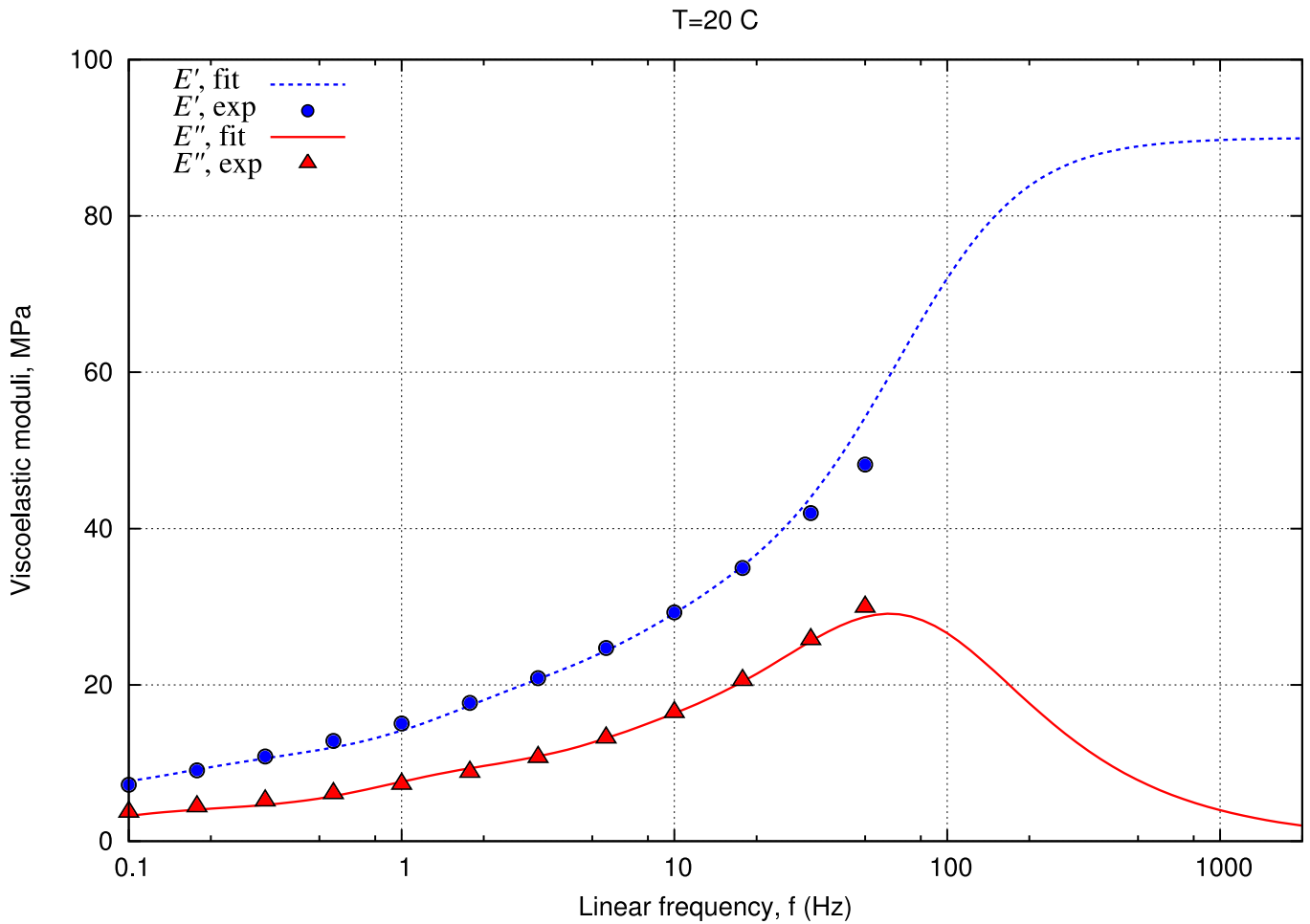
• Material model: example (FEA vs Analytics)





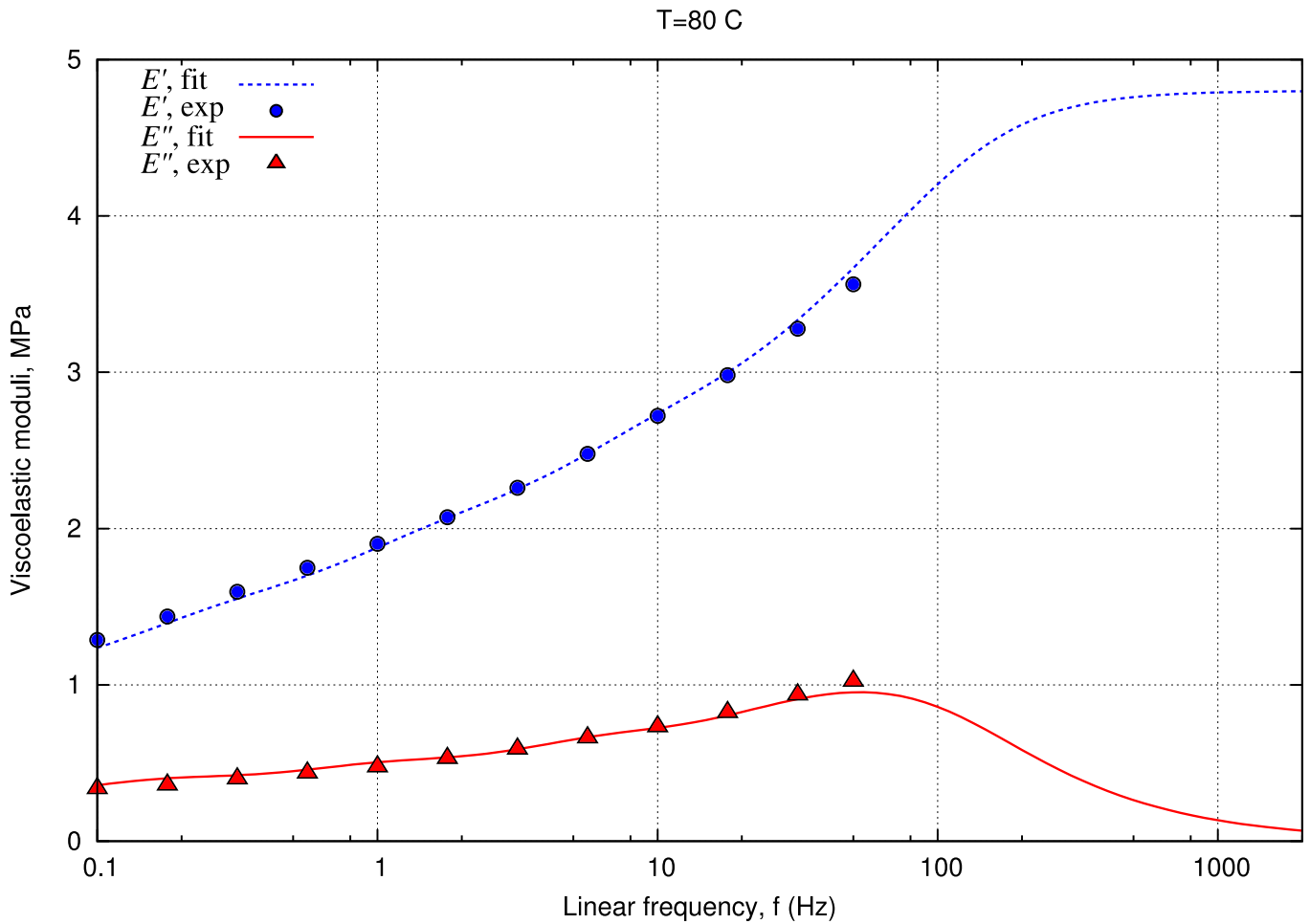
Linear frequency $f = 10^8$ Hz

Viscoelastic sliding: bulk friction



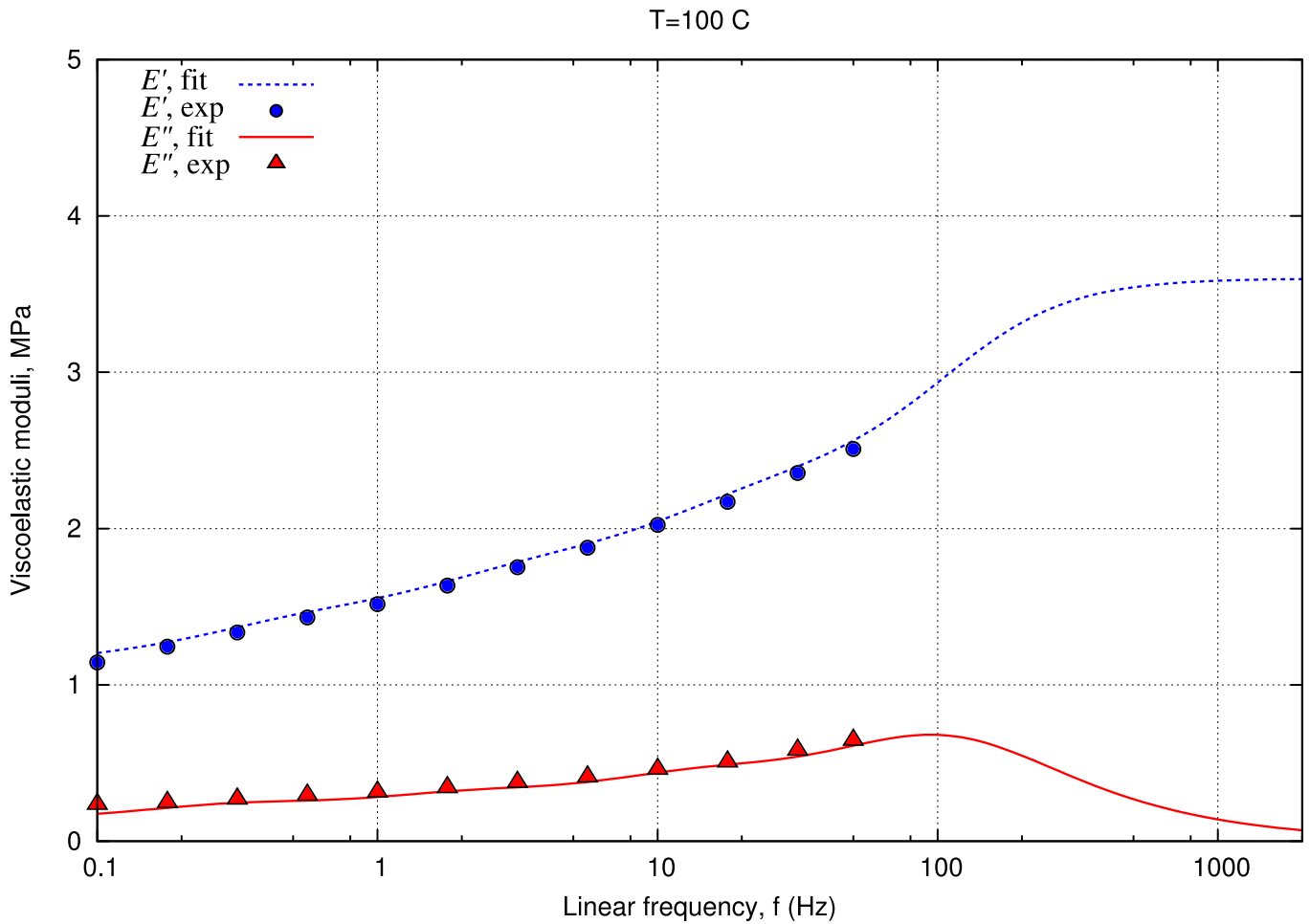
Fitting generalized Maxwell model for rubber to experimental data at $T = 20\ ^{o}C$

Viscoelastic sliding: bulk friction



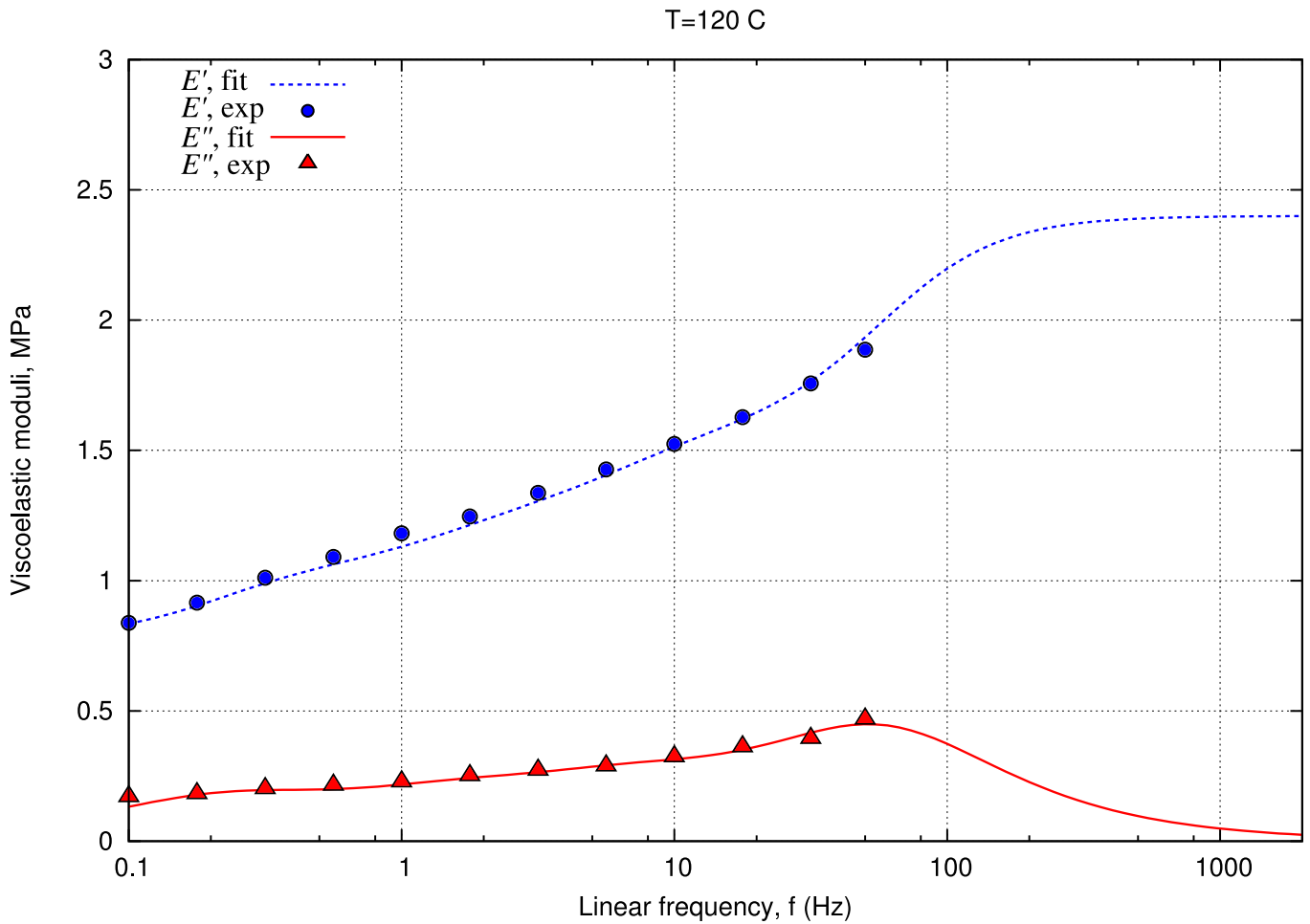
Fitting generalized Maxwell model for rubber to experimental data at $T = 80\ ^{o}C$

Viscoelastic sliding: bulk friction



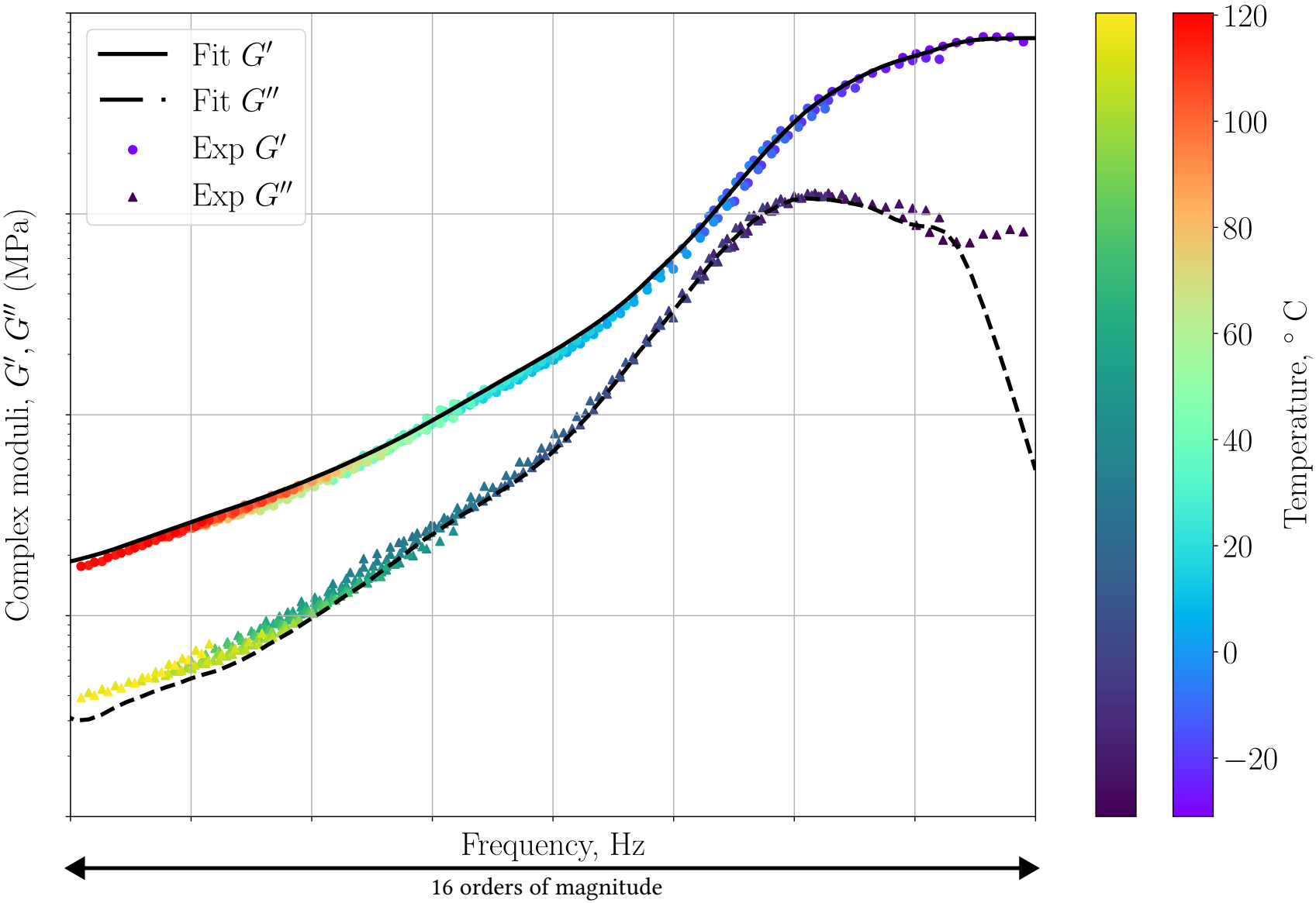
Fitting generalized Maxwell model for rubber to experimental data at $T = 100\ ^{o}C$

Viscoelastic sliding: bulk friction



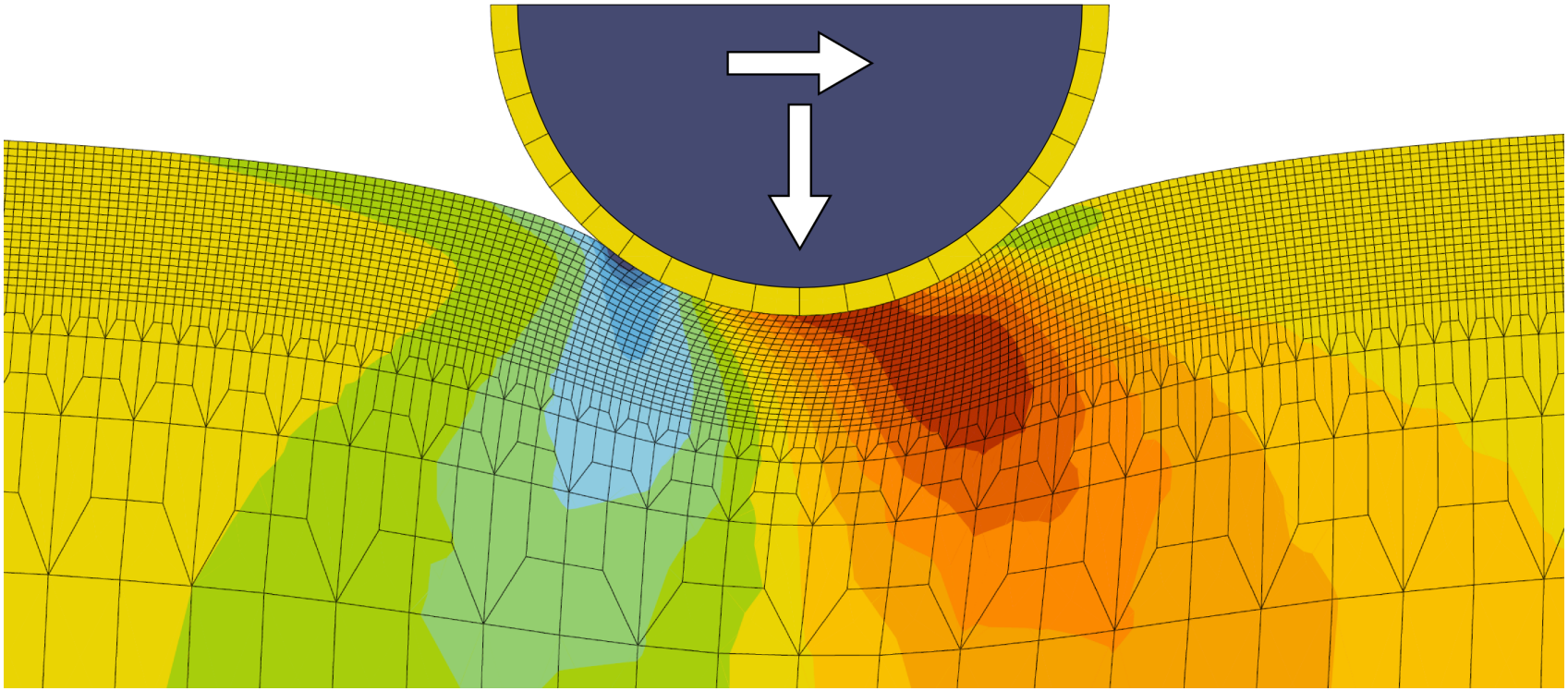
Fitting generalized Maxwell model for rubber to experimental data at $T = 120\ ^{o}C$

Viscoelastic sliding: bulk friction



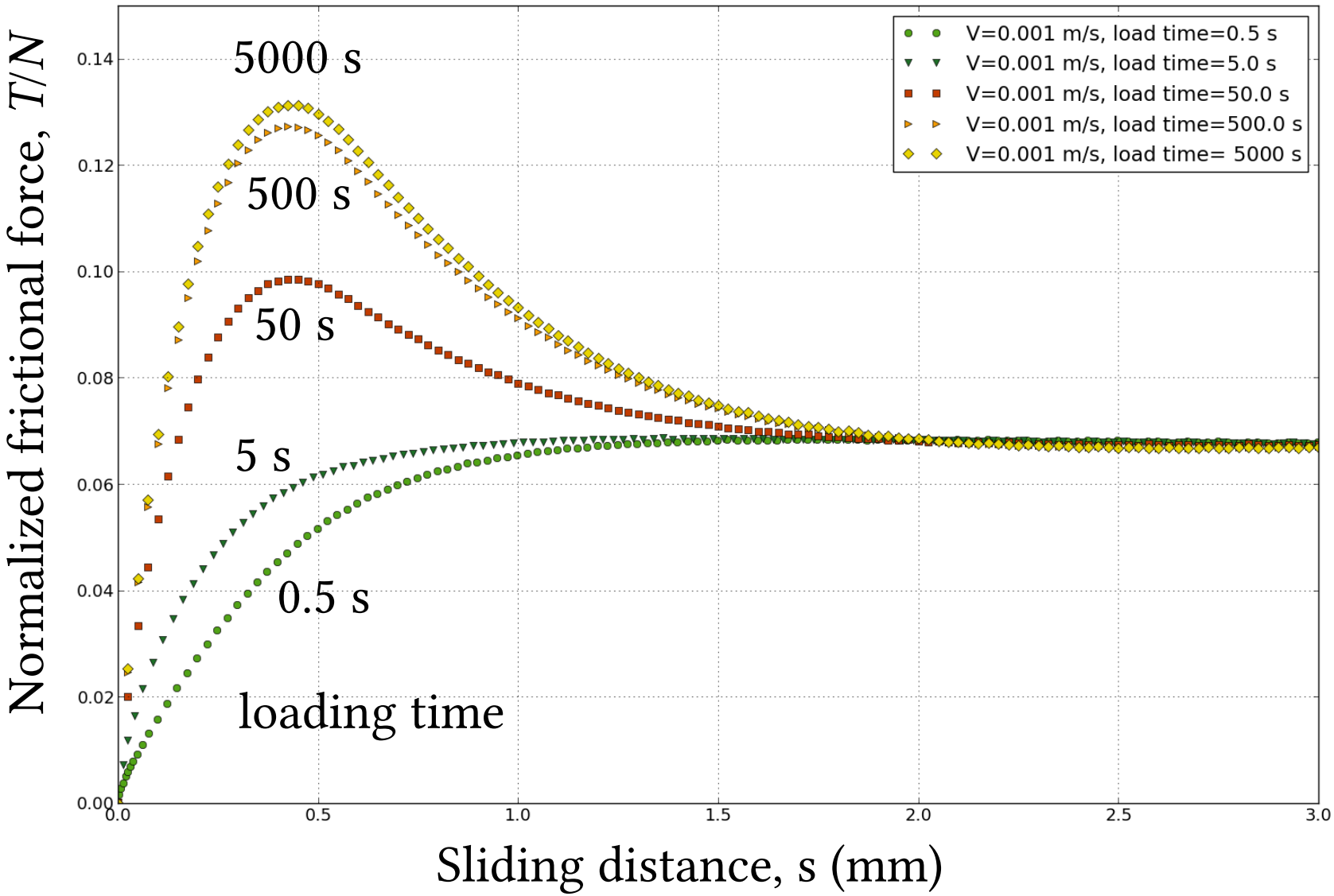
Full data on complex moduli of a sample charged rubber

Viscoelastic sliding: bulk friction



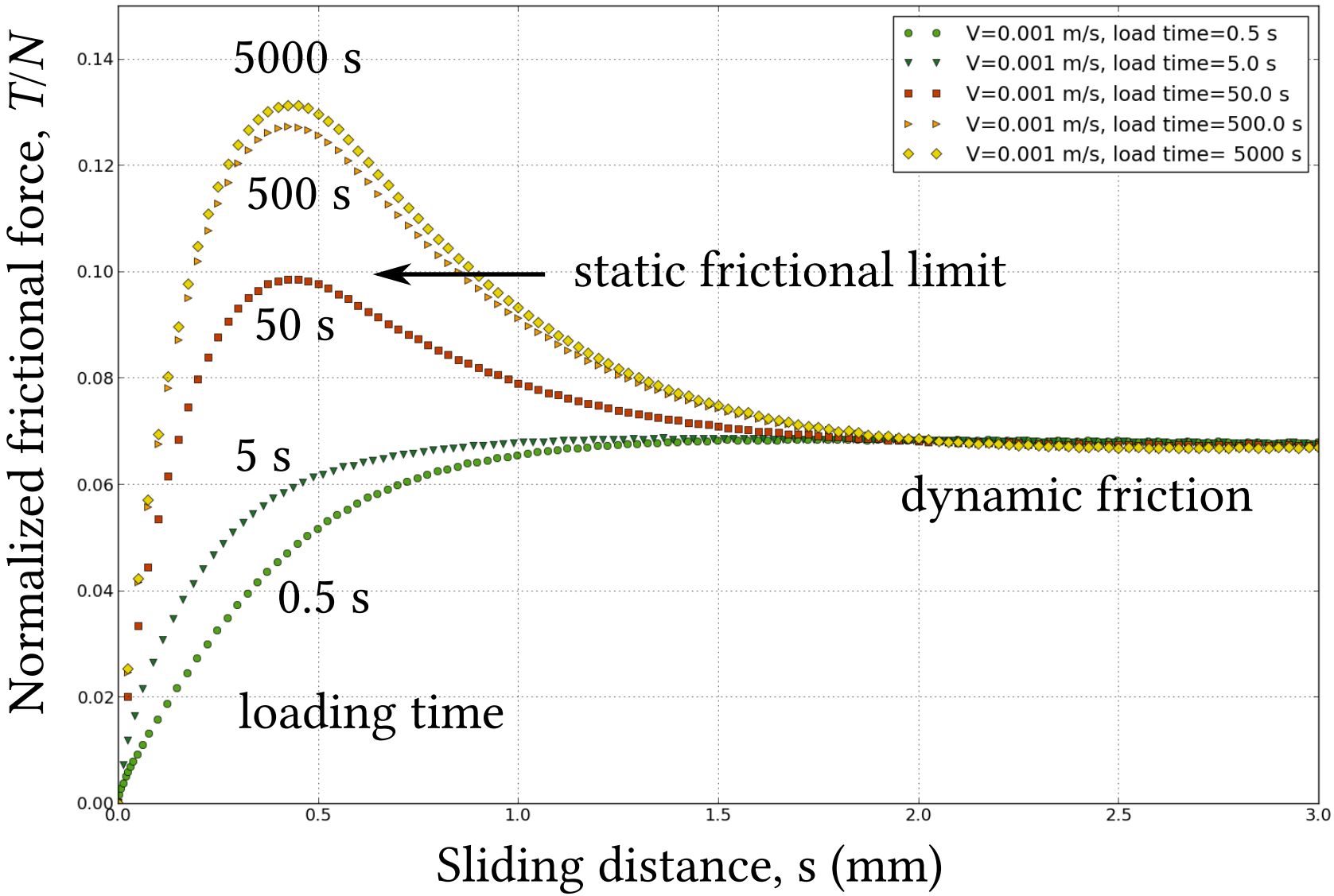
Simulation sketch

Viscoelastic sliding: bulk friction



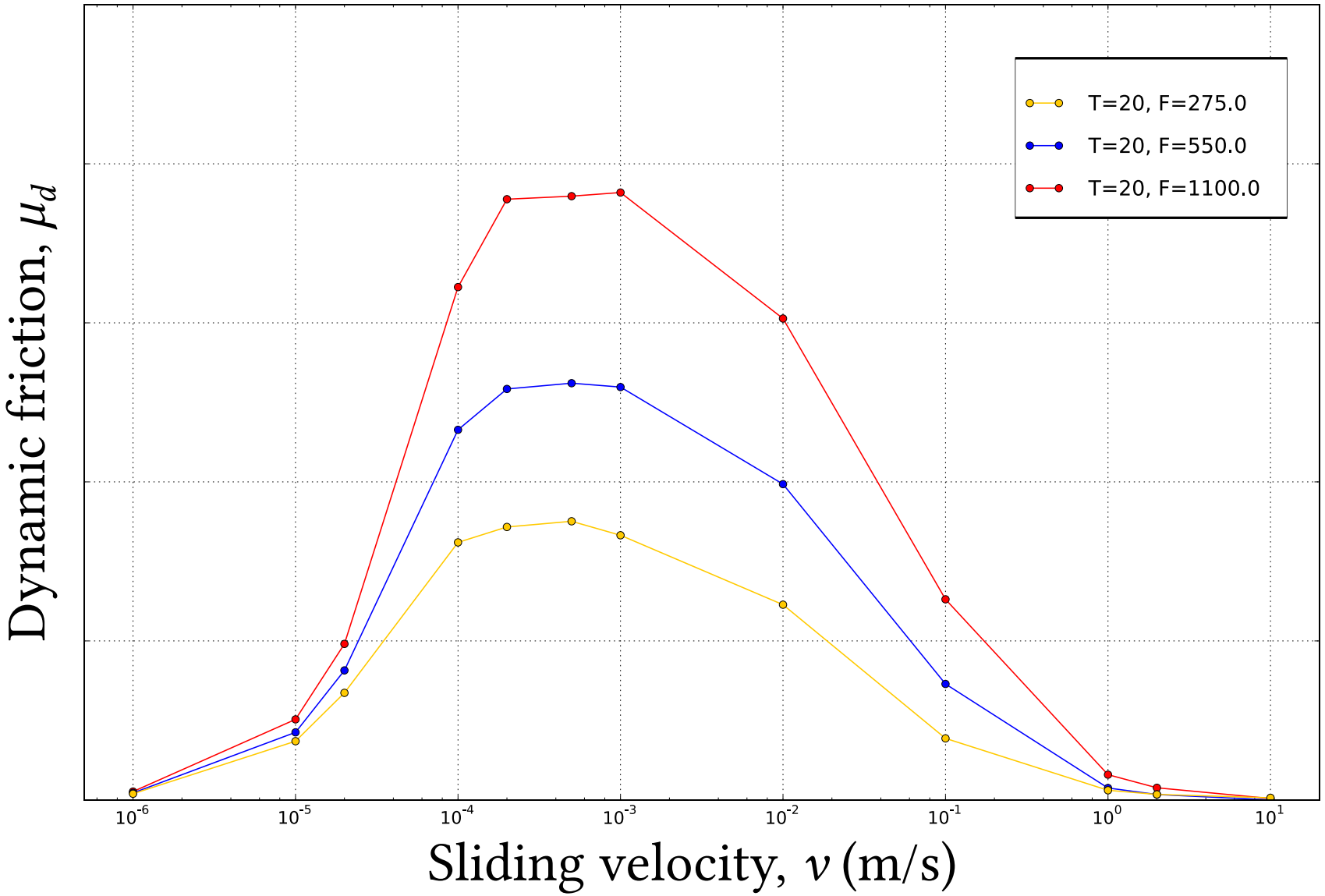
Effect of normal-loading rate on the frictional force evolution

Viscoelastic sliding: bulk friction



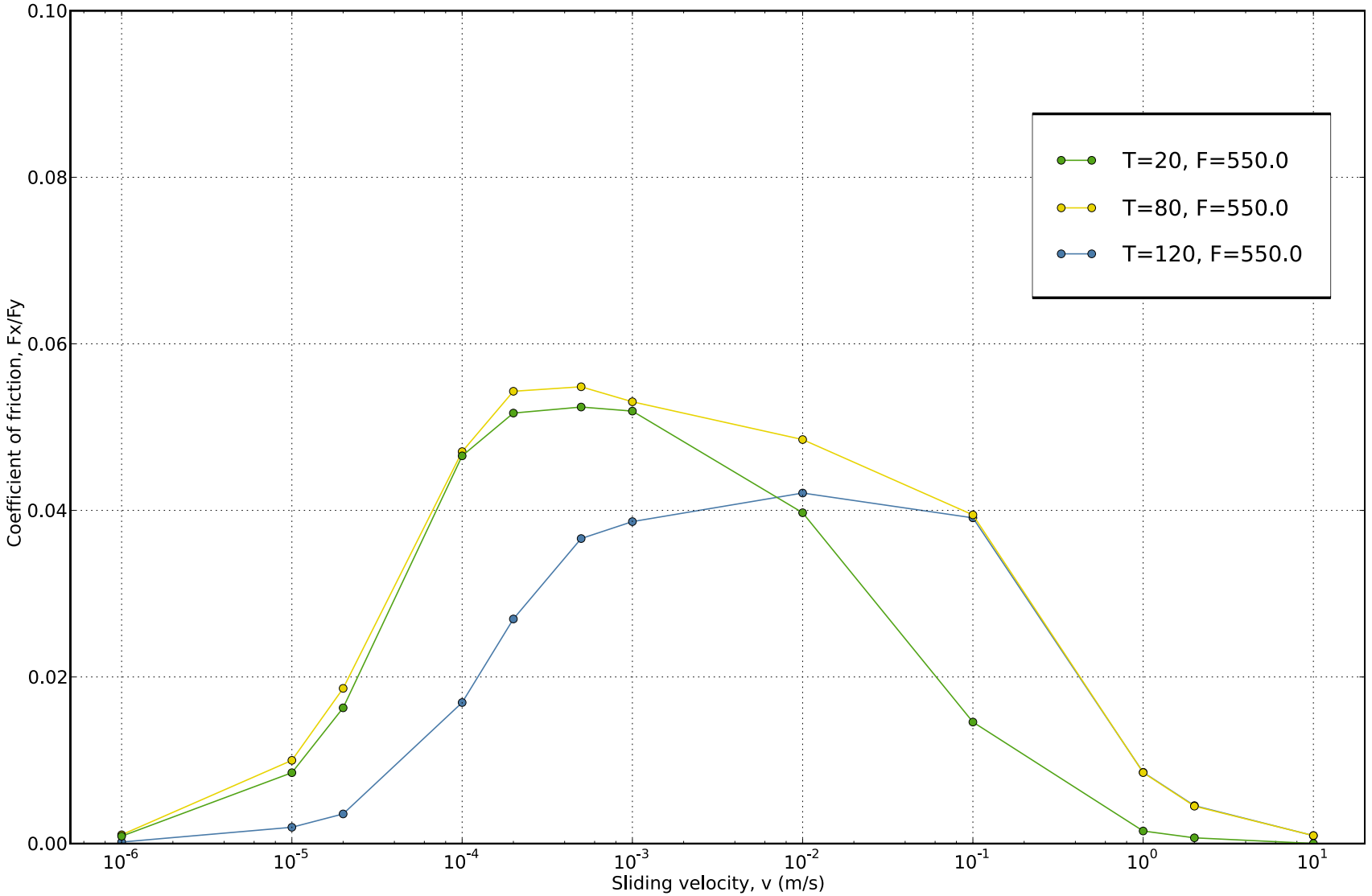
Effect of normal-loading rate on the frictional force evolution

Viscoelastic sliding: bulk friction



Frictional force at different slip velocity (effect of force)

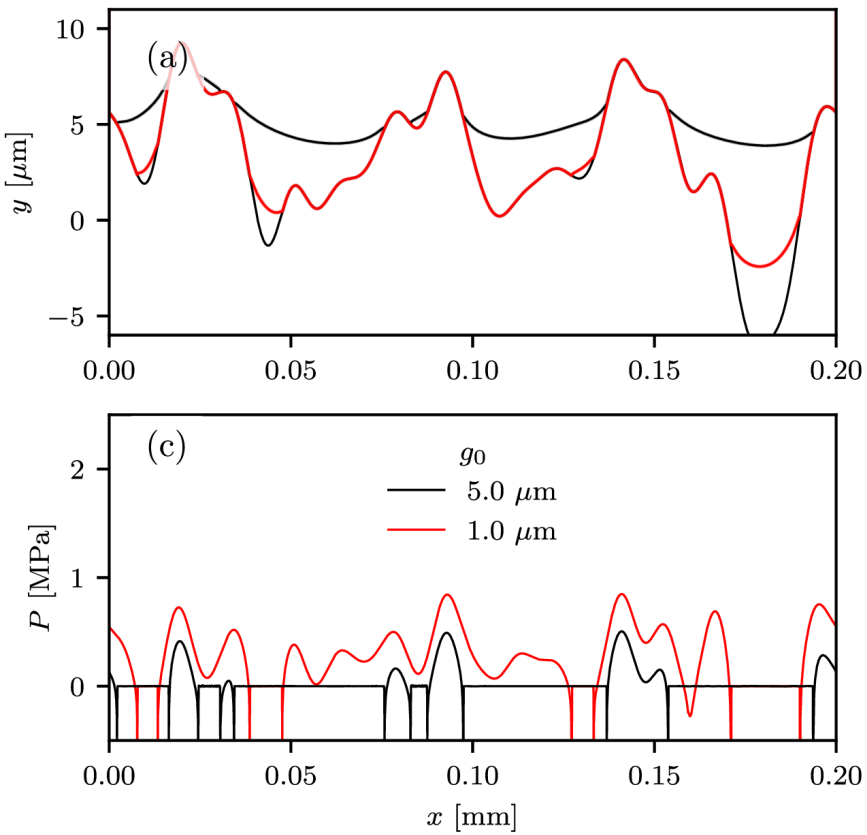
Viscoelastic sliding: bulk friction



Frictional force at different slip velocity (effect of temperature)

Viscoelastic friction: recent progress

- Account for multiscales[1,2]
- Account for adhesion[1,2]
- Account for large deformations[3]



Simulation results[1]

[1] Plagge, J. and Hentschke, R., (2022) Numerical solution of the adhesive rubber-solid contact problem and friction coefficients using a scale-splitting approach. Tribology International, 173:107622.

[2] Carbone, G., Mandriota, C. and Menga, N., (2022) Theory of viscoelastic adhesion and friction. Extreme Mechanics Letters, 56:101877.

[3] Lengiewicz, J., de Souza, M., Lahmar, M.A., Courbon, C., Dalmas, D., Stupkiewicz, S. and Scheibert, J., (2020) Finite deforma- tions govern the anisotropic shear-induced area reduction of soft elastic 2contacts. Journal of the Mechanics and Physics of Solids, 143:104056

Viscoelastic friction: recent progress

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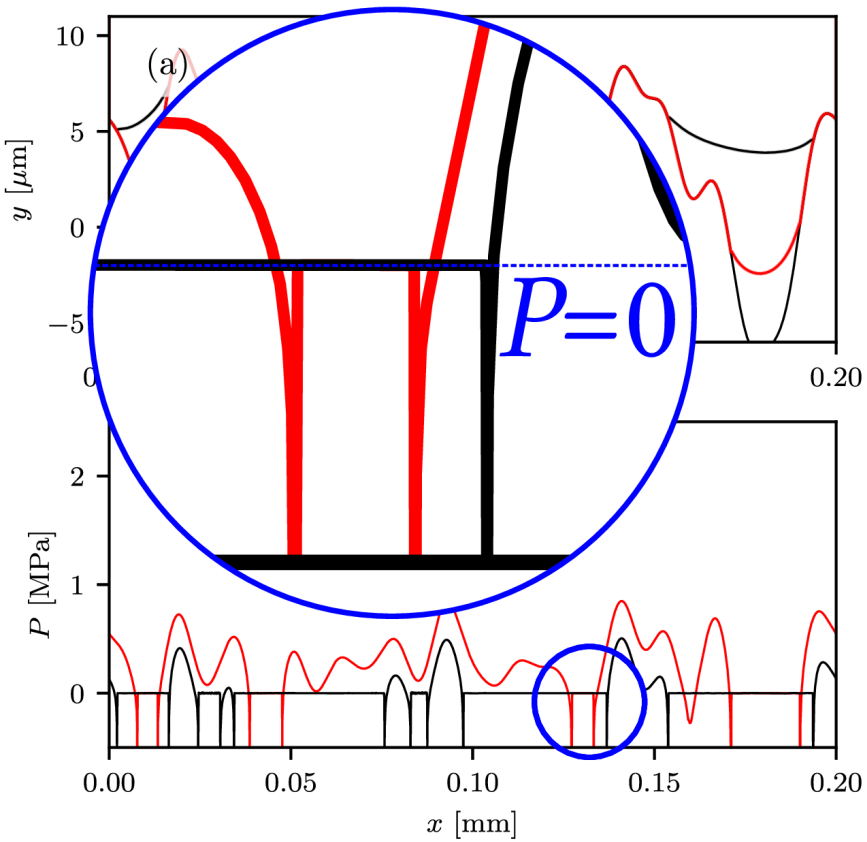

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Viscoelastic friction: recent progress

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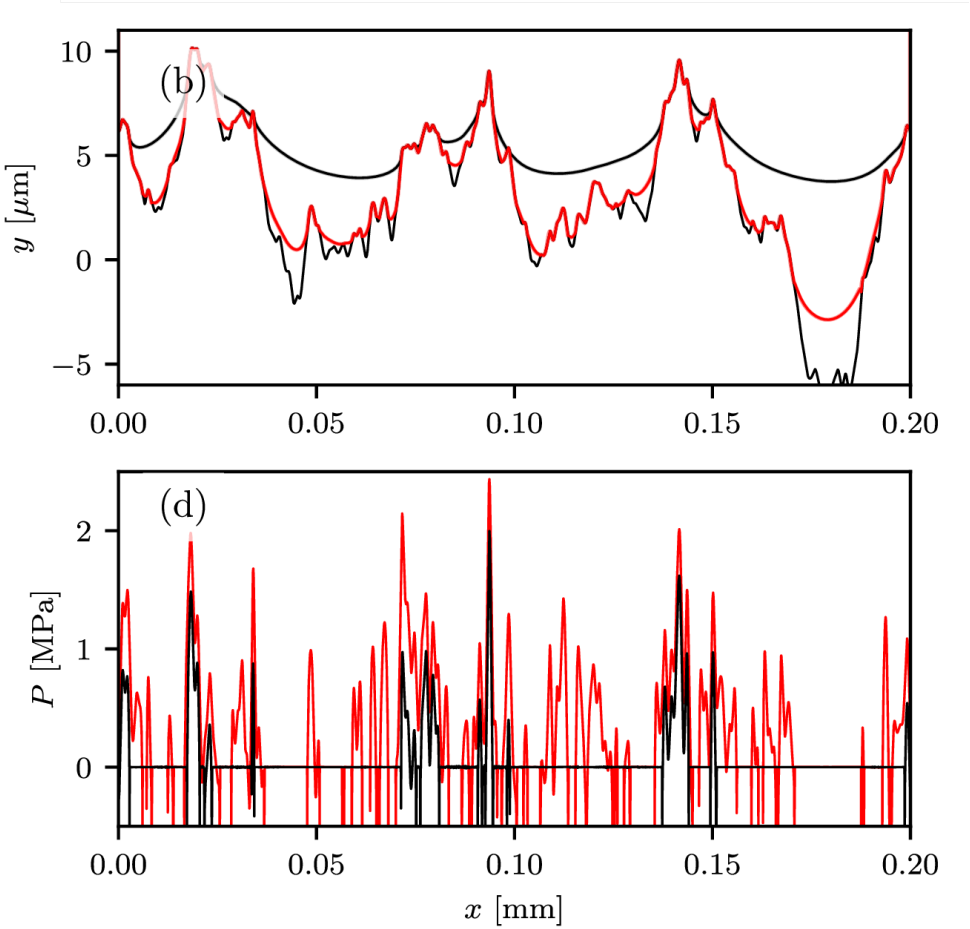

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Viscoelastic friction: recent progress

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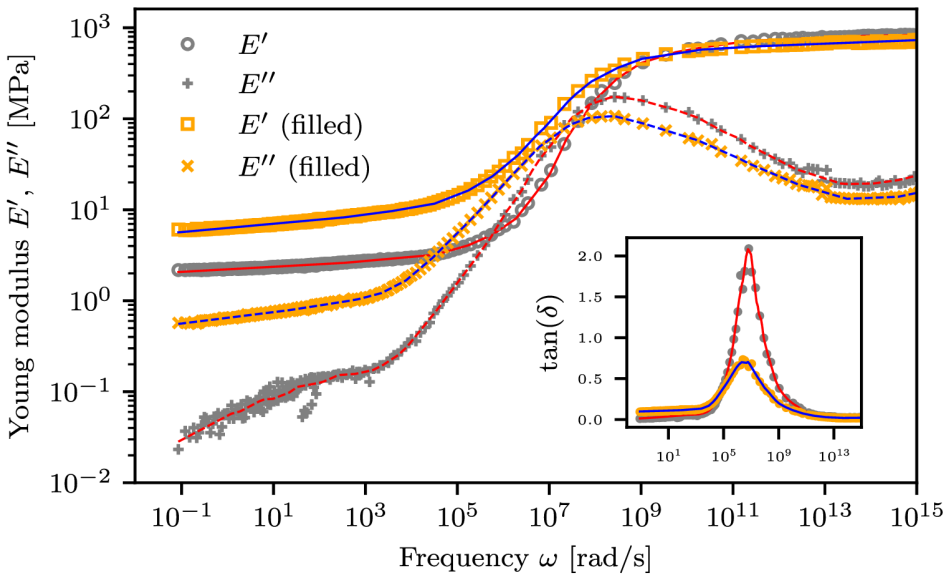

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Viscoelastic friction: recent progress

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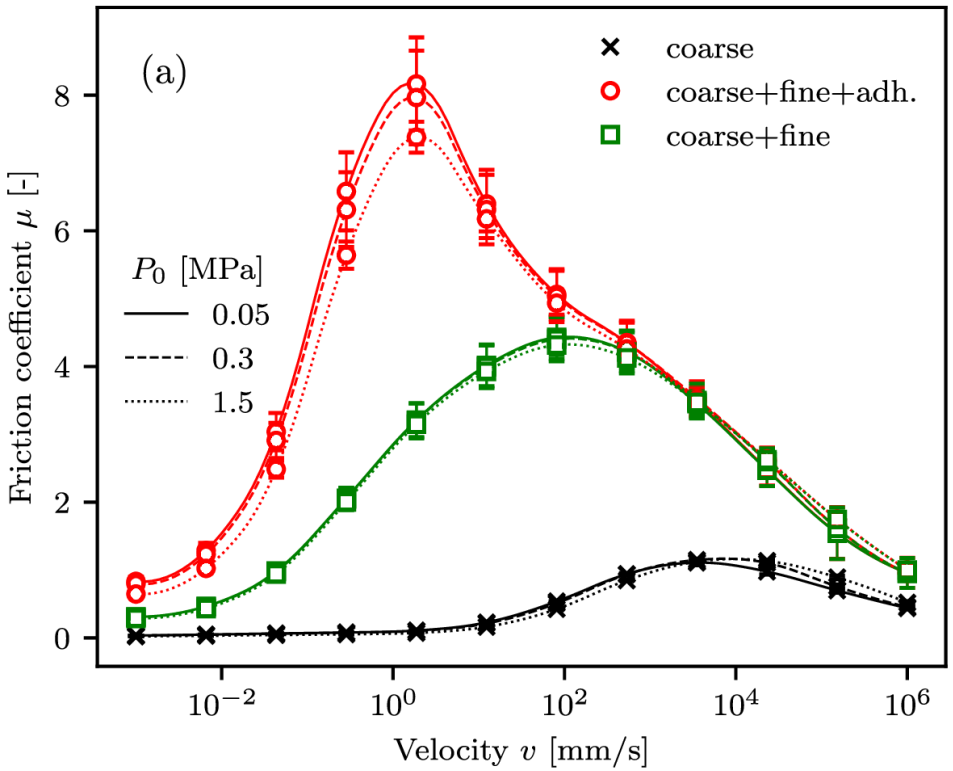

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Time-Dependent Friction

Check a separate presentation

Thank you for your attention!